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# Streams and caustics: the fine-grained structure of ACDM halos

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Aquarius-A-1 Springel et al 2008 Well *after* CDM particles become nonrelativistic, but *before* they dominate the cosmic density, their distribution function is

$$f(x, v, t) = \rho(t) [1 + \delta(x)] N [\{v - V(x)\}/\sigma]$$

where  $\rho(t)$  is the mean mass density of CDM,  $\delta(x)$  is a Gaussian random field with finite variance  $\ll 1$ ,  $V(x) = \nabla \psi(x)$  where  $\nabla^2 \psi(x) \propto \delta(x)$ and *N* is standard normal with  $\sigma^2 \ll \langle |\mathbf{V}|^2 \rangle$ 

CDM occupies a thin 3-D 'sheet' within the full 6-D phase-space and its projection onto x-space is near-uniform.

 $Df/Dt = 0 \longrightarrow$  only a 3-D subspace is occupied at later times. Nonlinear evolution leads to a complex, multi-stream structure.

# **Evolution of CDM structure**

Consequences of 
$$Df/Dt = 0$$

- The 3-D phase sheet can be stretched and folded but not torn
- At least 1 sheet must pass through every point **x**
- In nonlinear objects there are typically many sheets at each **x**
- Stretching which reduces a sheet's density must also reduce its velocity dispersions to maintain f = const.
- At a caustic, at least one velocity dispersion must  $\longrightarrow \infty$
- All these processes can be followed in fully general simulations by tracking the phase-sheet local to each simulation particle

# The geodesic deviation equation

Particle equation of motion: 
$$\dot{\mathbf{X}} = \begin{bmatrix} \mathbf{x} \\ \mathbf{\dot{v}} \end{bmatrix} = \begin{bmatrix} \mathbf{v} \\ -\nabla \phi \end{bmatrix}$$
  
Offset to a neighbor:  $\delta \dot{\mathbf{X}} = \begin{bmatrix} \delta \mathbf{v} \\ T \cdot \delta \mathbf{x} \end{bmatrix} = \begin{bmatrix} 0 & I \\ T & 0 \end{bmatrix} \cdot \delta \mathbf{X} ; T = -$ 

Write  $\delta X(t) = D(X_0, t) \cdot \delta X_0$ , then differentiating w.r.t. time gives,

$$\dot{\mathbf{D}} = \begin{bmatrix} 0 & \mathbf{I} \\ \mathbf{T} & \mathbf{0} \end{bmatrix} \cdot \mathbf{D} \text{ with } \mathbf{D}_0 = \mathbf{I}$$

- Integrating this equation together with each particle's trajectory gives the evolution of its local phase-space distribution
- No symmetry or stationarity assumptions are required
- det(D) = 1 at all times by Liouville's theorem
- For CDM,  $1/|det(D_{xx})|$  gives the decrease in local 3D space density of each particle's phase sheet. Switches sign and is infinite at caustics.



# Similarity solution for spherical collapse in CDM

#### Bertschinger 1985





## Simulation from self-similar spherical initial conditions

Geodesic deviation equation — phase-space structure local to each particle



### Simulation from self-similar spherical initial conditions



Vogelsberger et al 2009

The radial orbit instability leads to a system which is strongly prolate in the inner nonlinear regions

# Caustic crossing counts in a ACDM Milky Way halo



# Caustic crossing counts in a ACDM Milky Way halo



#### **Convergence of density profiles for Aquarius halos**



## **Caustic count profiles for Aquarius halos**



#### Stream density distribution in Aquarius halos



## Stream number profiles for Aquarius halos



#### Stream density distribution at the Sun

#### Vogelsberger & White 2010



Cumulative stream density distribution for particles with 7 kpc < r < 13 kpc

# Radial distribution of peak density at caustics

Vogelsberger & White 2010



Initial velocity dispersion assumes a standard WIMP with  $m = 100 \text{ GeV/c}^2$ 

## Fraction of annihilation luminosity from caustics



- Integration of the GDE can augment the ability of ΛCDM simulations to resolve fine-grained structure by 15 to 20 orders of magnitude
- Fine-grained streams and their associated caustics will have no significant effect on direct and indirect Dark Matter detection experiments
- The most massive stream at the Sun should contain roughly 0.001 of the local DM density and so might be detectable in an axion detector