

Cosmology with Large-scale Structure of the Universe

Eiichiro Komatsu (Texas Cosmology Center, UT Austin)
Korean Young Cosmologists Workshop, June 27, 2011

Cosmology Update: WMAP 7-year+

● Standard Model

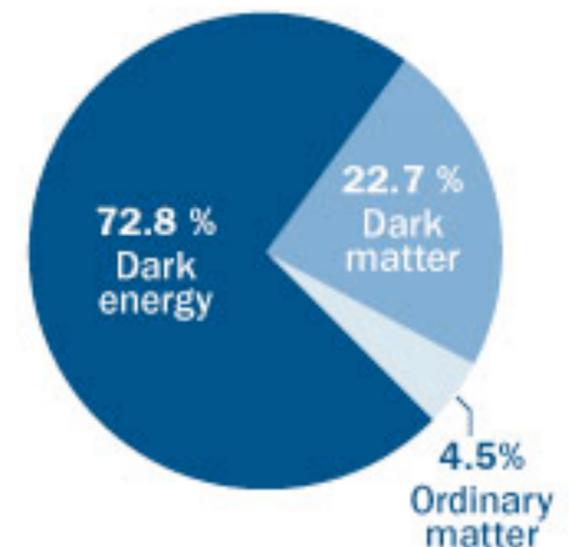
- H&He = 4.58% ($\pm 0.16\%$)
- **Dark Matter = 22.9%** ($\pm 1.5\%$)
- **Dark Energy = 72.5%** ($\pm 1.6\%$)
- $H_0 = 70.2 \pm 1.4$ km/s/Mpc
- Age of the Universe = 13.76 billion years (± 0.11 billion years)

Universal Stats

Age of the universe today
13.75 billion years

Age of the cosmos at
time of reionization
457 million years

Universe composition



*“ScienceNews” article on
the WMAP 7-year results*

Cosmology: Next Decade?

- Astro2010: Astronomy & Astrophysics Decadal Survey
 - Report from *Cosmology and Fundamental Physics* Panel (Panel Report, Page T-3):

TABLE I Summary of Science Frontiers Panels' Findings

Panel		Science Questions
Cosmology and Fundamental Physics	CFP 1	How Did the Universe Begin?
	CFP 2	Why Is the Universe Accelerating?
	CFP 3	What Is Dark Matter?
	CFP 4	What Are the Properties of Neutrinos?

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	CFP 2	Why Is the Universe Accelerating? <i>Dark Energy</i>
	CFP 3	What Is Dark Matter? <i>Dark Matter</i>
	CFP 4	What Are the Properties of Neutrinos? <i>Neutrino Mass</i>

Cosmology: Next Decade?

Large-scale structure of the universe has a potential to give us valuable information on all of these items.

Cosmology and
Fundamental Physics

CFP 1

How Did the Universe Begin *Inflation*

CFP 2

Why Is the Universe Accelerating? *Dark Energy*

CFP 3

What Is Dark Matter? *Dark Matter*

CFP 4

What Are the Properties of Neutrinos? *Neutrino Mass*

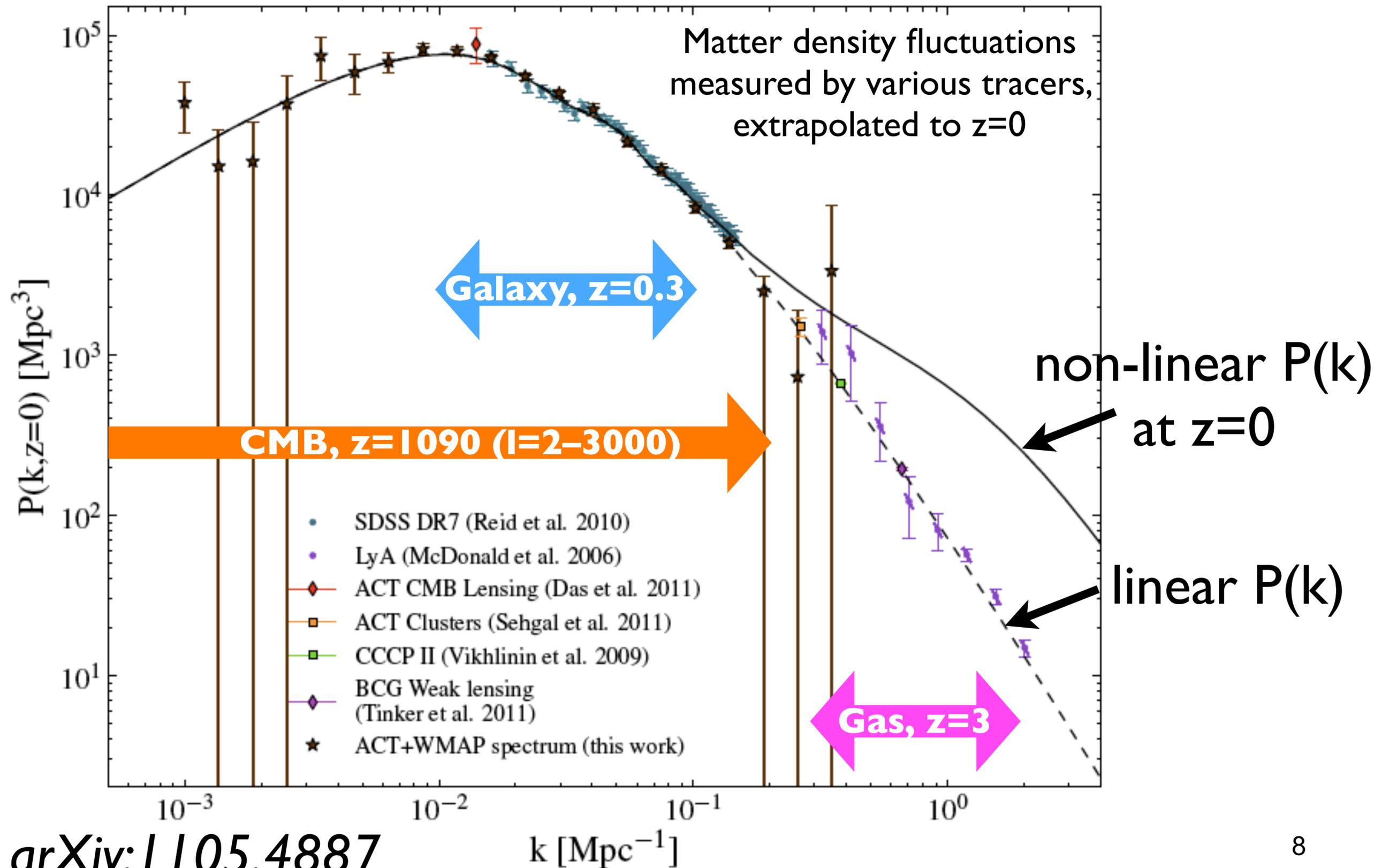
What to measure?

- **Inflation**
 - Shape of the initial power spectrum (n_s ; $dn_s/d\ln k$; etc)
 - Non-Gaussianity (3pt $f_{\text{NL}}^{\text{local}}$; 4pt $\tau_{\text{NL}}^{\text{local}}$; etc)
- **Dark Energy**
 - Angular diameter distances over a wide redshift range
 - Hubble expansion rates over a wide redshift range
 - Growth of linear density fluctuations over a wide redshift range
 - Shape of the matter power spectrum (modified grav)

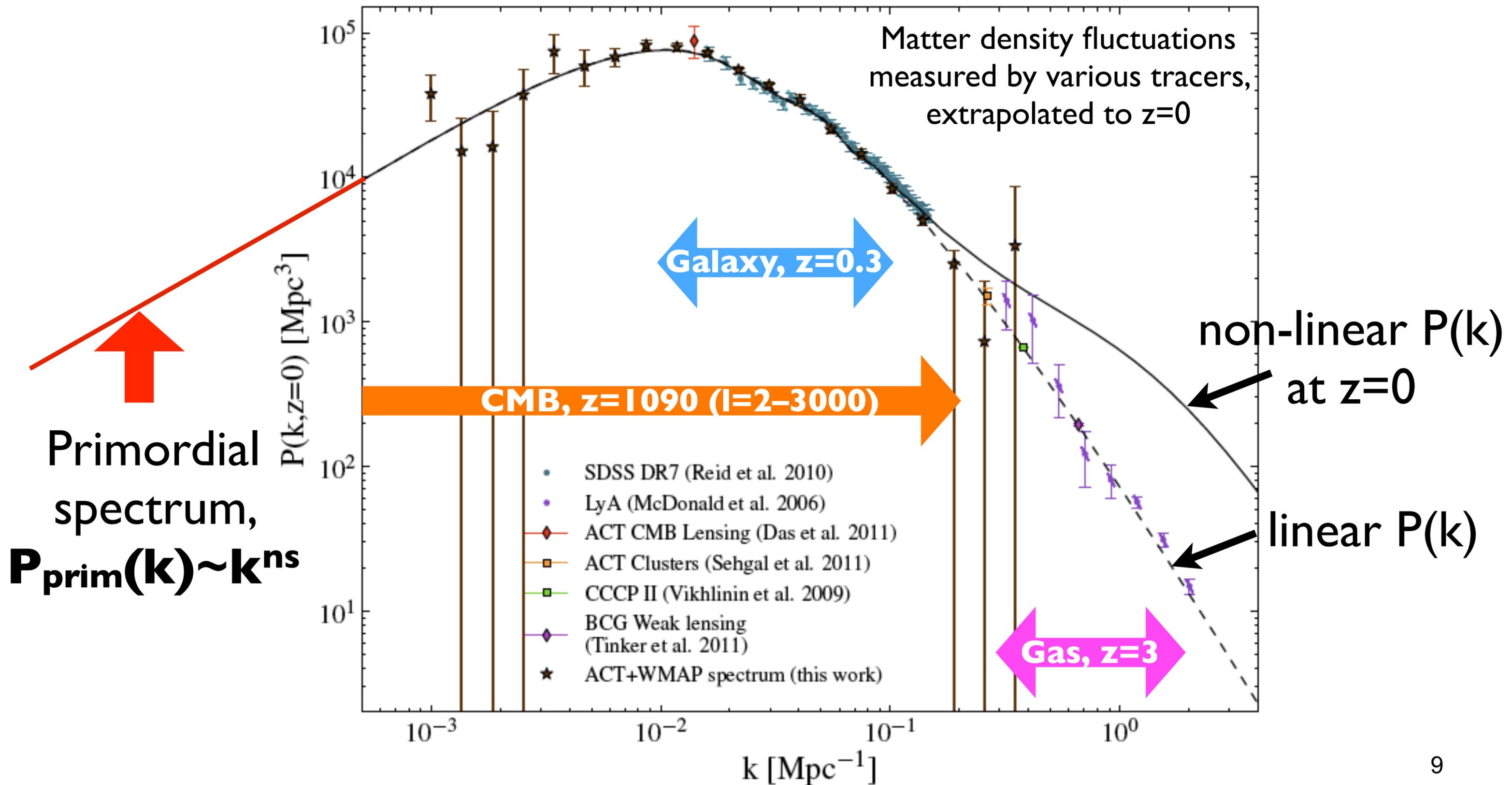
What to measure?

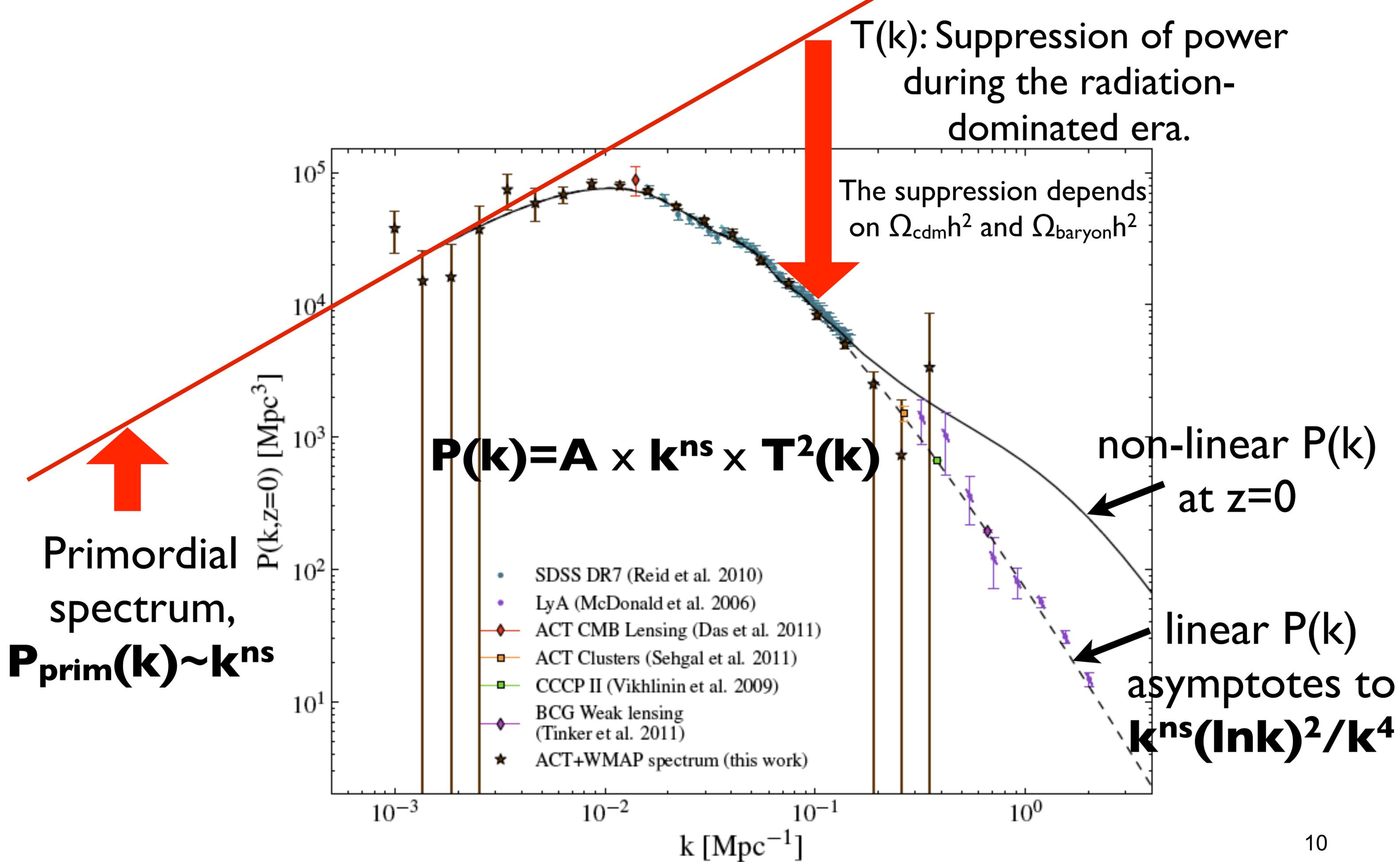
- **Neutrino Mass**
 - Shape of the matter power spectrum
- **Dark Matter**
 - Shape of the matter power spectrum (warm/hot DM)
 - Large-scale structure traced by γ -ray photons

Shape of the Power Spectrum, $P(k)$



Shape of the Power Spectrum, $P(k)$



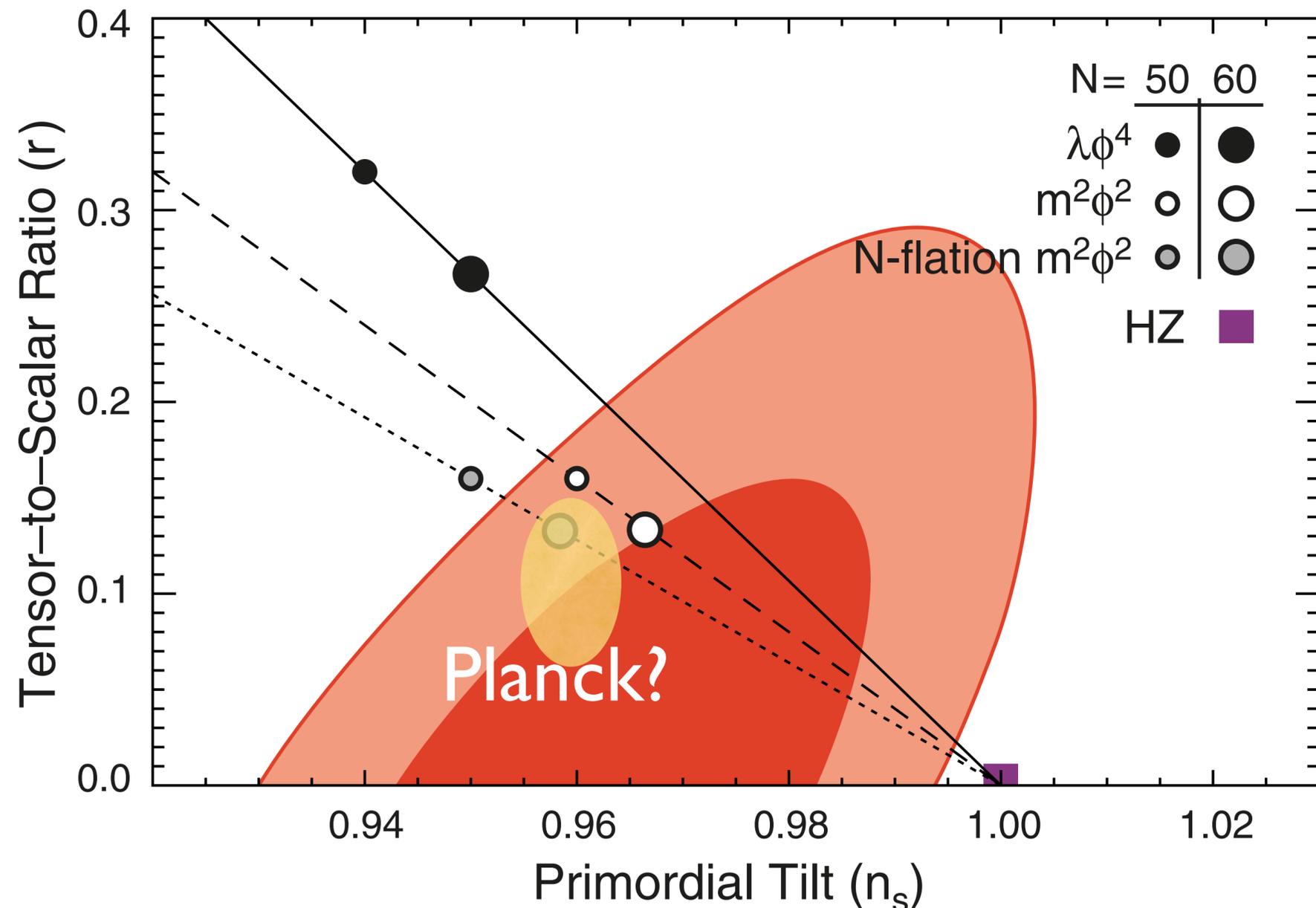


Current Limit on n_s

- Limit on the tilt of the power spectrum:
 - **$n_s = 0.968 \pm 0.012$** (68%CL; Komatsu et al. 2011)
 - Precision is dominated by the WMAP 7-year data
- Planck's CMB data are expected to improve the error bar by a factor of ~ 4 .

Probing Inflation (2-point Function)

$$r = (\text{gravitational waves})^2 / (\text{gravitational potential})^2$$

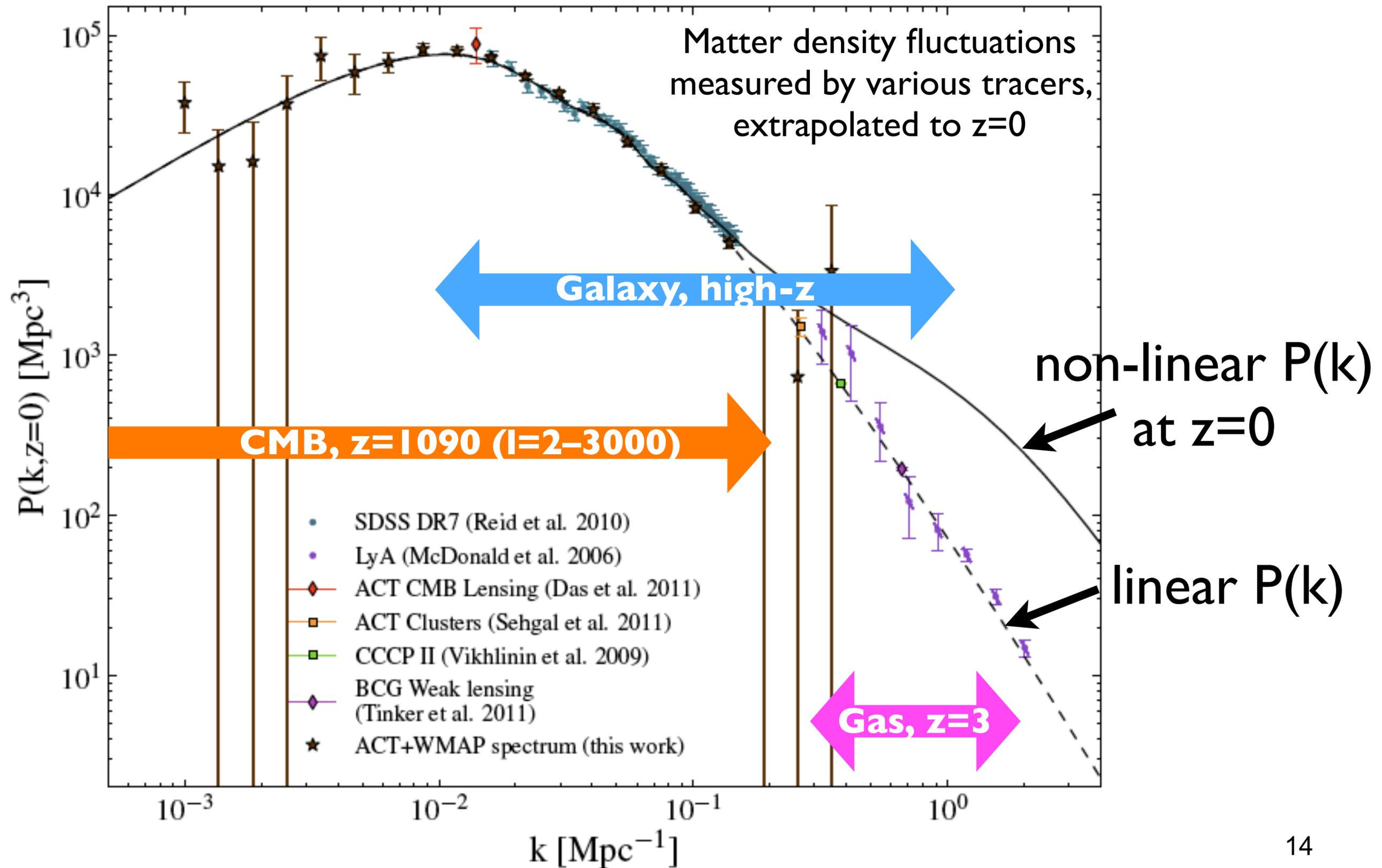


- Joint constraint on the primordial tilt, n_s , and the tensor-to-scalar ratio, r .
- Not so different from the 5-year limit.
- $r < 0.24$ (95%CL)
- Limit on the tilt of the power spectrum:
 $n_s = 0.968 \pm 0.012$ (68%CL)

Role of the Large-scale Structure of the Universe

- However, CMB data can't go much beyond $k=0.2 \text{ Mpc}^{-1}$ ($l=3000$).
- Large-scale structure data are required to go to smaller scales.

Shape of the Power Spectrum, $P(k)$



Measuring a scale-dependence of $n_s(k)$

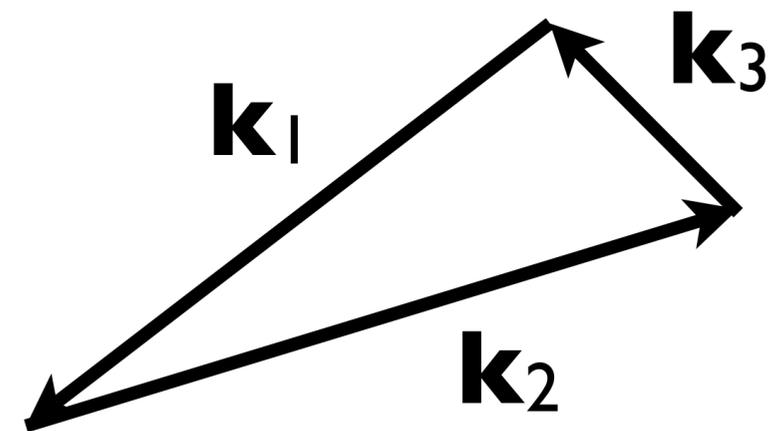
- As far as the value of n_s is concerned, CMB is probably enough.
- However, if we want to measure the scale-dependence of n_s , i.e., deviation of $P_{\text{prim}}(k)$ from a pure power-law, then we need the small-scale data.
 - This is where the large-scale structure data become quite powerful (Takada, Komatsu & Futamase 2006)
- Schematically:
 - $dn_s/d\ln k = [n_s(\text{CMB}) - n_s(\text{LSS})]/(\ln k_{\text{CMB}} - \ln k_{\text{LSS}})$

Probing Inflation (3-point Function)

Can We Rule Out Inflation?

- Inflation models predict that primordial fluctuations are very close to Gaussian.
- In fact, **ALL SINGLE-FIELD** models predict a particular form of **3-point function** to have the amplitude of $f_{\text{NL}}^{\text{local}}=0.02$.
- Detection of $f_{\text{NL}} > 1$ would rule out ALL single-field models!

Bispectrum



- Three-point function!

- $B_\zeta(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3)$

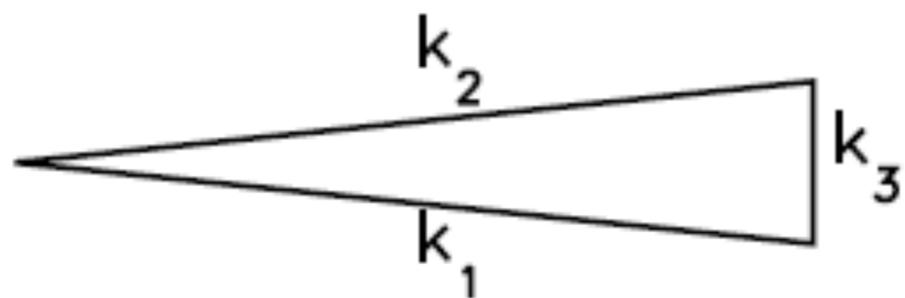
$$= \langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \zeta_{\mathbf{k}_3} \rangle = (\text{amplitude}) \times (2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) F(k_1, k_2, k_3)$$

model-dependent function

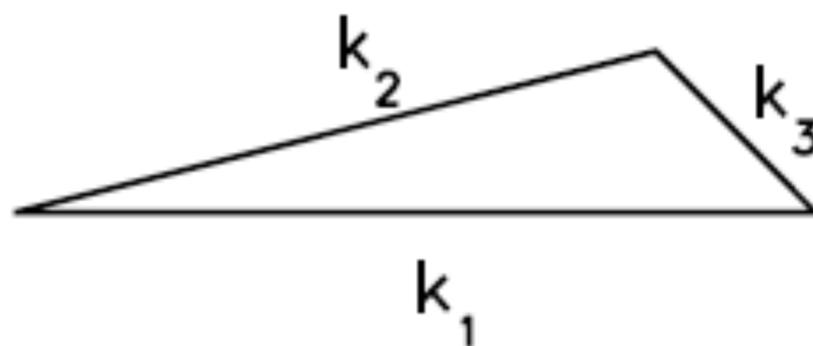


Primordial fluctuation

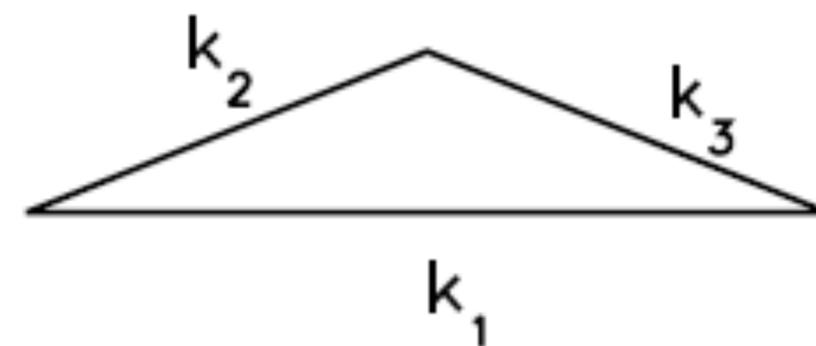
(a) squeezed triangle
($k_1 \approx k_2 \gg k_3$)



(b) elongated triangle
($k_1 = k_2 + k_3$)

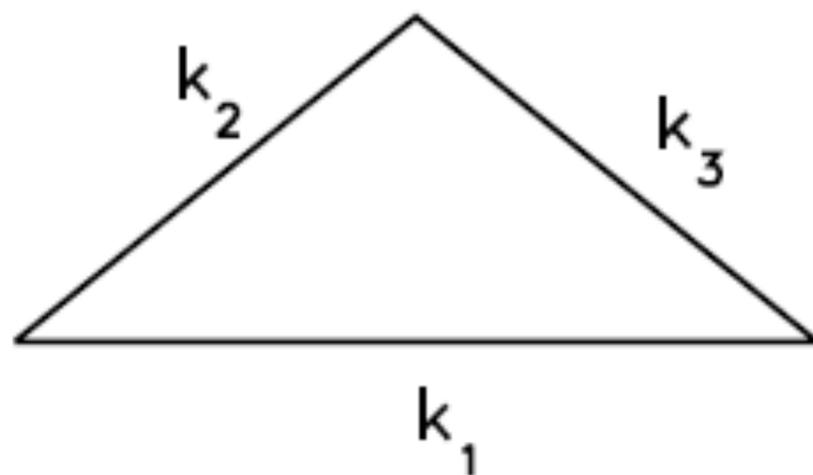


(c) folded triangle
($k_1 = 2k_2 = 2k_3$)

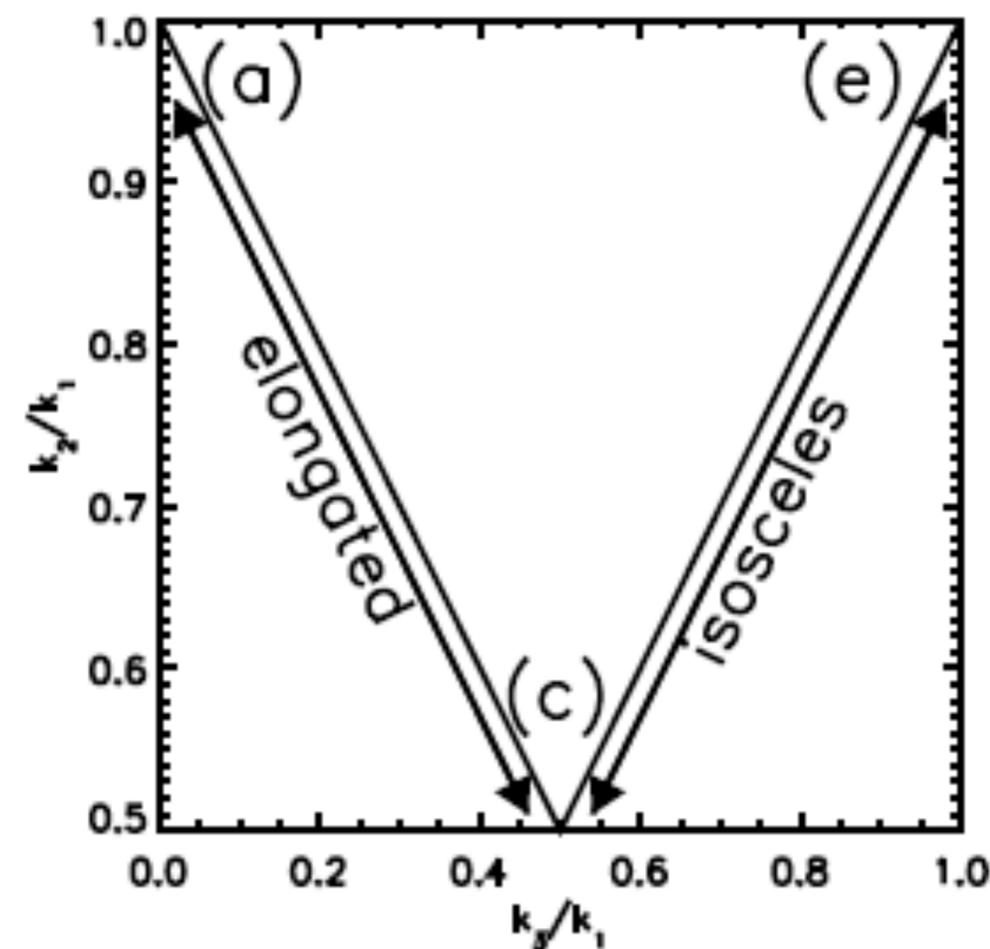
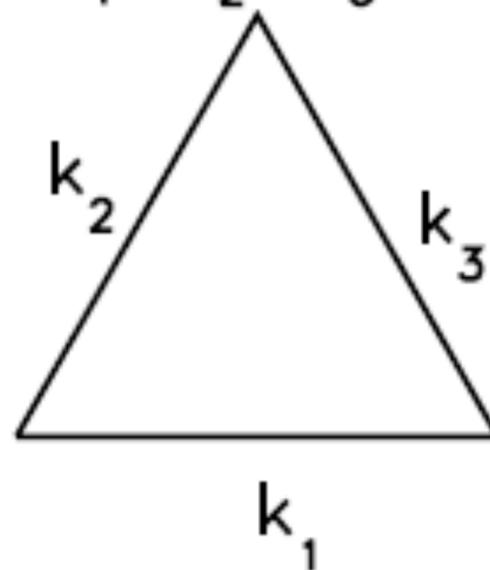


MOST IMPORTANT

(d) isosceles triangle
($k_1 > k_2 = k_3$)



(e) equilateral triangle
($k_1 = k_2 = k_3$)



Single-field Theorem (Consistency Relation)

- For **ANY** single-field models*, the bispectrum in the squeezed limit is given by
- $B_{\zeta}(\mathbf{k}_1 \sim \mathbf{k}_2 \ll \mathbf{k}_3) \approx (1 - n_s) \times (2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) \times P_{\zeta}(k_1) P_{\zeta}(k_3)$
- Therefore, all single-field models predict $f_{\text{NL}} \approx (5/12)(1 - n_s)$.
- With the current limit $n_s = 0.968$, f_{NL} is predicted to be 0.01.

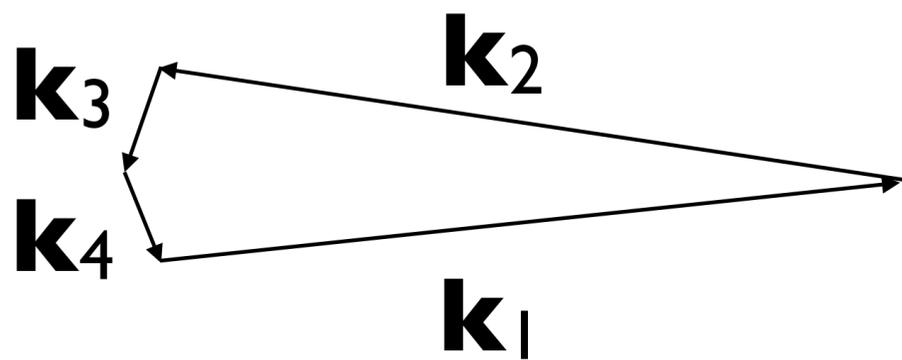
* for which the single field is solely responsible for driving inflation and generating observed fluctuations.

Probing Inflation (3-point Function)

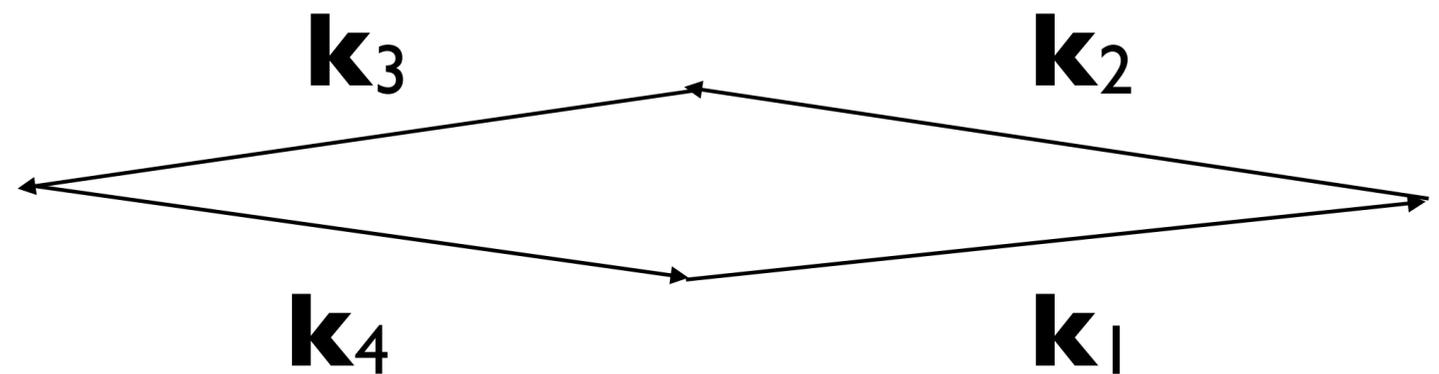
- No detection of 3-point functions of primordial curvature perturbations. The 95% CL limit is:
 - $-10 < f_{\text{NL}}^{\text{local}} < 74$
- The 68% CL limit: $f_{\text{NL}}^{\text{local}} = 32 \pm 21$
 - The WMAP data are consistent with the prediction of **simple single-field inflation** models: $1 - n_s \approx r \approx f_{\text{NL}}$
- The Planck's expected 68% CL uncertainty: $\Delta f_{\text{NL}}^{\text{local}} = 5$

Trispectrum

- $T_{\zeta}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3, \mathbf{k}_4) = (2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4) \{ g_{NL} [(54/25) P_{\zeta}(k_1) P_{\zeta}(k_2) P_{\zeta}(k_3) + \text{cyc.}] + T_{NL} [P_{\zeta}(k_1) P_{\zeta}(k_2) (P_{\zeta}(|\mathbf{k}_1 + \mathbf{k}_3|) + P_{\zeta}(|\mathbf{k}_1 + \mathbf{k}_4|)) + \text{cyc.}] \}$

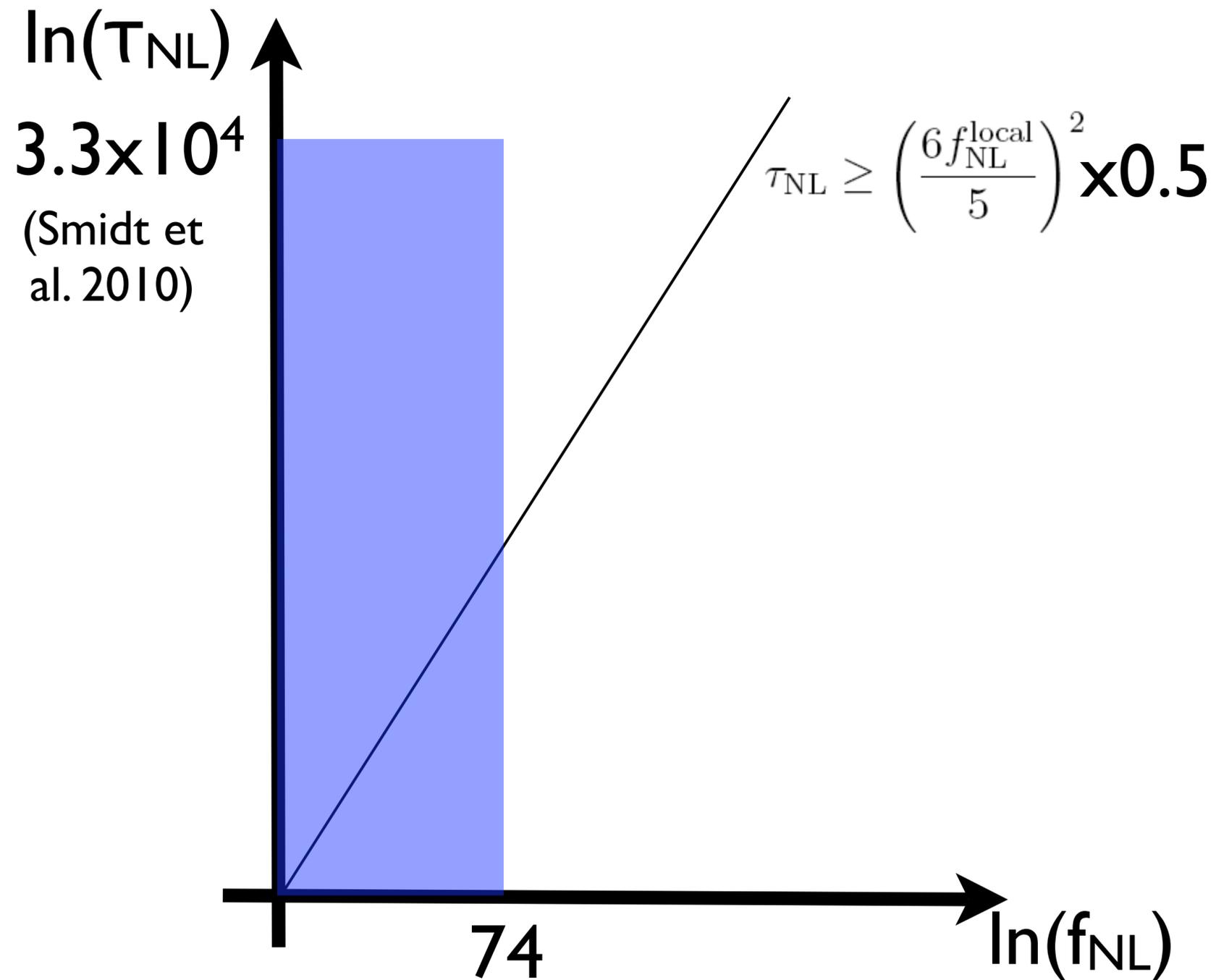


g_{NL}



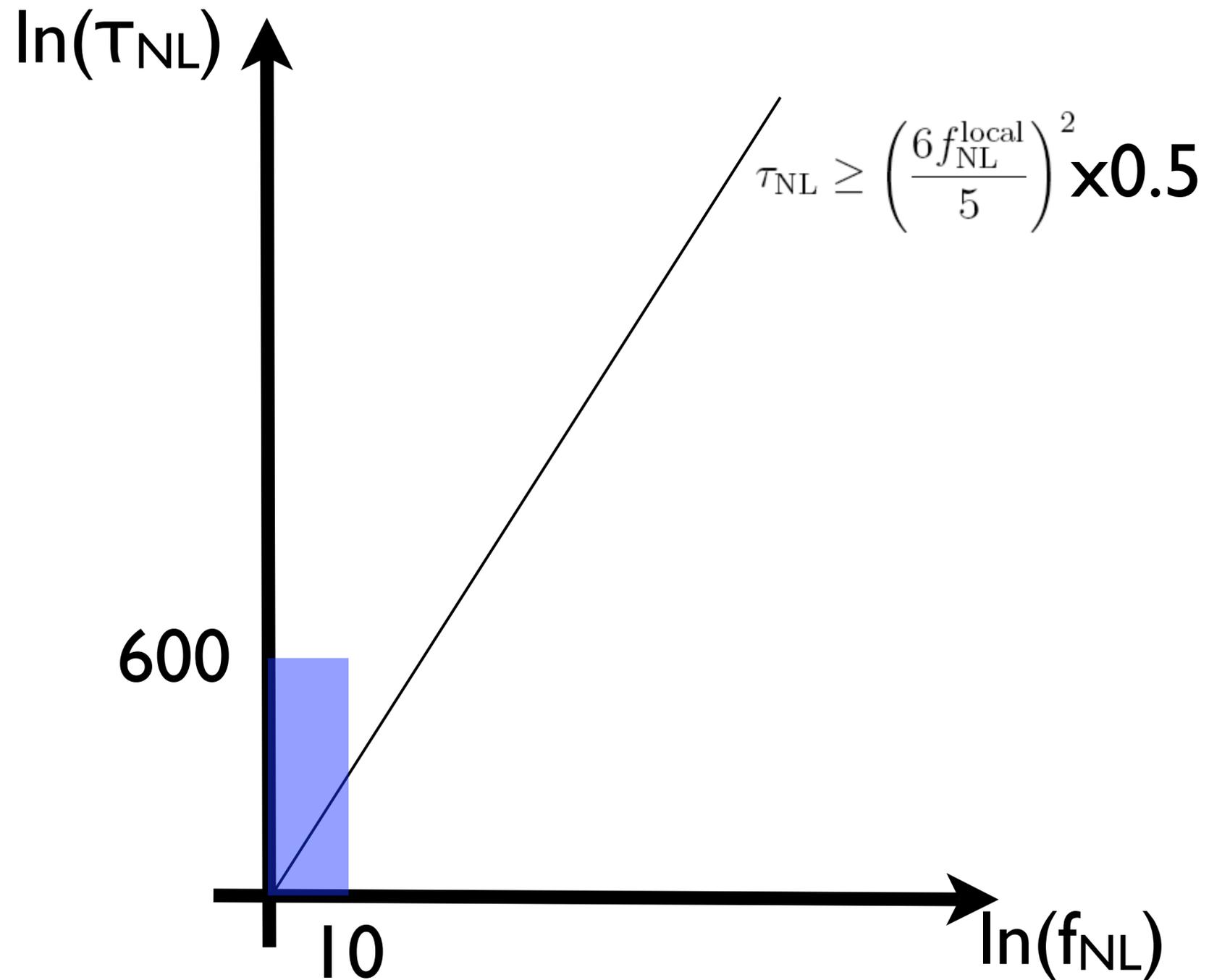
T_{NL}

$\tau_{\text{NL}}^{\text{local}} - f_{\text{NL}}^{\text{local}}$ Diagram



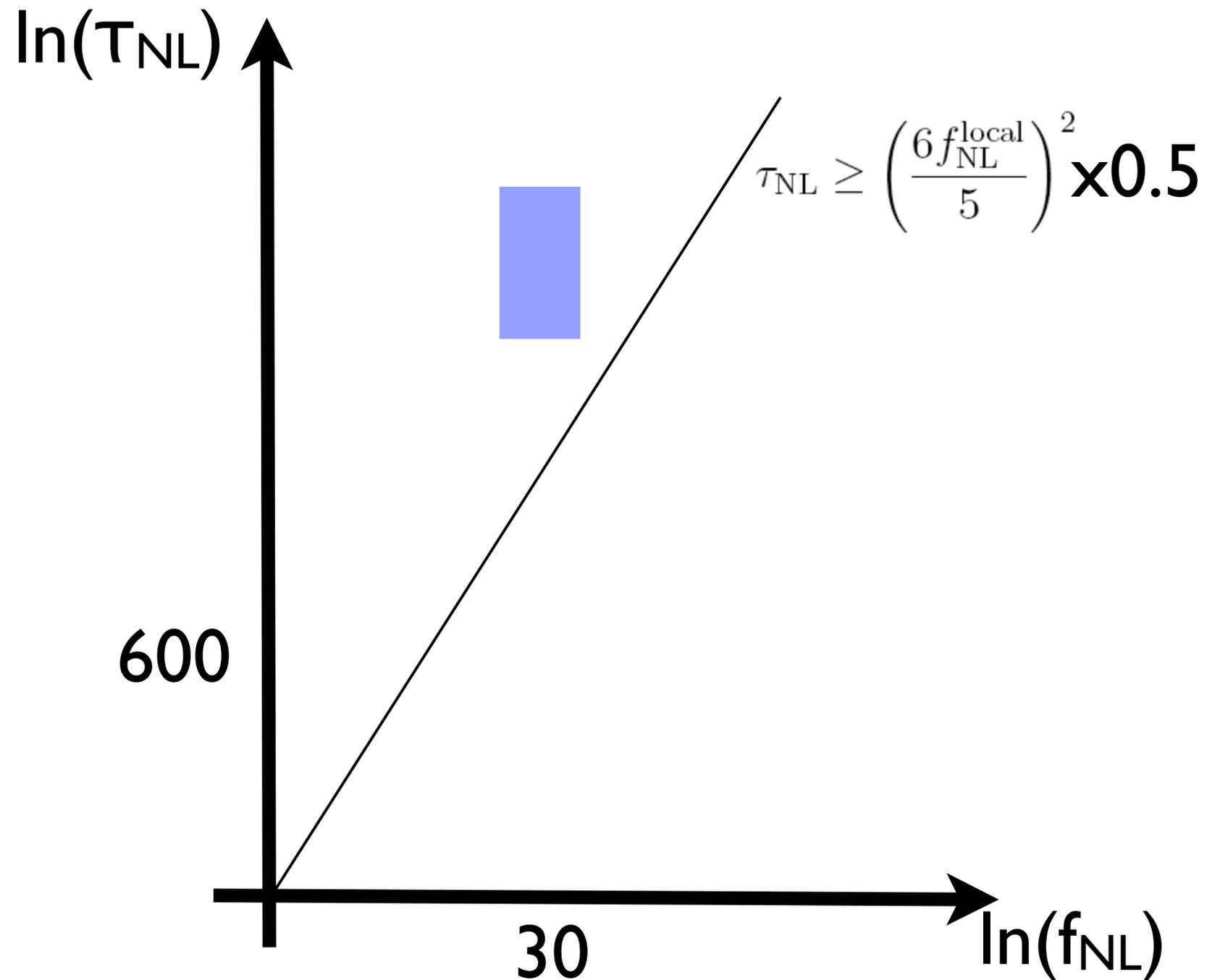
- The current limits from WMAP 7-year are consistent with single-field or multi-field models.
- So, let's play around with the future.

Case A: Single-field Happiness



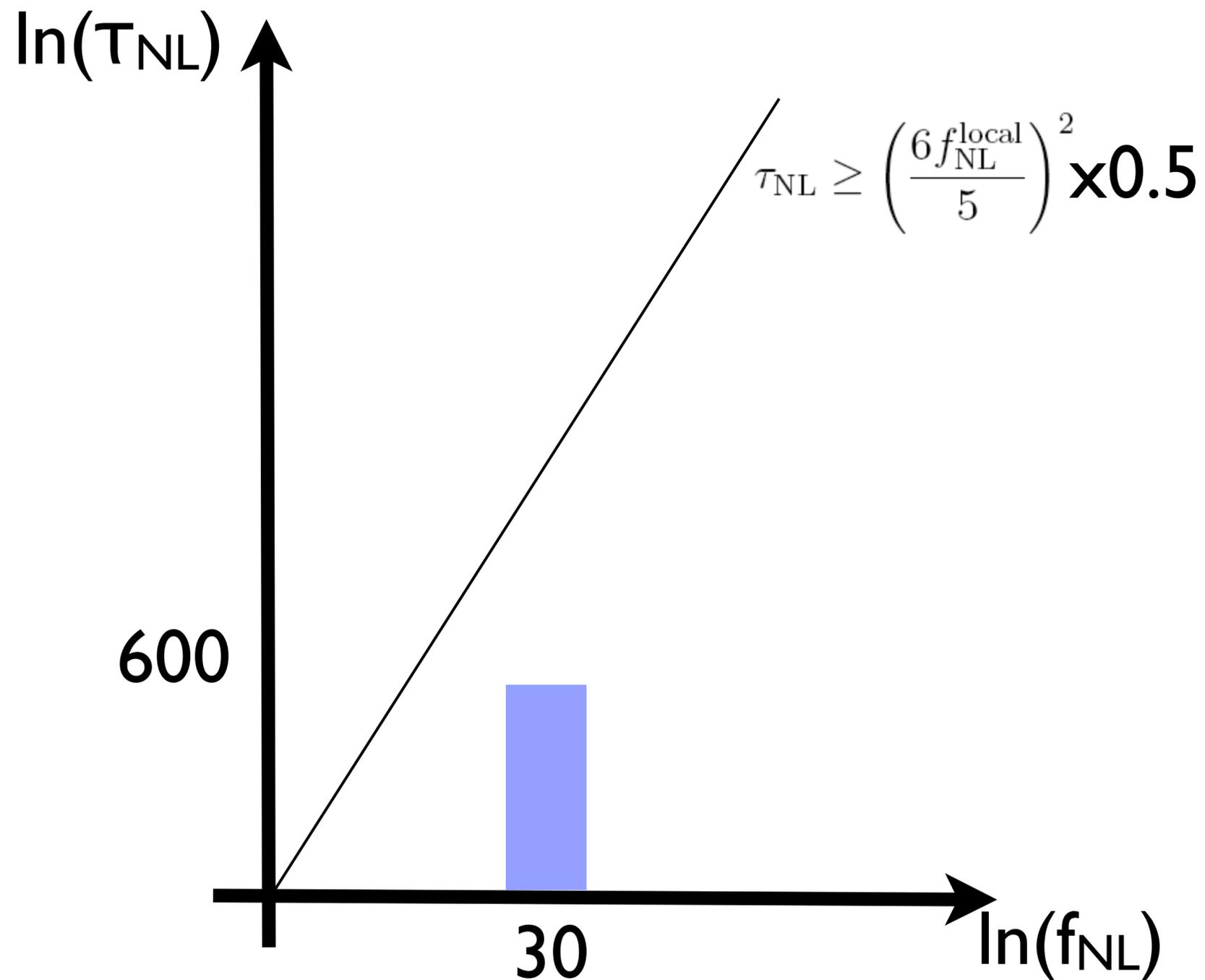
- No detection of anything after Planck. Single-field survived the test (for the moment: the future galaxy surveys can improve the limits by a factor of ten).

Case B: Multi-field Happiness



- f_{NL} is detected. Single-field is dead.
- But, τ_{NL} is also detected, in accordance with multi-field models: $\tau_{\text{NL}} > 0.5 \left(\frac{6f_{\text{NL}}}{5}\right)^2$ [Sugiyama, Komatsu & Futamase (2011)]

Case C: Madness



- f_{NL} is detected. Single-field is dead.
- But, τ_{NL} is **not** detected, inconsistent with the multi-field bound.
- (With the caveat that this bound may not be completely general) BOTH the single-field and multi-field are gone.

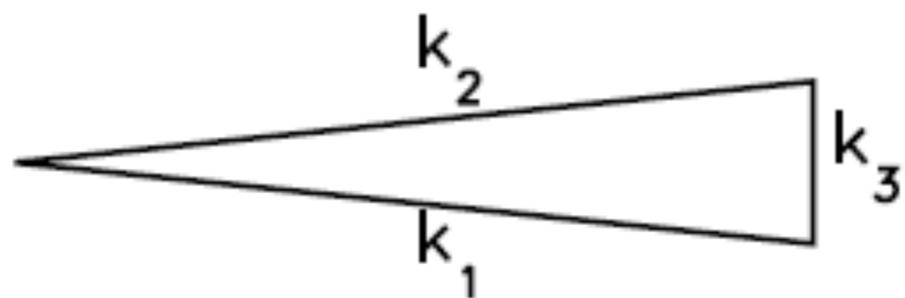
Beyond CMB: Large-scale Structure!

- In principle, the large-scale structure of the universe offers a lot more statistical power, because we can get 3D information. (CMB is 2D, so the number of Fourier modes is limited.)

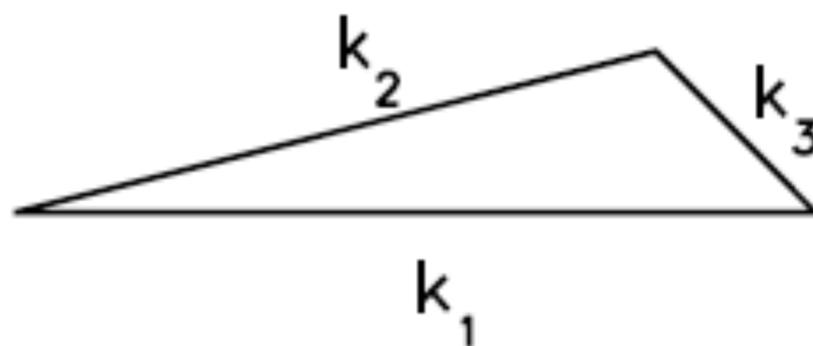
Beyond CMB: Large-scale Structure?

- Statistics is great, but the large-scale structure is non-linear, so perhaps it is less clean?
- Not necessarily.

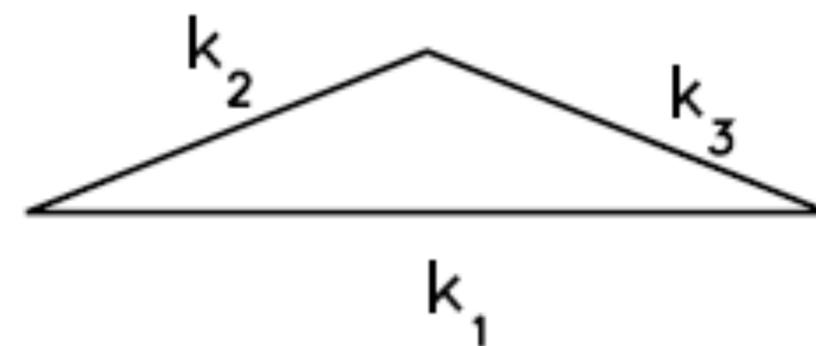
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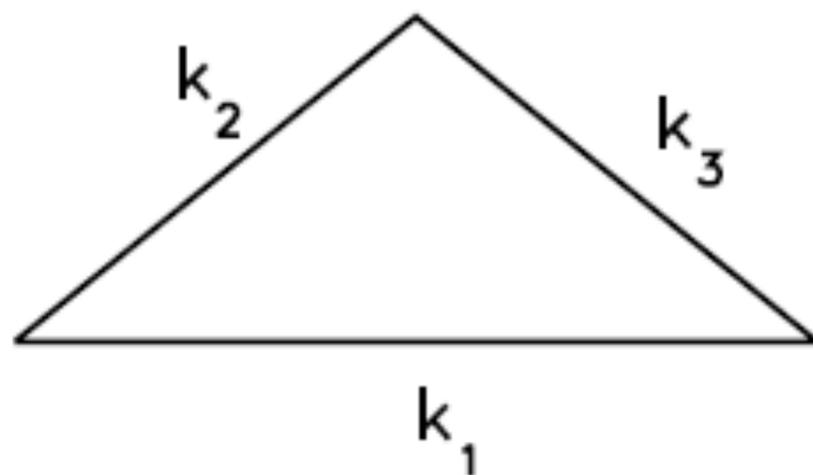


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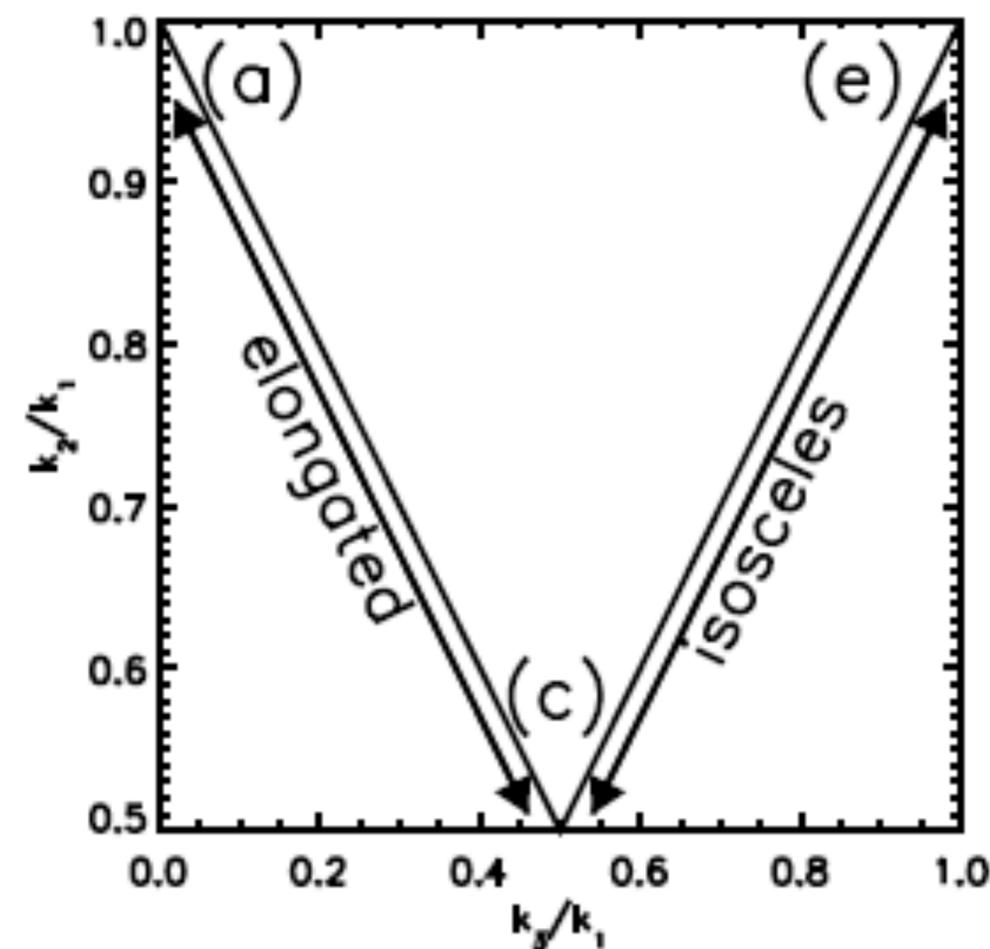
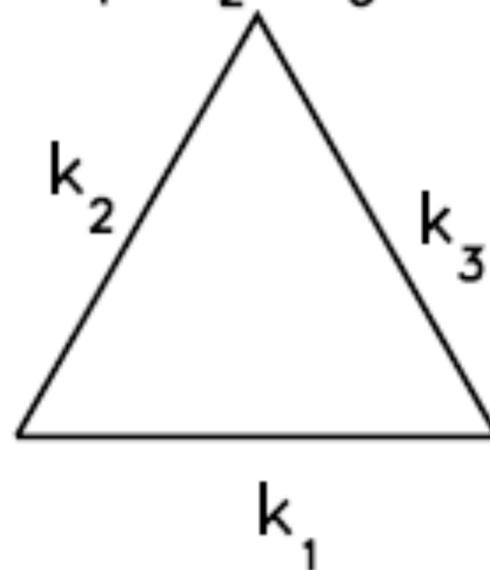


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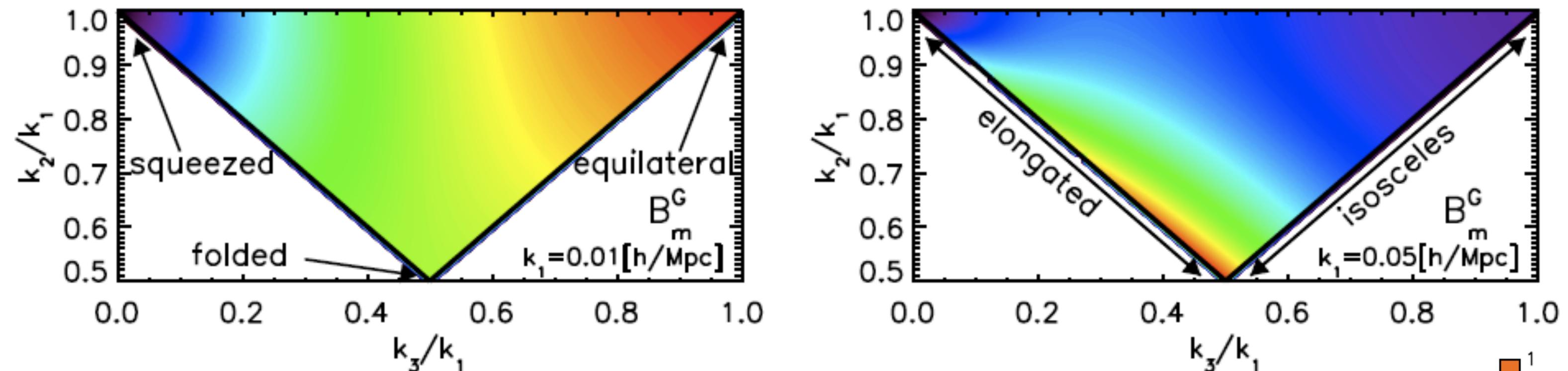
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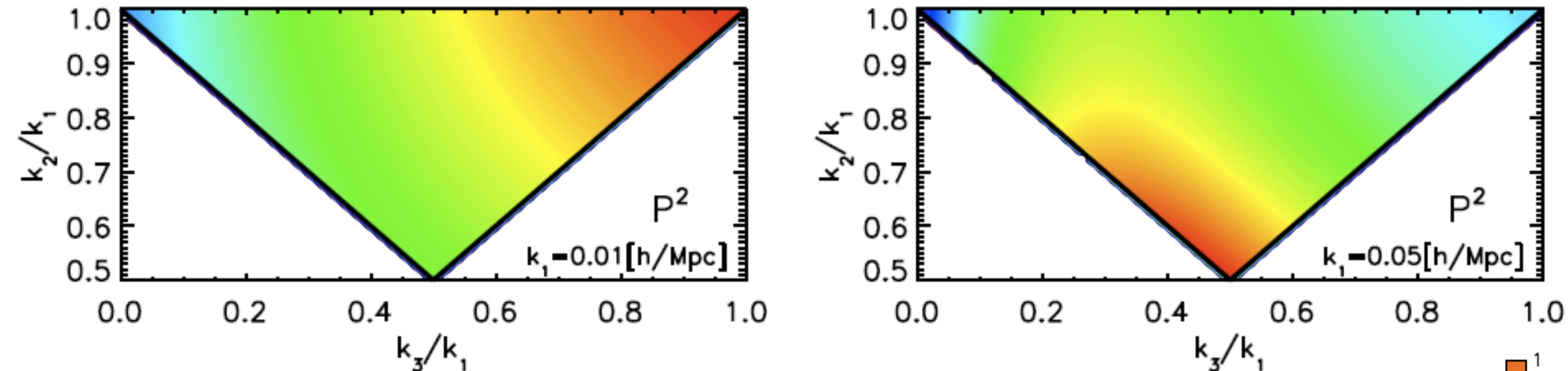
Non-linear Gravity



$$2b_1^3 \left[F_2^{(s)}(\mathbf{k}_1, \mathbf{k}_2) P_m(k_1, z) P_m(k_2, z) + (\text{cyclic}) \right]$$

- For a given k_1 , vary k_2 and k_3 , with $k_3 \leq k_2 \leq k_1$
- $F_2(k_2, k_3)$ vanishes in the squeezed limit, and peaks at the elongated triangles.

Non-linear Galaxy Bias

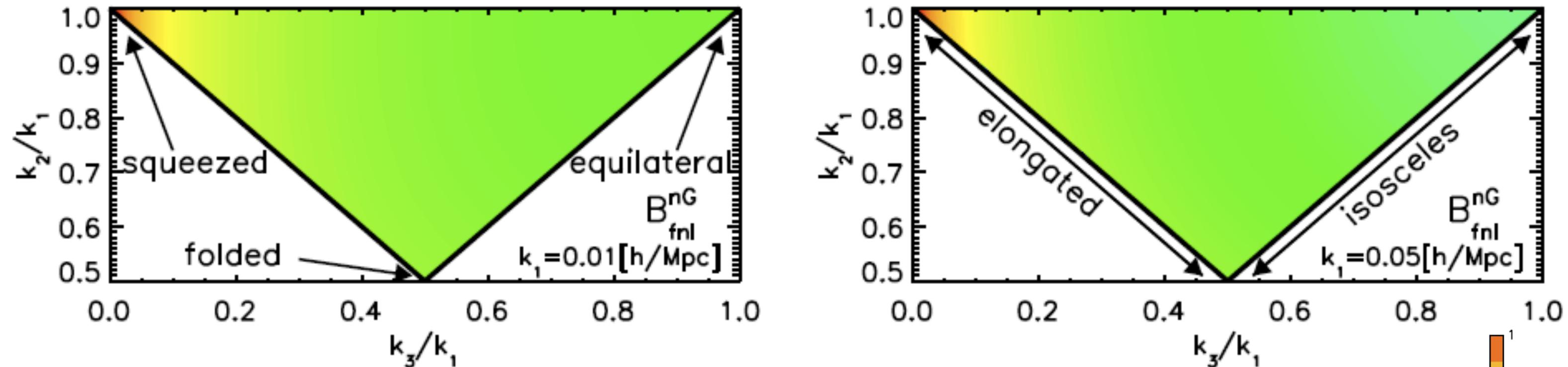


$$b_1^2 b_2 [P_m(k_1, z) P_m(k_2, z) + (\text{cyclic})]$$

- There is no F_2 : less suppression at the squeezed, and less enhancement along the elongated triangles.
- Still peaks at the equilateral or elongated forms.



Primordial Non-Gaussianity



$$3b_1^3 f_{\text{NL}} \Omega_m H_0^2 \left[\frac{P_m(k_1, z)}{k_1^2 T(k_1)} \frac{P_m(k_2, z)}{k_2^2 T(k_2)} \frac{k_3^2 T(k_3)}{D(z)} + (\text{cyclic}) \right]$$

- This gives the peaks at the squeezed configurations, clearly distinguishable from other non-linear/astrophysical effects.

Bispectrum is powerful

- $f_{\text{NL}}^{\text{local}} \sim O(1)$ is quite possible with the bispectrum method. (See Donghui Jeong's talk)
- This needs to be demonstrated by the real data! (e.g., SDSS-LRG)

Need For Dark “Energy”

- First of all, DE does not even need to be an energy.
- At present, *anything* that can explain the observed
 - (1) **Luminosity Distances** (Type Ia supernovae)
 - (2) **Angular Diameter Distances** (BAO, CMB)

simultaneously is qualified for being called “Dark Energy.”
- The candidates in the literature include: (a) energy, (b) modified gravity, and (c) extreme inhomogeneity.
- Measurements of the (3) **growth of structure** break degeneracy. (The best data right now is the X-ray clusters.)

$H(z)$: Current Knowledge

- $H^2(z) = H^2(0) [\Omega_r(1+z)^4 + \Omega_m(1+z)^3 + \Omega_k(1+z)^2 + \Omega_{de}(1+z)^{3(1+w)}]$
- (expansion rate) $H(0) = 70.2 \pm 1.4 \text{ km/s/Mpc}$
- (radiation) $\Omega_r = (8.4 \pm 0.3) \times 10^{-5}$
- (matter) $\Omega_m = 0.275 \pm 0.016$
- (curvature) $\Omega_k < 0.008$ (95%CL)
- (dark energy) $\Omega_{de} = 0.725 \pm 0.015$
- (DE equation of state) $w = -1.00 \pm 0.06$

H(z) to Distances

- Comoving Distance
 - $\chi(z) = c \int^z [dz'/H(z')]$
- Luminosity Distance
 - $D_L(z) = (1+z)\chi(z) [1 - (k/6)\chi^2(z)/R^2 + \dots]$
 - $R = (\text{curvature radius of the universe}); k = (\text{sign of curvature})$
 - WMAP 7-year limit: $R > 2\chi(\infty)$; justify the Taylor expansion
- Angular Diameter Distance
 - $D_A(z) = [\chi(z)/(1+z)] [1 - (k/6)\chi^2(z)/R^2 + \dots]$

$$D_A(z) = (1+z)^{-2} D_L(z)$$

$D_L(z)$

Type Ia Supernovae

$D_A(z)$

Galaxies (BAO)

CMB

0.02

0.2

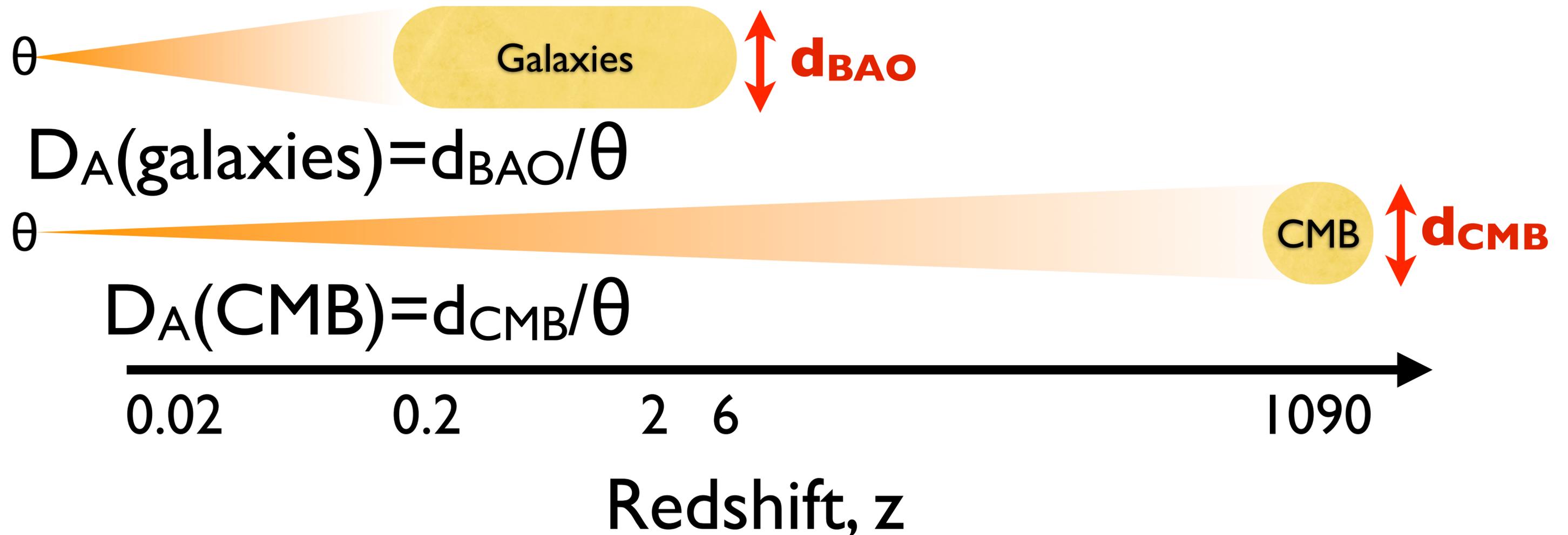
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Redshift, z

- To measure $D_A(z)$, we need to know the intrinsic size.
- What can we use as the *standard ruler*?

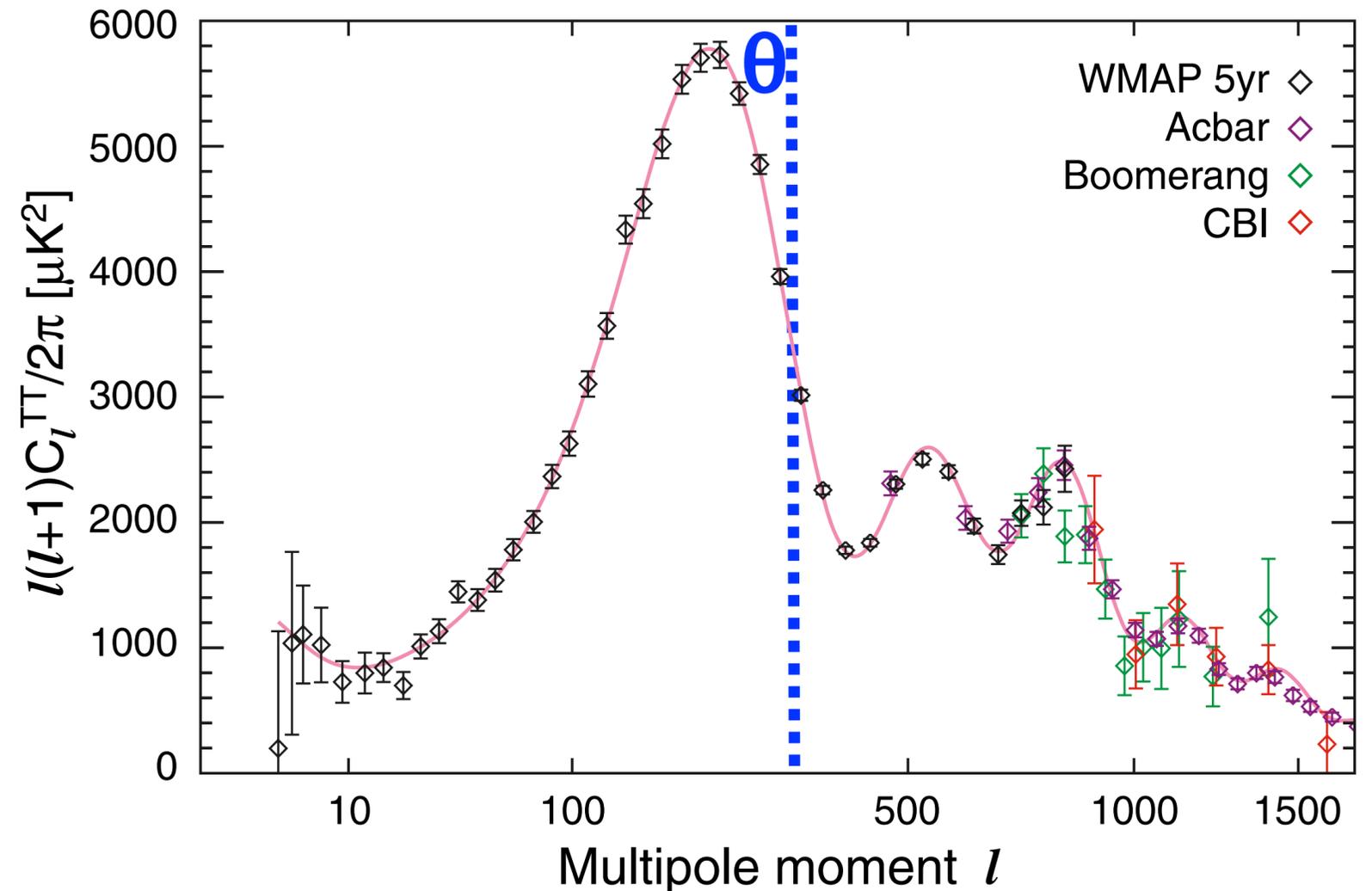
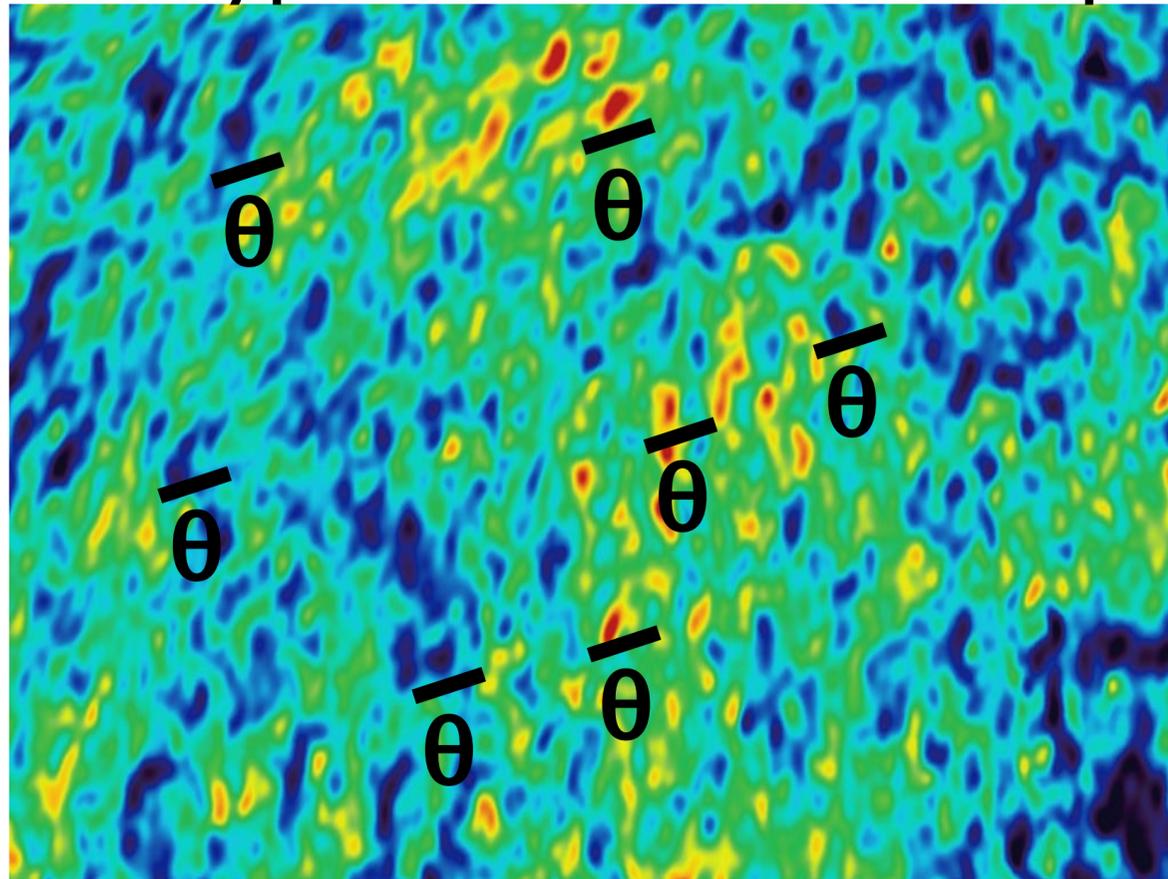
How Do We Measure $D_A(z)$?



- If we know the intrinsic physical sizes, d , we can measure D_A . What determines d ?

CMB as a Standard Ruler

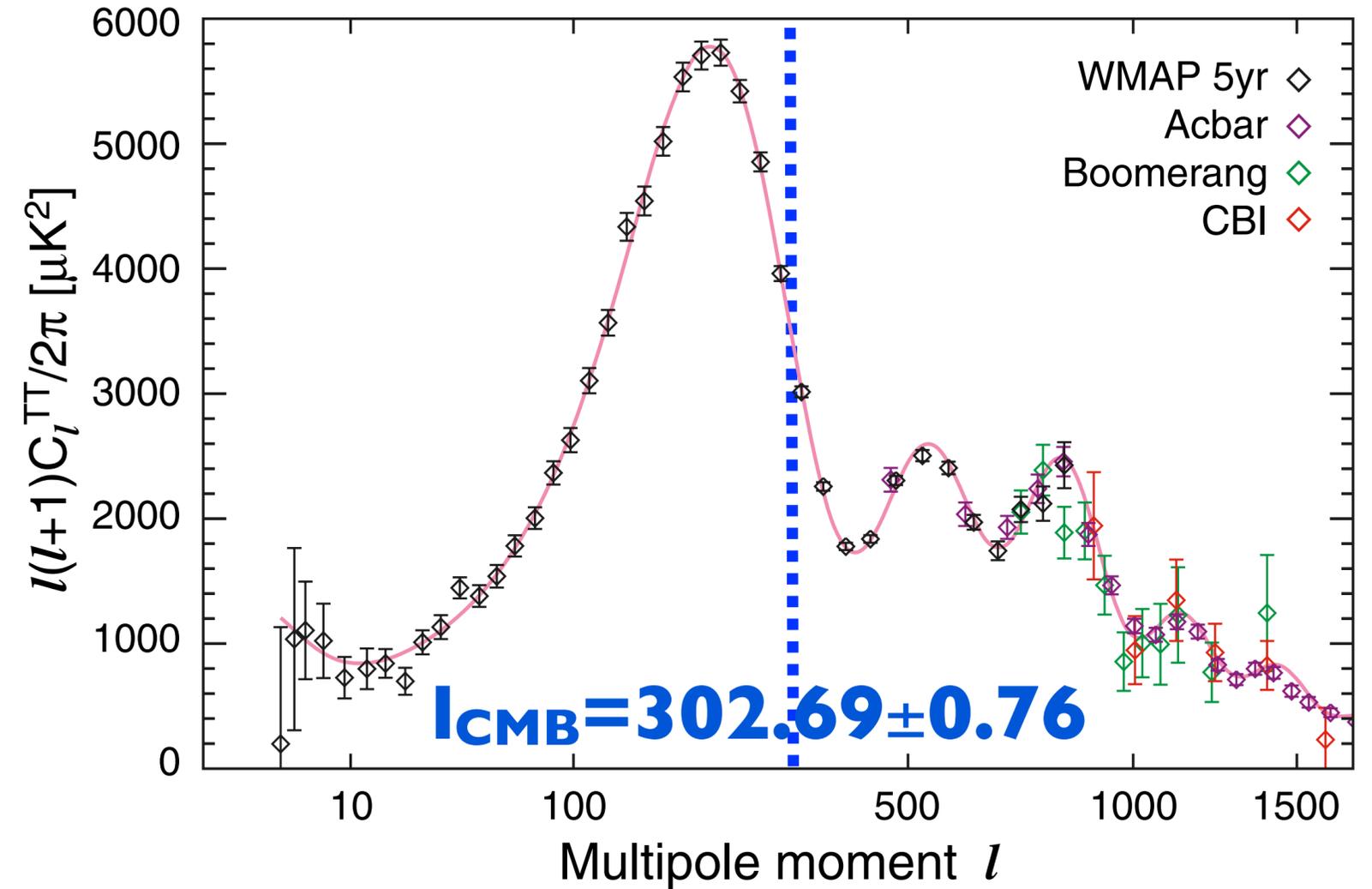
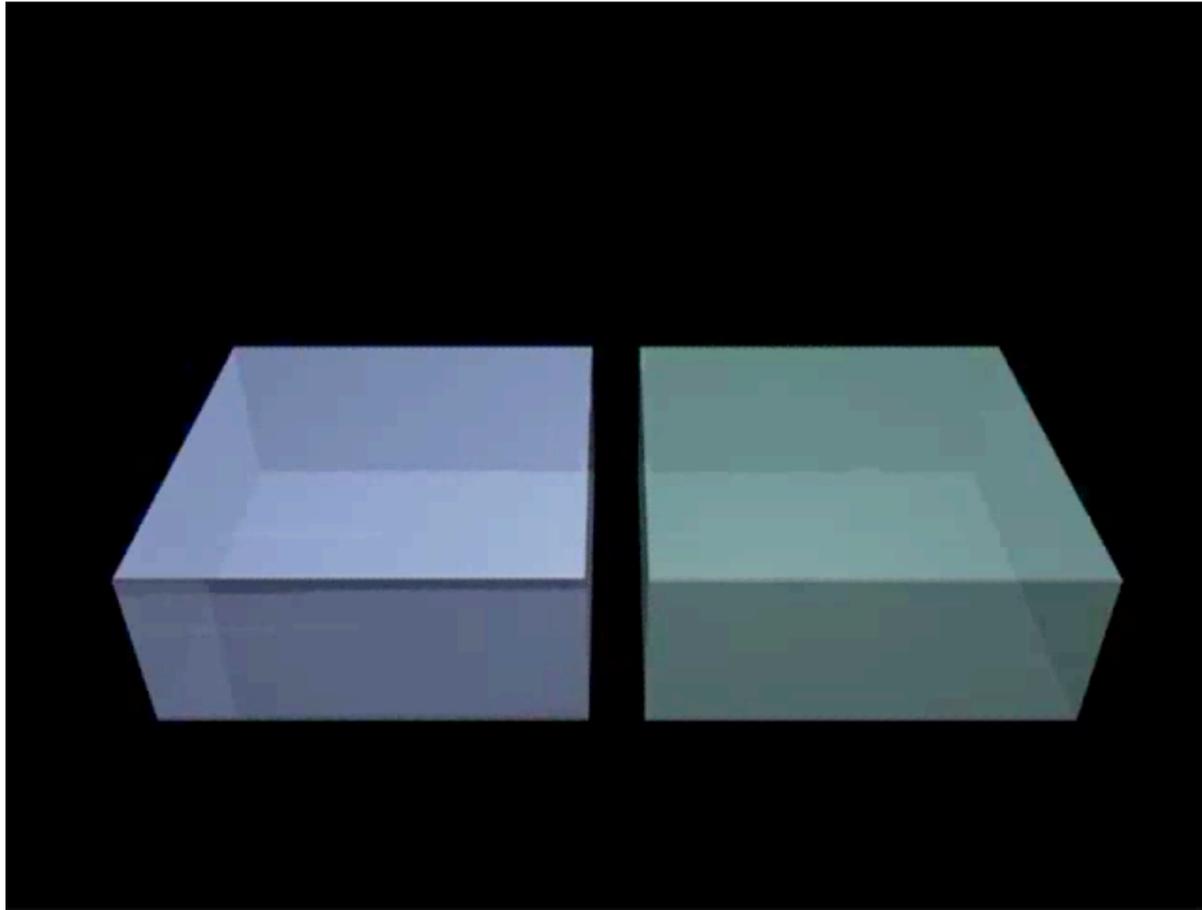
θ ~ the typical size of hot/cold spots



- The existence of typical spot size in image space yields oscillations in harmonic (Fourier) space. What determines the physical size of typical spots, d_{CMB} ? ³⁸

Sound Horizon

- The typical spot size, d_{CMB} , is determined by the **physical distance traveled by the sound wave** from the Big Bang to the decoupling of photons at $z_{\text{CMB}} \sim 1090$ ($t_{\text{CMB}} \sim 380,000$ years).
- The causal horizon (photon horizon) at t_{CMB} is given by
 - $d_{\text{H}}(t_{\text{CMB}}) = a(t_{\text{CMB}}) * \text{Integrate} [\mathbf{c} \, dt/a(t), \{t, 0, t_{\text{CMB}}\}]$.
- The sound horizon at t_{CMB} is given by
 - $d_{\text{s}}(t_{\text{CMB}}) = a(t_{\text{CMB}}) * \text{Integrate} [\mathbf{c}_{\text{s}}(\mathbf{t}) \, dt/a(t), \{t, 0, t_{\text{CMB}}\}]$, where $c_{\text{s}}(t)$ is the time-dependent **speed of sound of photon-baryon fluid**.



- The WMAP 7-year values:

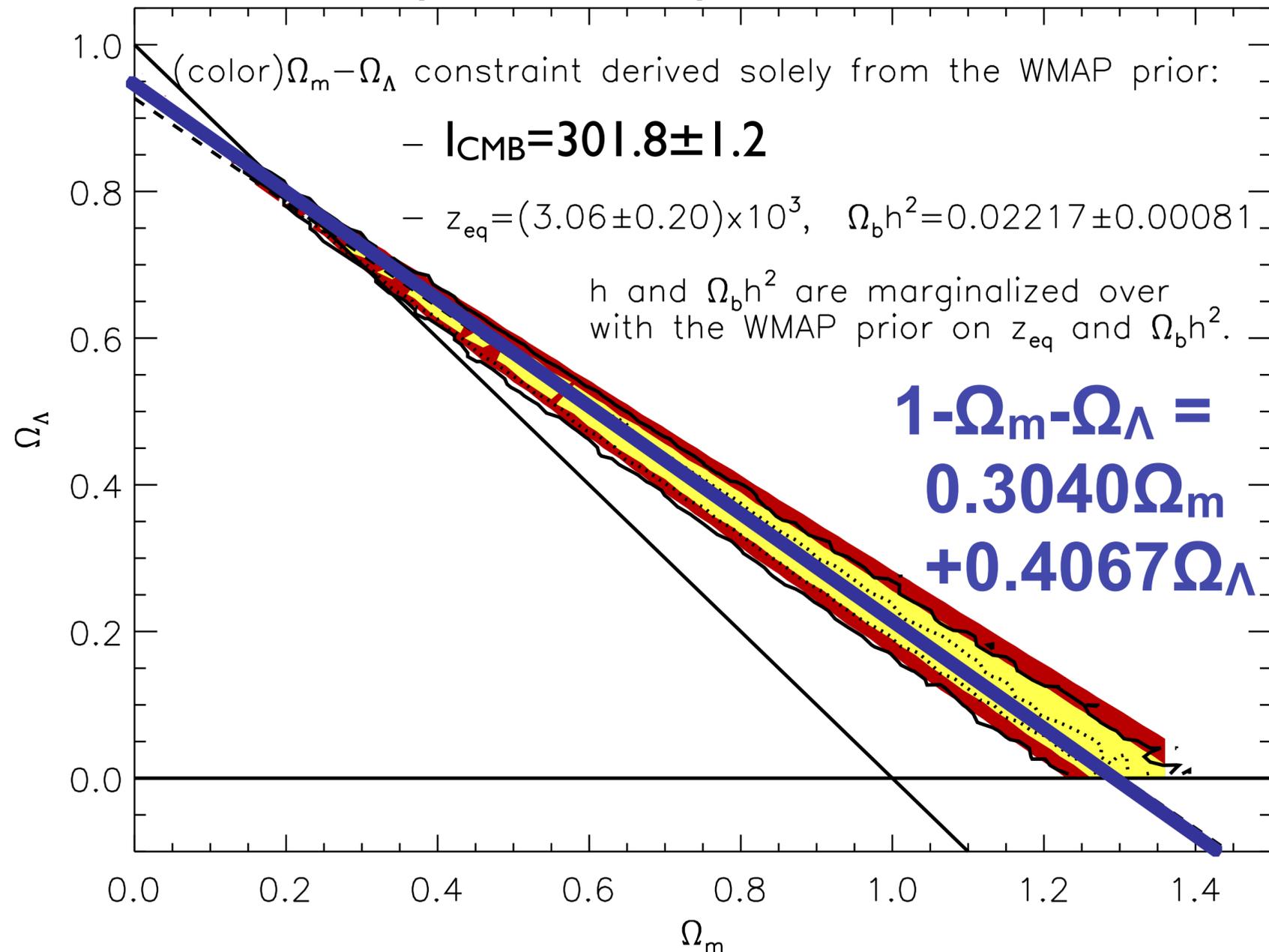
- $l_{\text{CMB}} = \pi/\theta = \pi D_A(z_{\text{CMB}})/d_s(z_{\text{CMB}}) = 302.69 \pm 0.76$

- CMB data constrain the ratio, $D_A(z_{\text{CMB}})/d_s(z_{\text{CMB}})$.

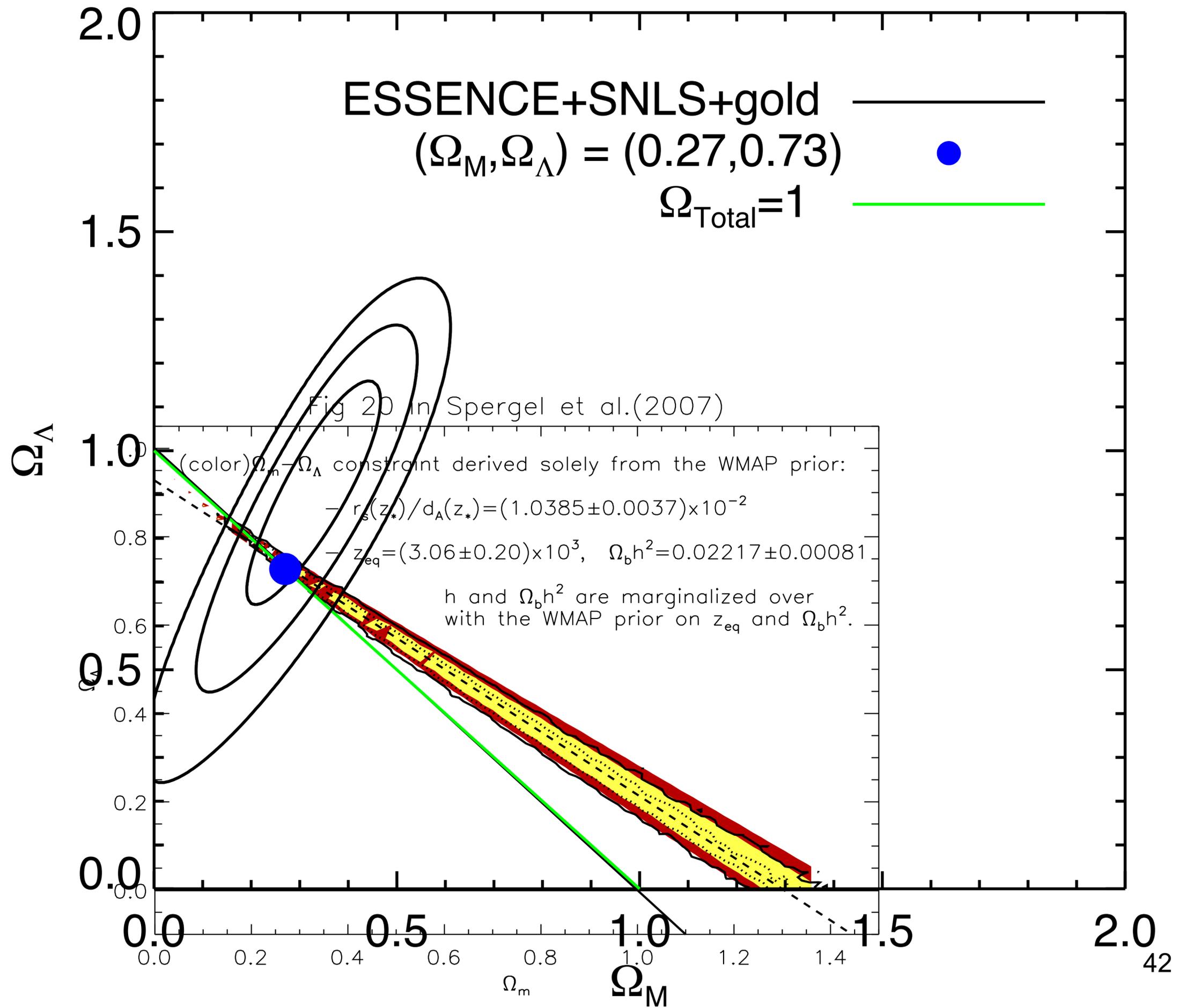
- $r_s(z_{\text{CMB}}) = (1+z_{\text{CMB}})d_s(z_{\text{CMB}}) = 146.6 \pm 1.6$ Mpc (comoving)

What $D_A(z_{\text{CMB}})/d_s(z_{\text{CMB}})$ Gives You (3-year example)

Fig 20 in Spergel et al.(2007)

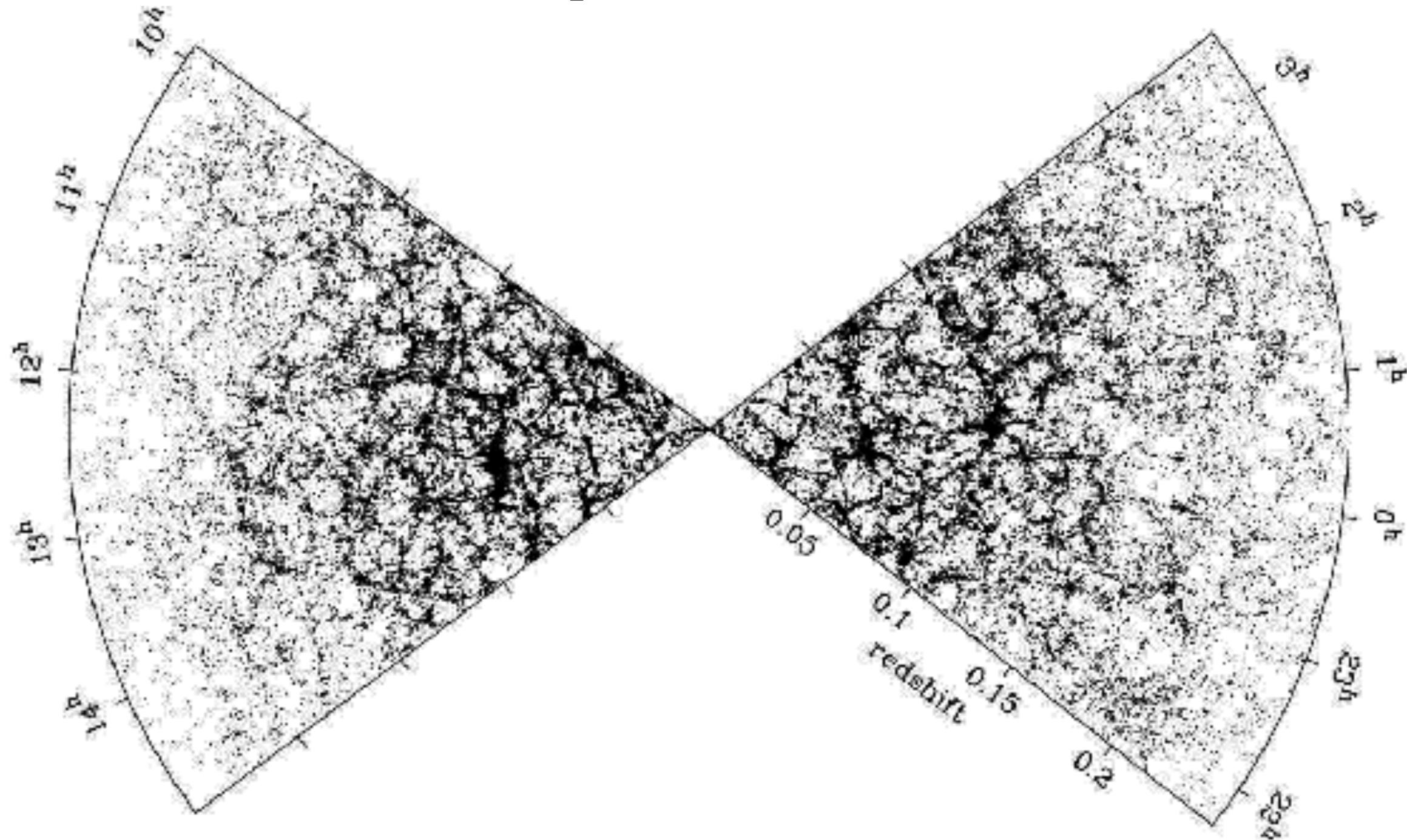


- **Color**: constraint from $l_{\text{CMB}} = \pi D_A(z_{\text{CMB}})/d_s(z_{\text{CMB}})$ with z_{EQ} & $\Omega_b h^2$.
- Black contours: Markov Chain from WMAP 3yr (Spergel et al. 2007)



BAO in Galaxy Distribution

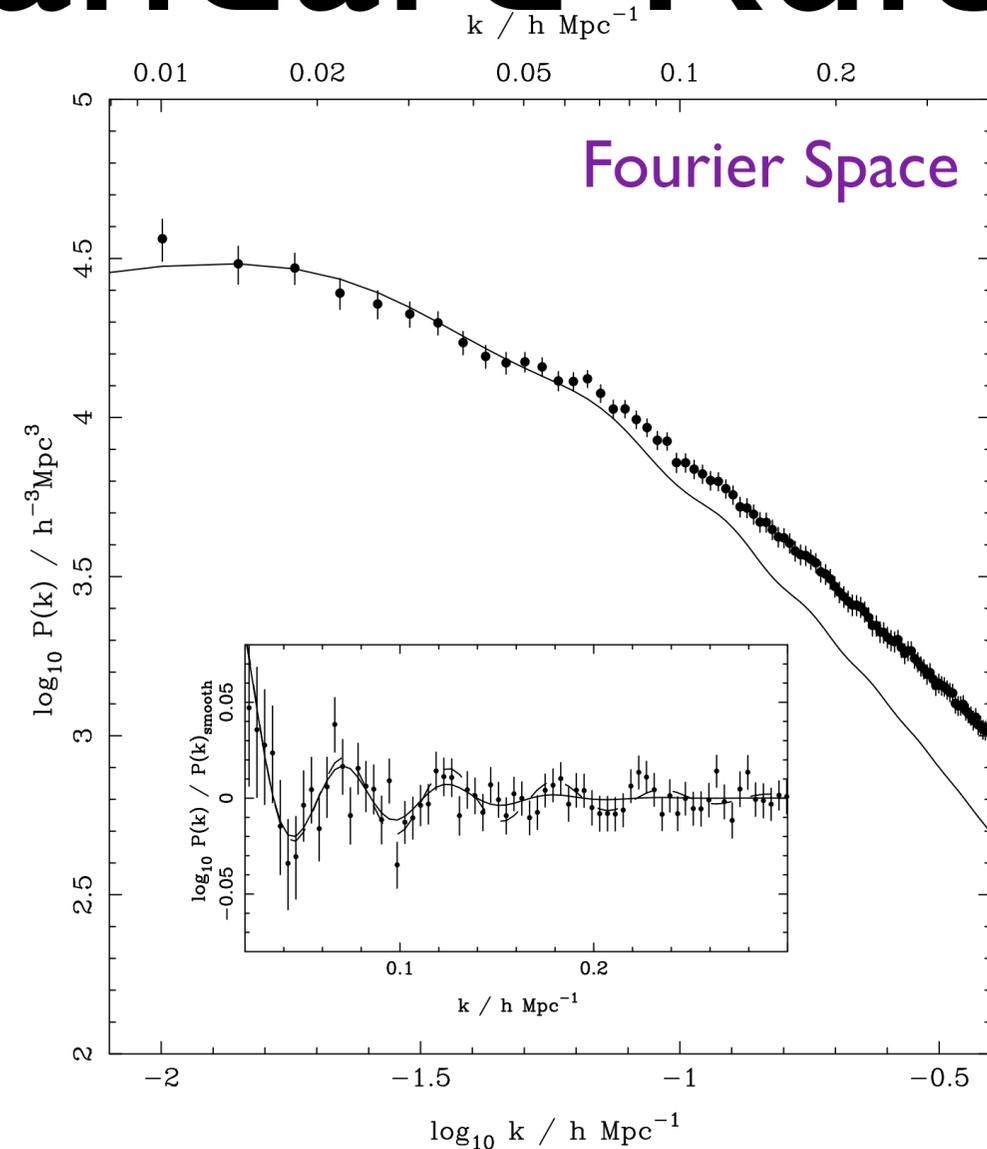
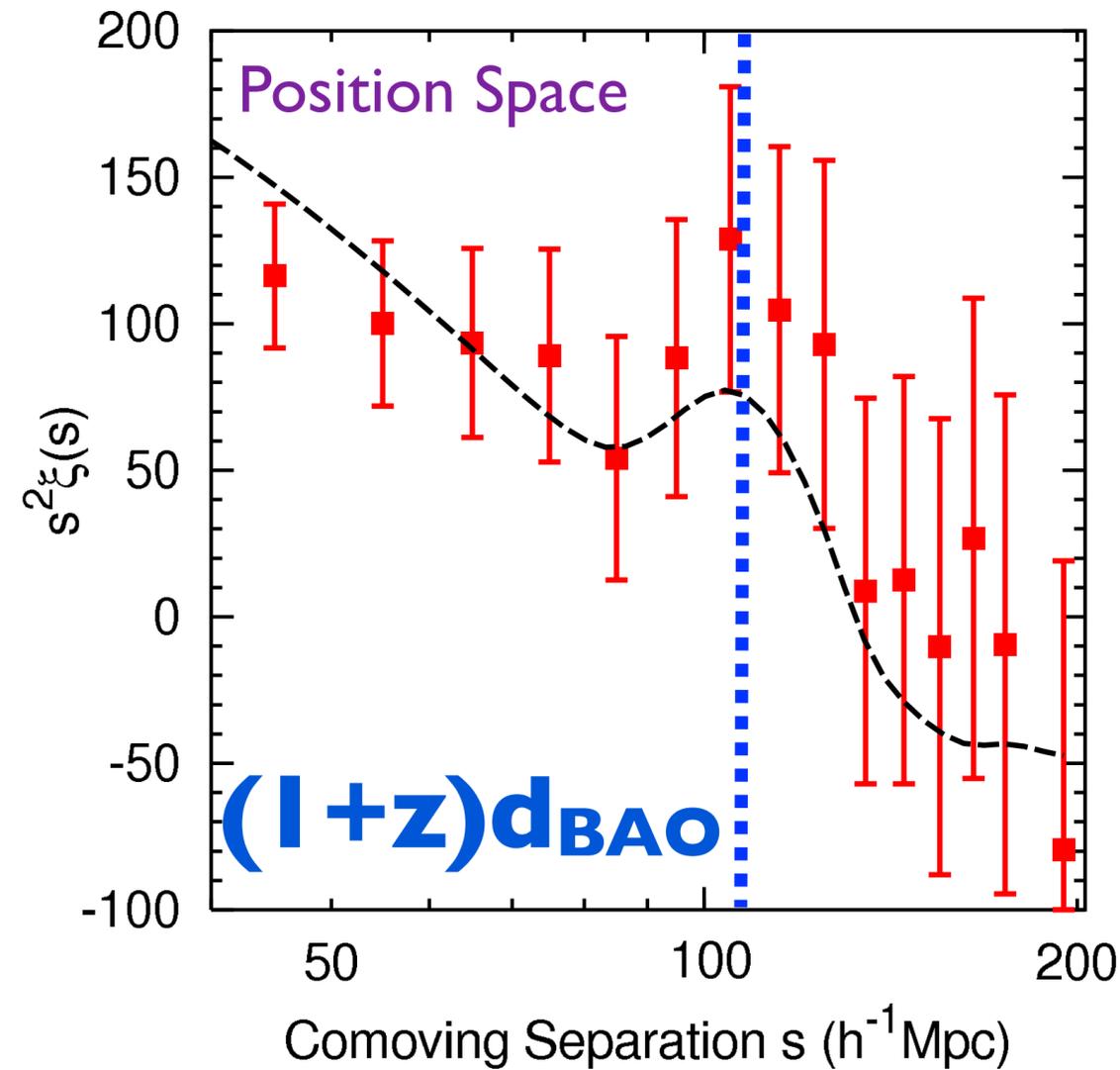
2dFGRS



- The same acoustic oscillations should be hidden in this galaxy distribution...

BAO as a Standard Ruler

Okumura et al. (2007)



Percival et al. (2006)

- The existence of a localized clustering scale in the 2-point function yields oscillations in Fourier space.

Sound Horizon Again

- The clustering scale, d_{BAO} , is given by the physical distance traveled by the sound wave from the Big Bang to the **decoupling of baryons** at $z_{\text{BAO}} = 1020.5 \pm 1.6$ (c.f., $z_{\text{CMB}} = 1091 \pm 1$).
- The baryons decoupled slightly later than CMB.
 - By the way, this is not universal in cosmology, but *accidentally* happens to be the case for our Universe.
 - If $3\rho_{\text{baryon}}/(4\rho_{\text{photon}}) = 0.64(\Omega_b h^2/0.022)(1090/(1+z_{\text{CMB}}))$ is greater than unity, $z_{\text{BAO}} > z_{\text{CMB}}$. Since our Universe happens to have $\Omega_b h^2 = 0.022$, $z_{\text{BAO}} < z_{\text{CMB}}$. (ie, $d_{\text{BAO}} > d_{\text{CMB}}$)

Standard Rulers in CMB & Matter

	Quantity	Eq.	5-year WMAP
CMB	z_*	(66)	1090.51 ± 0.95
CMB	$r_s(z_*)$	(6)	146.8 ± 1.8 Mpc
Matter	z_d	(3)	1020.5 ± 1.6
Matter	$r_s(z_d)$	(6)	153.3 ± 2.0 Mpc

- For flat LCDM, but very similar results for $w \neq -1$ and curvature $\neq 0$!

Not Just $D_A(z)$...

- A really nice thing about BAO at a given redshift is that it can be used to measure not only $D_A(z)$, but also the expansion rate, $H(z)$, directly, at **that** redshift.

- BAO perpendicular to l.o.s

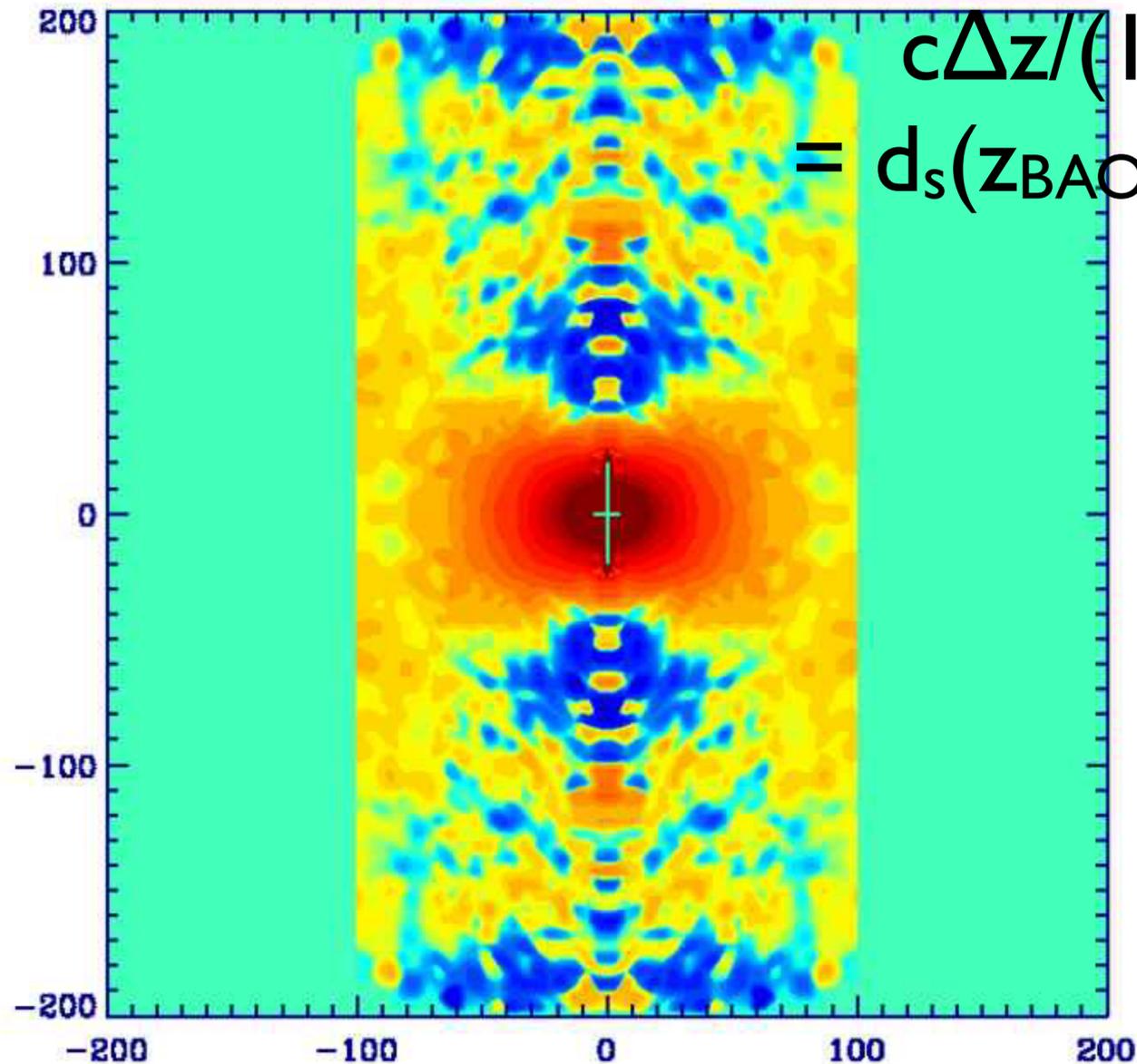
$$\Rightarrow D_A(z) = d_s(z_{\text{BAO}})/\theta$$

- BAO parallel to l.o.s

$$\Rightarrow \mathbf{H(z) = c\Delta z / [(1+z)d_s(z_{\text{BAO}})]}$$

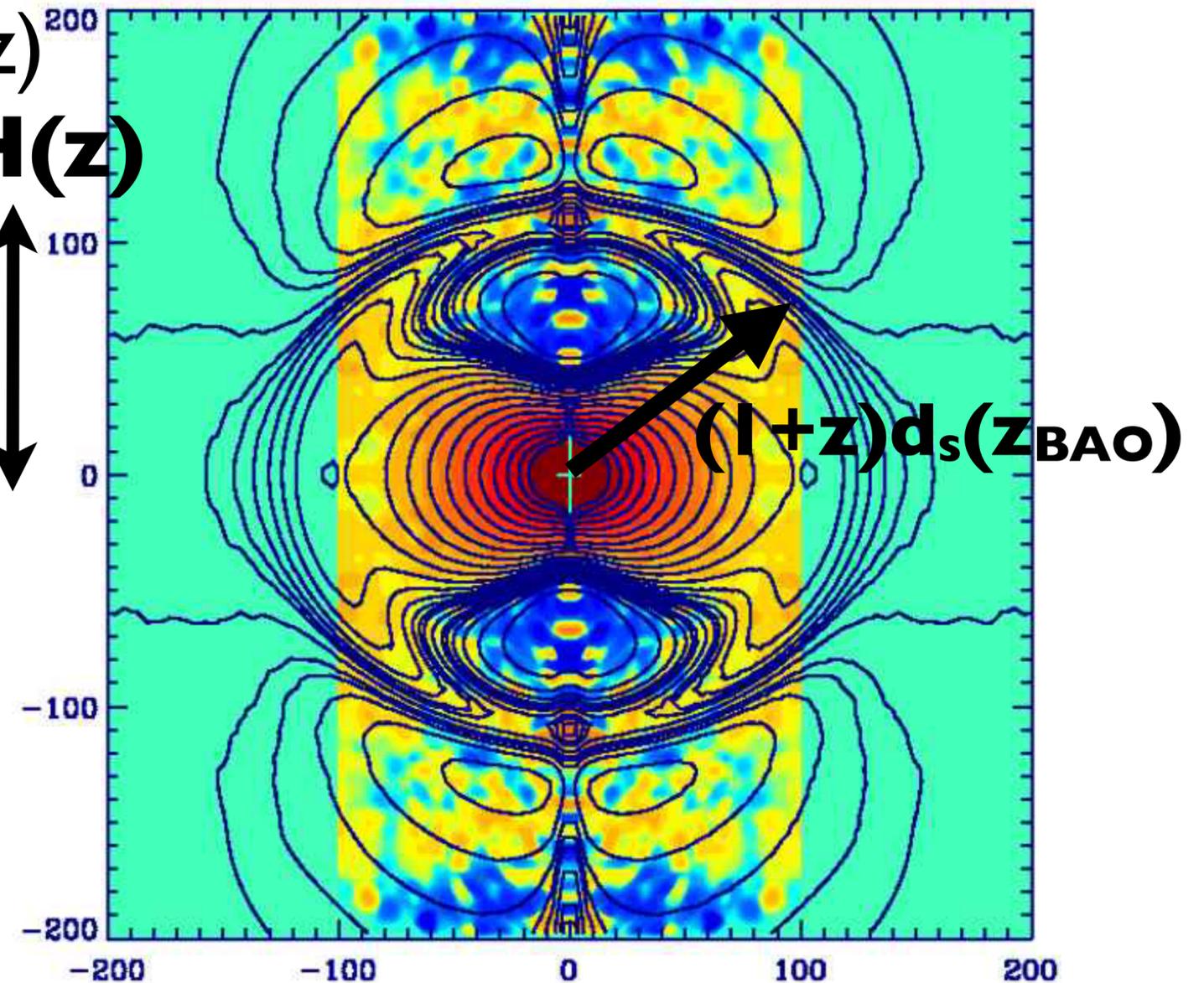
Transverse= $D_A(z)$; Radial= $H(z)$

SDSS Data
DR6



$$\frac{c\Delta z}{(1+z)} = d_s(z_{\text{BAO}}) \mathbf{H}(\mathbf{z})$$

Linear Theory
DR6 + best model

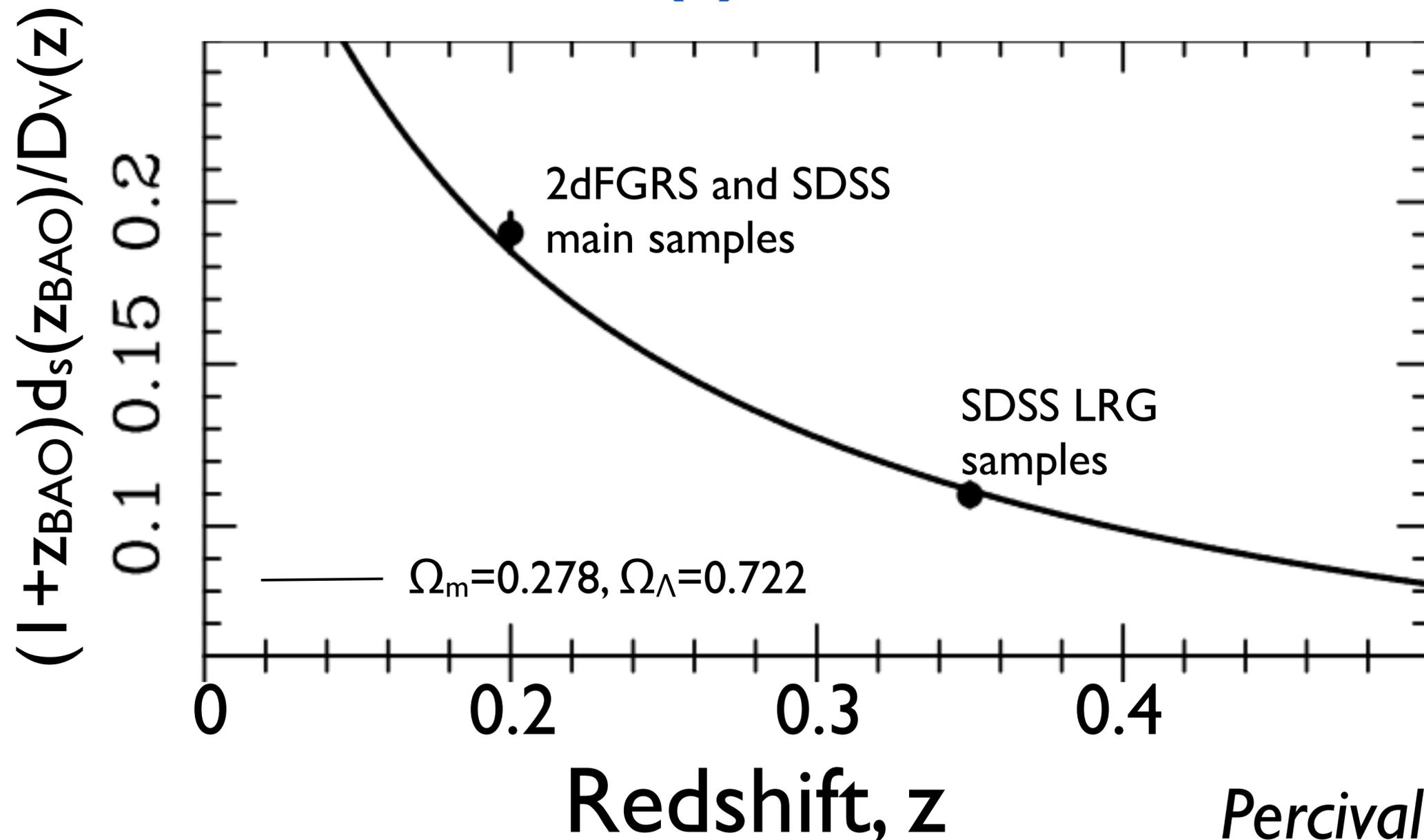


$$\theta = d_s(z_{\text{BAO}}) / \mathbf{D}_A(\mathbf{z})$$

Two-point correlation function measured from the SDSS Luminous Red Galaxies (Gaztanaga, Cabre & Hui 2008)

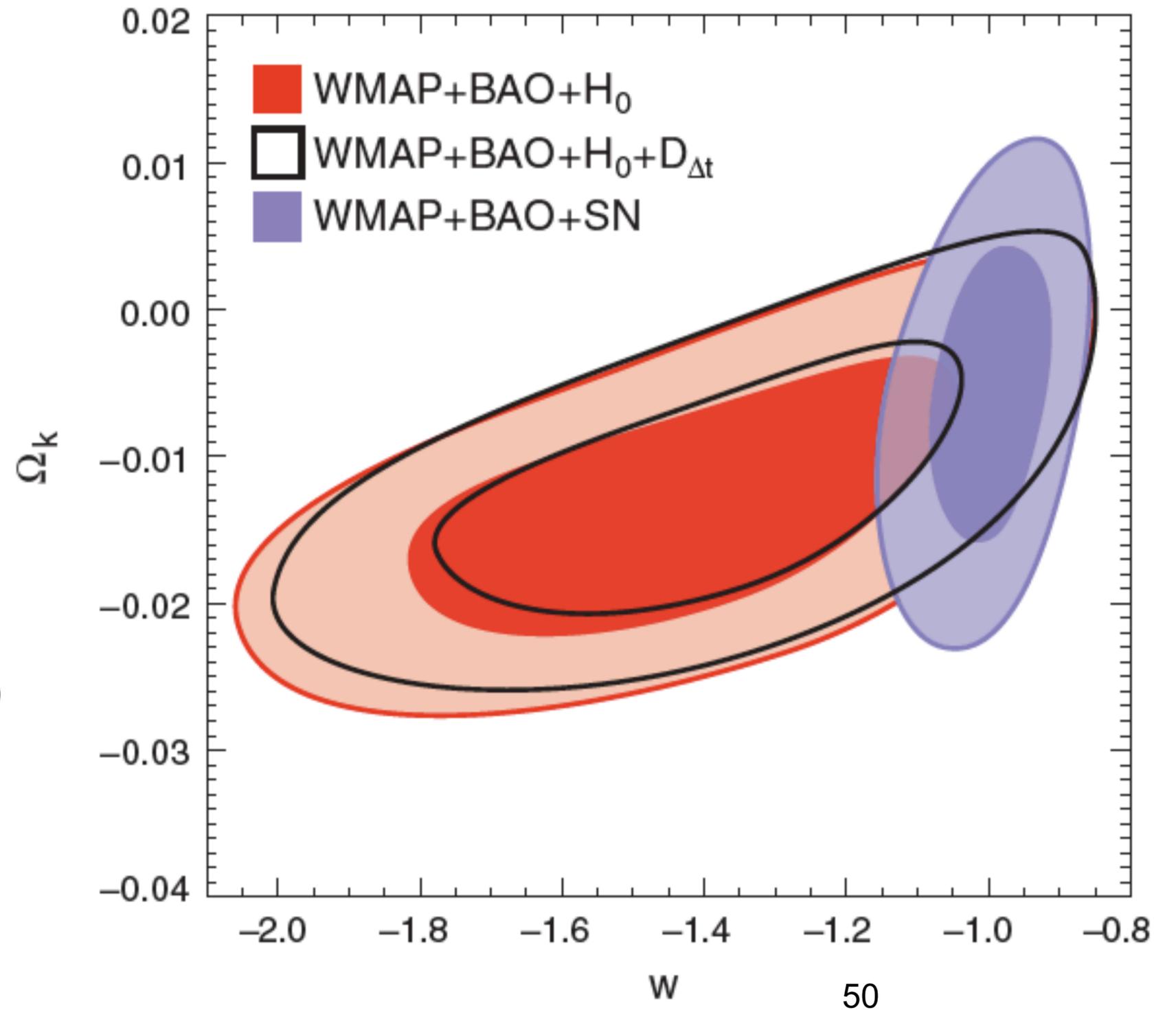
$$D_V(z) = \left\{ (1+z)^2 D_A^2(z) [cz/H(z)] \right\}^{1/3}$$

Since the current data are not good enough to constrain $D_A(z)$ and $H(z)$ separately, a combination distance, $D_V(z)$, has been constrained.



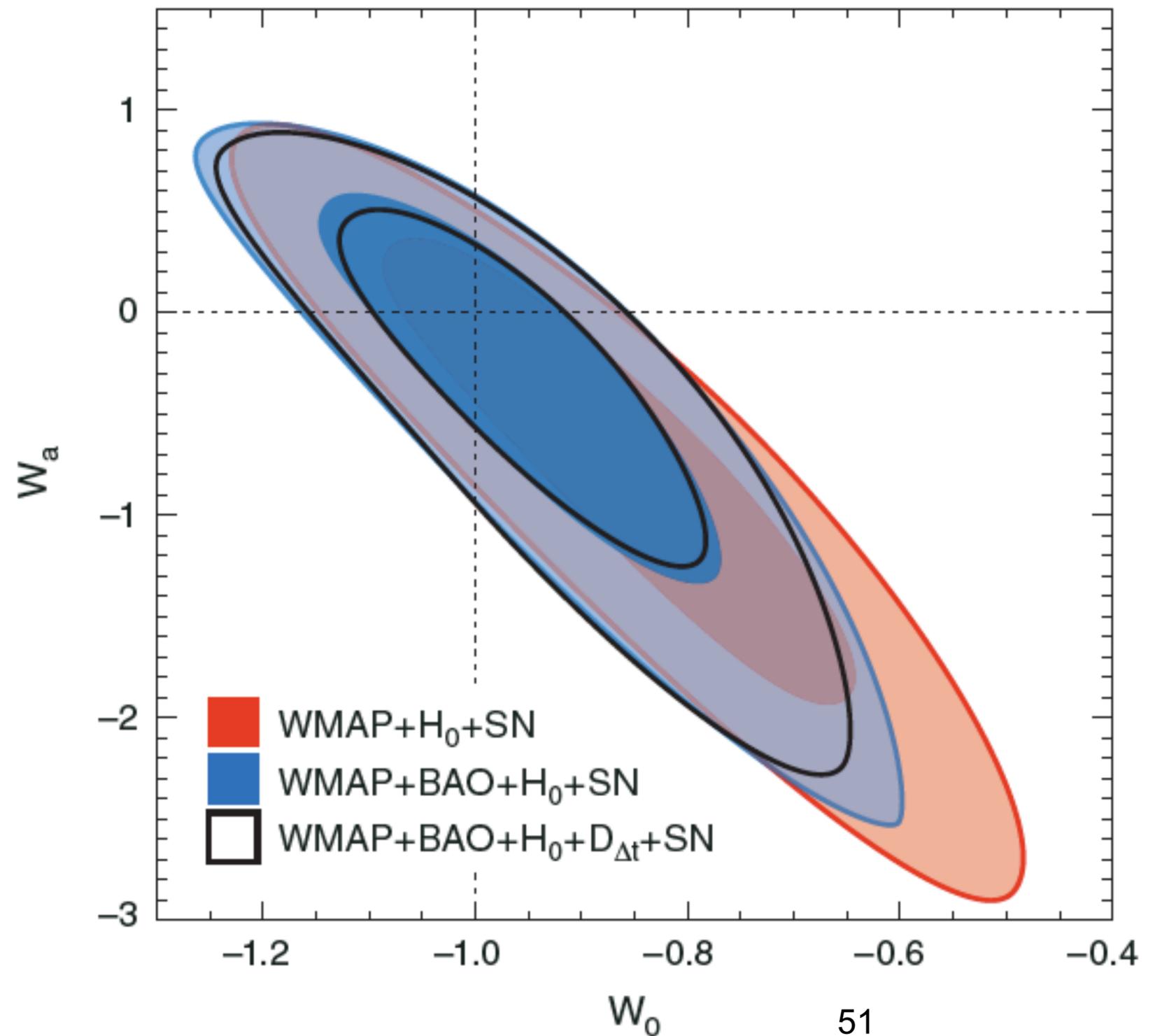
WMAP7+BAO+...

- At the moment, BAO is great for fixing curvature, but not good for fixing w
- We still need supernovae for fixing w , but this would change as more BAO data (especially at higher redshifts) become available.

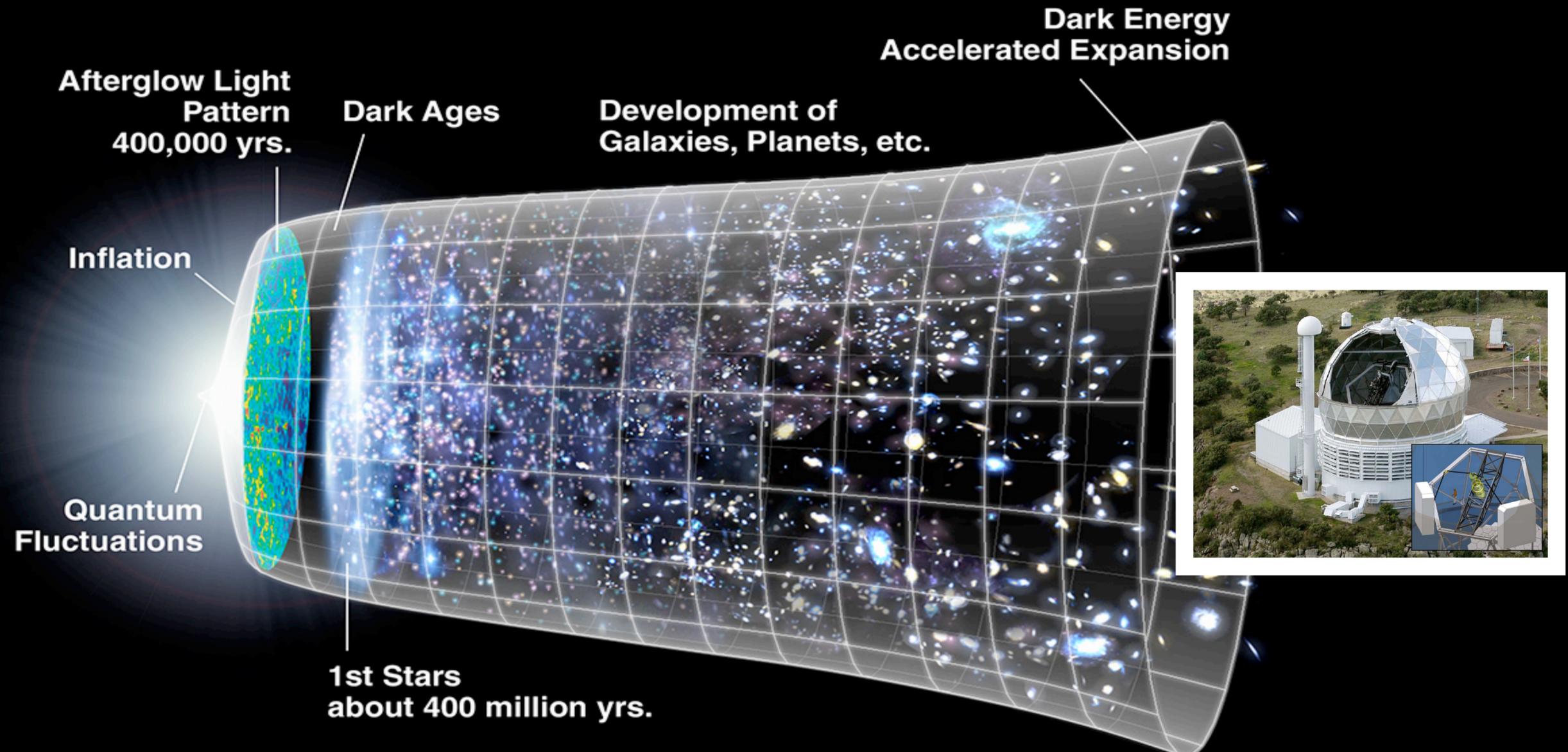


$$w(z) = w_0 + w_a * z / (1 + z)$$

- Cosmological constant, $w_0 = -1$ and $w_a = 0$, are perfectly consistent with data.
- Of course we all want this to change at some point...

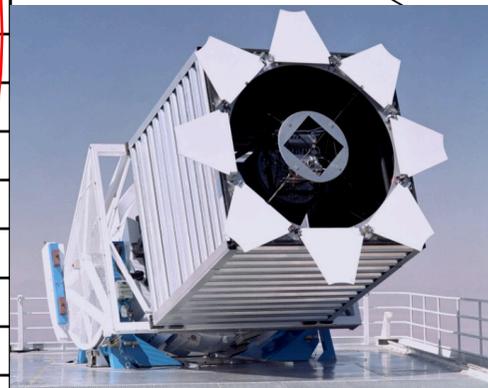
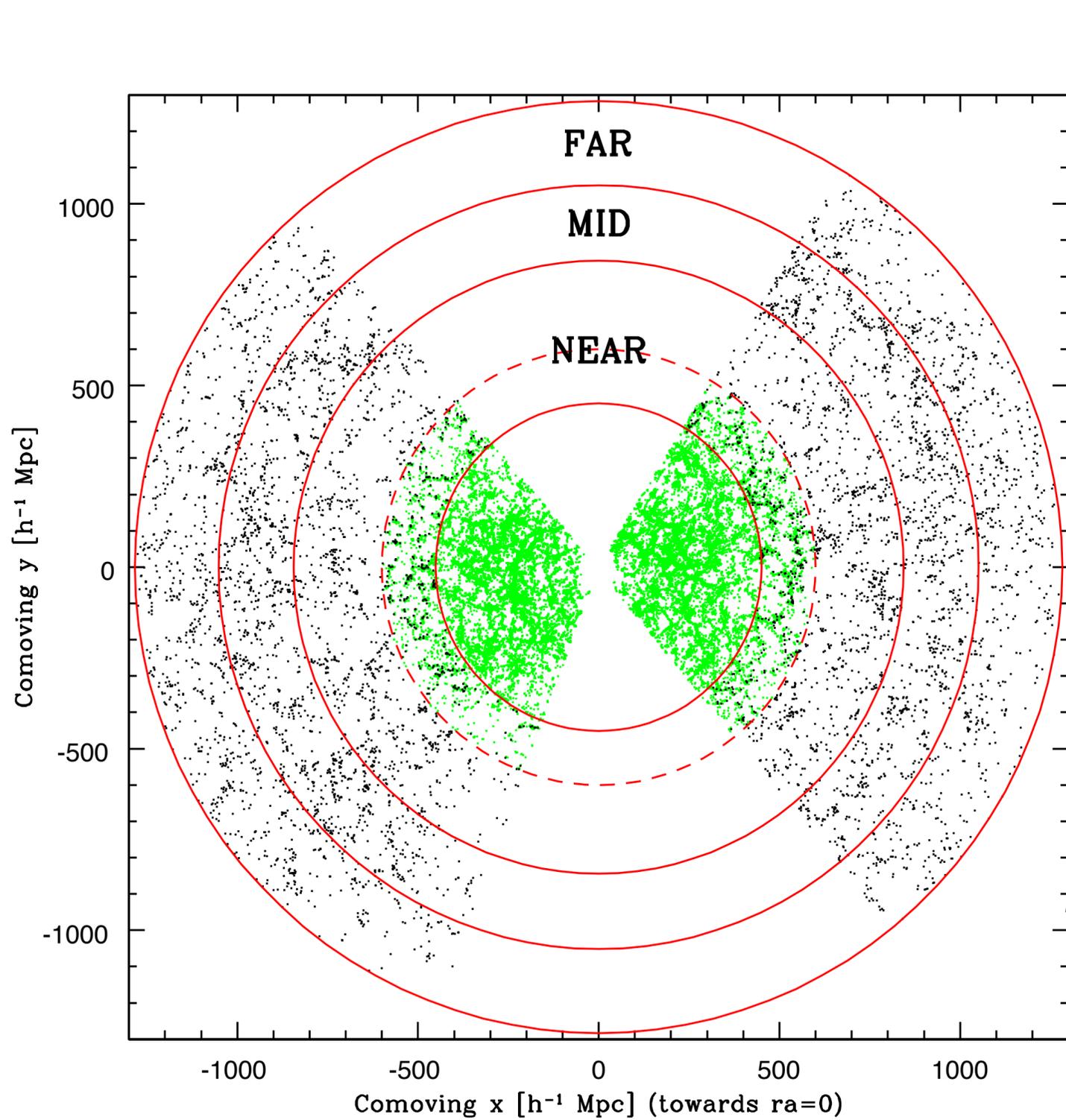


Hobby-Eberly Telescope Dark Energy Experiment (HETDEX)

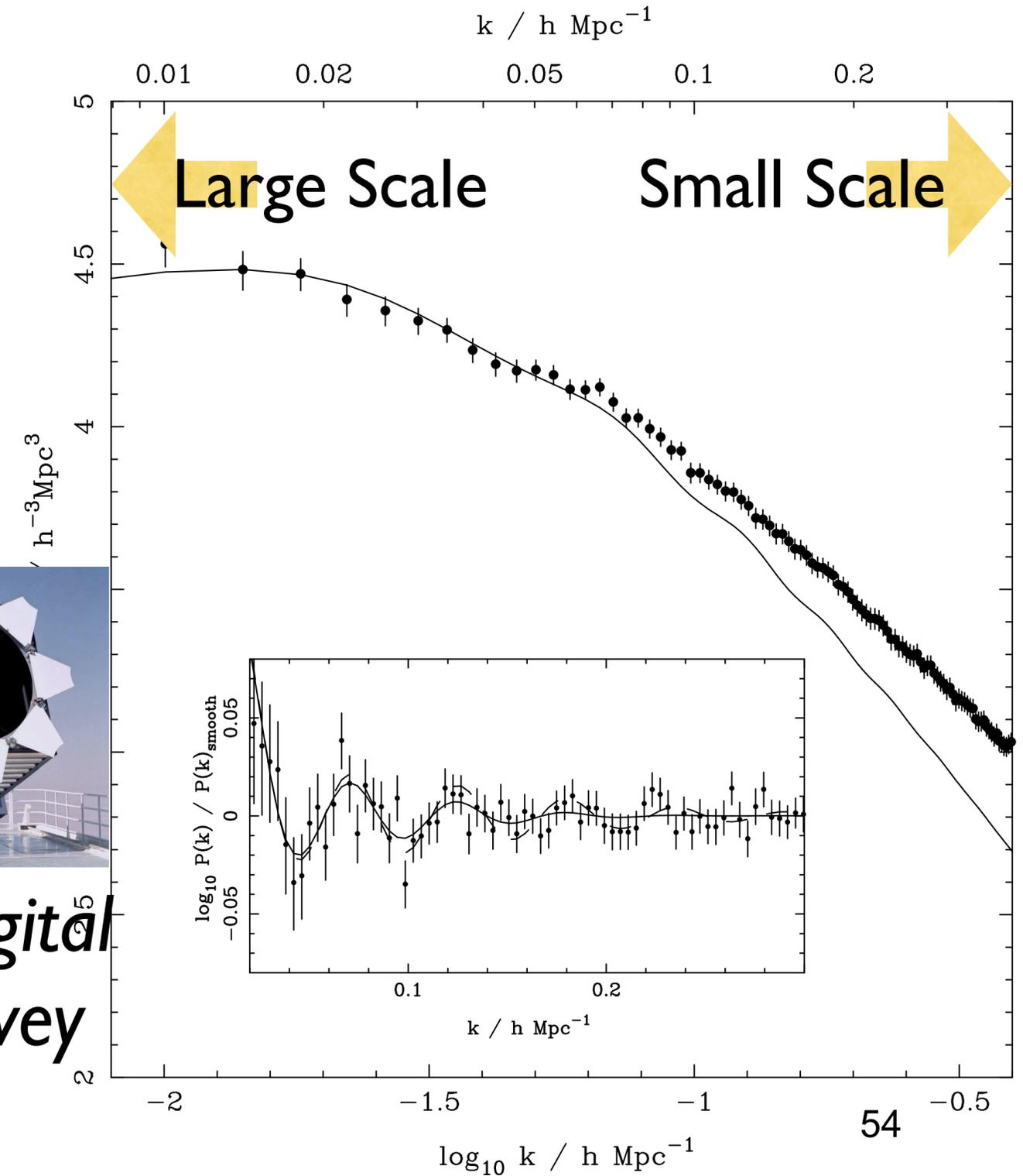


**Use 9.2-m HET to map the universe using
0.8M Lyman-alpha emitting galaxies
in $z=1.9-3.5$**

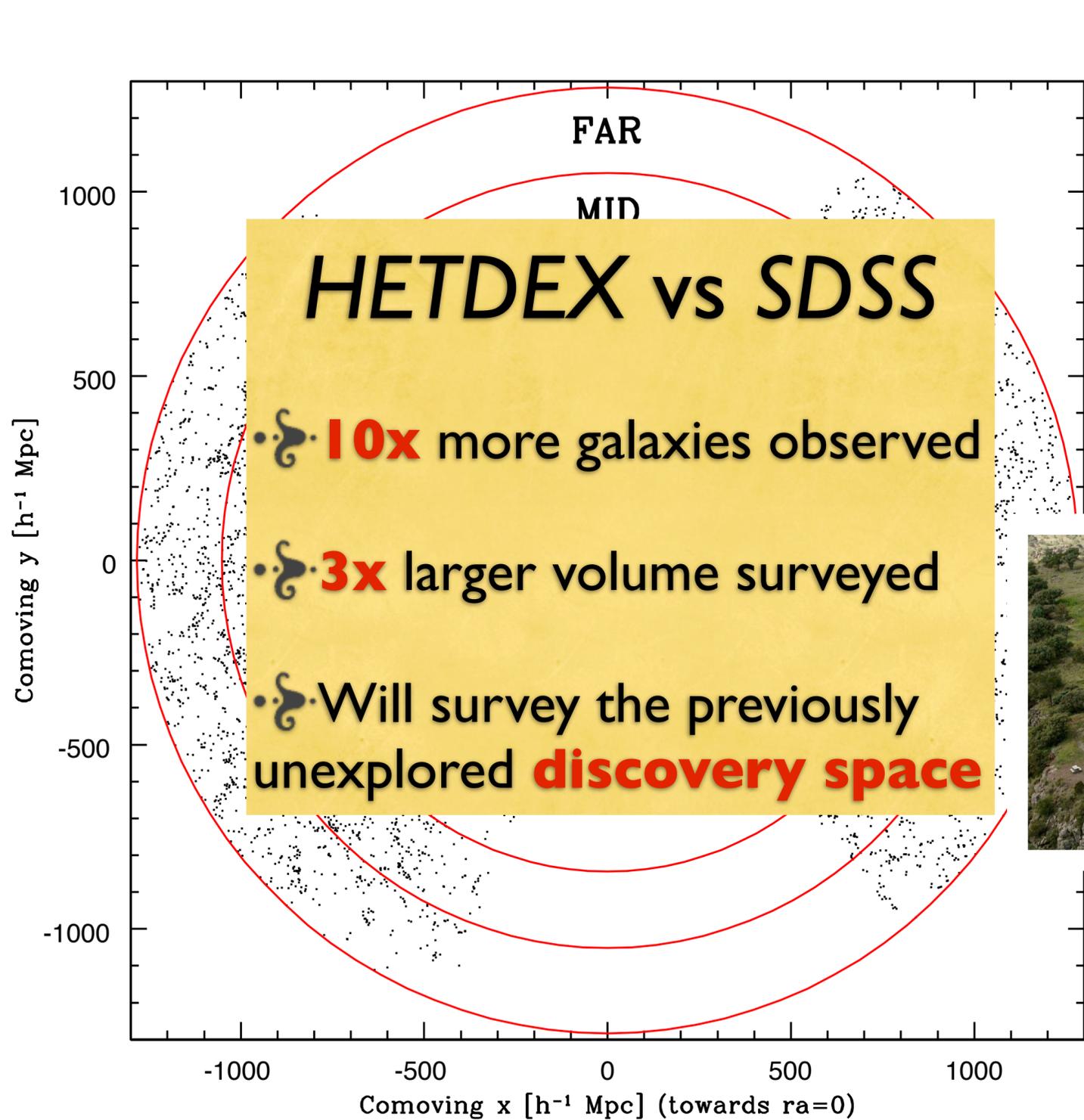
HETDEX: Sound Waves in the Distribution of Galaxies



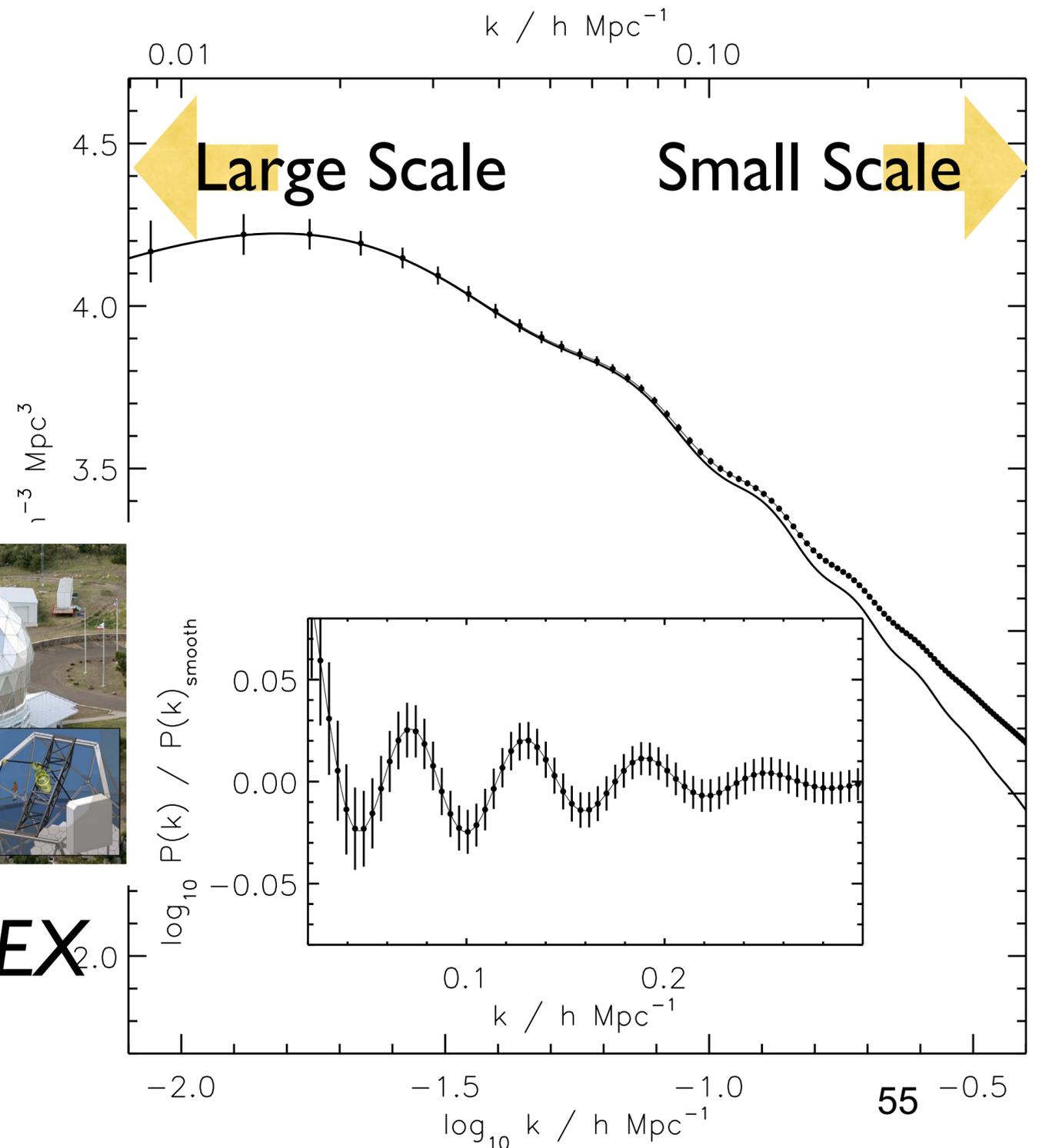
Sloan Digital Sky Survey



HETDEX: Sound Waves in the Distribution of Galaxies



HETDEX

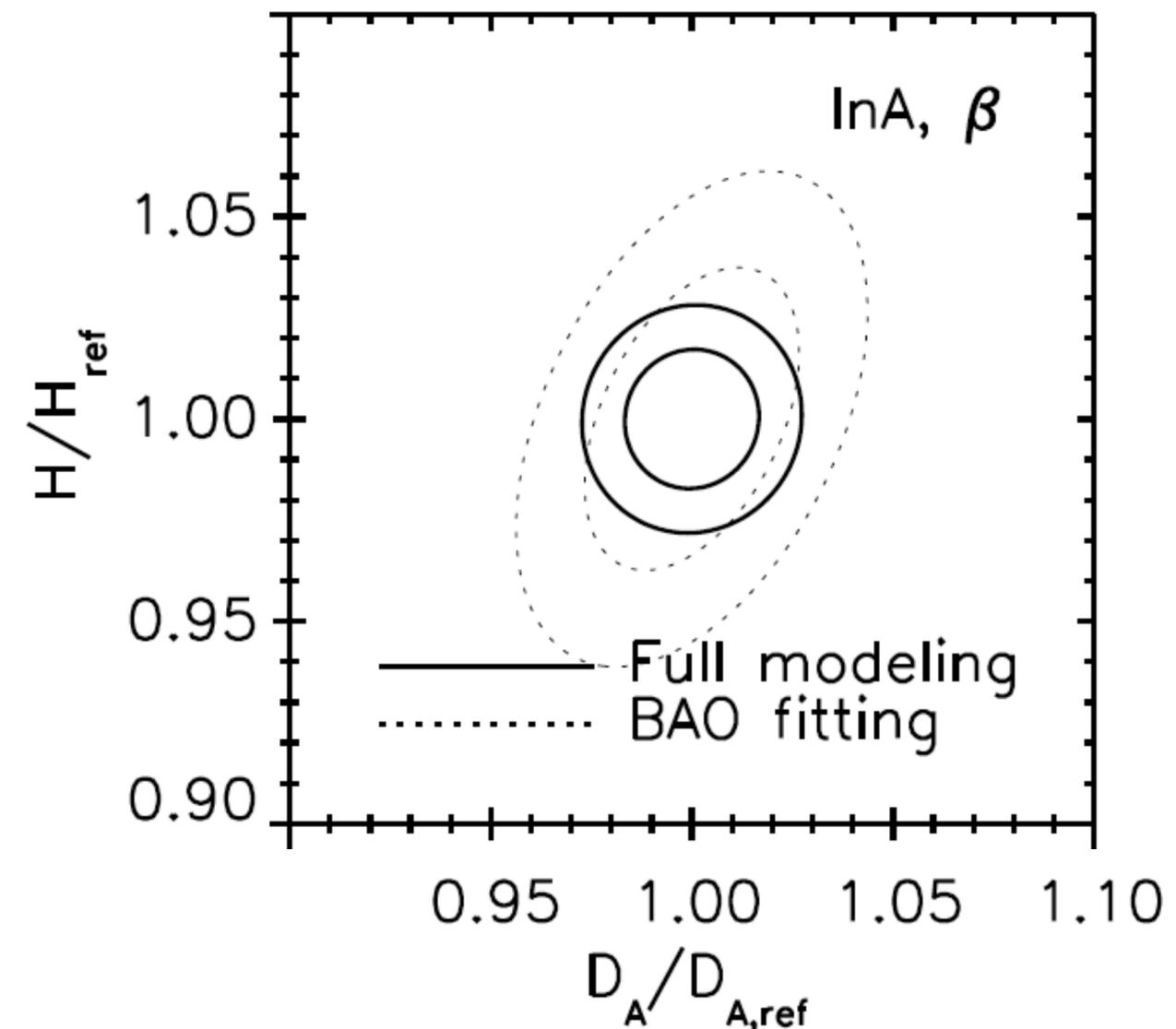


Beyond BAO

- BAOs capture only a **fraction** of the information contained in the galaxy power spectrum!
- The full usage of the 2-dimensional power spectrum leads to a *substantial* improvement in the precision of distance and expansion rate measurements.

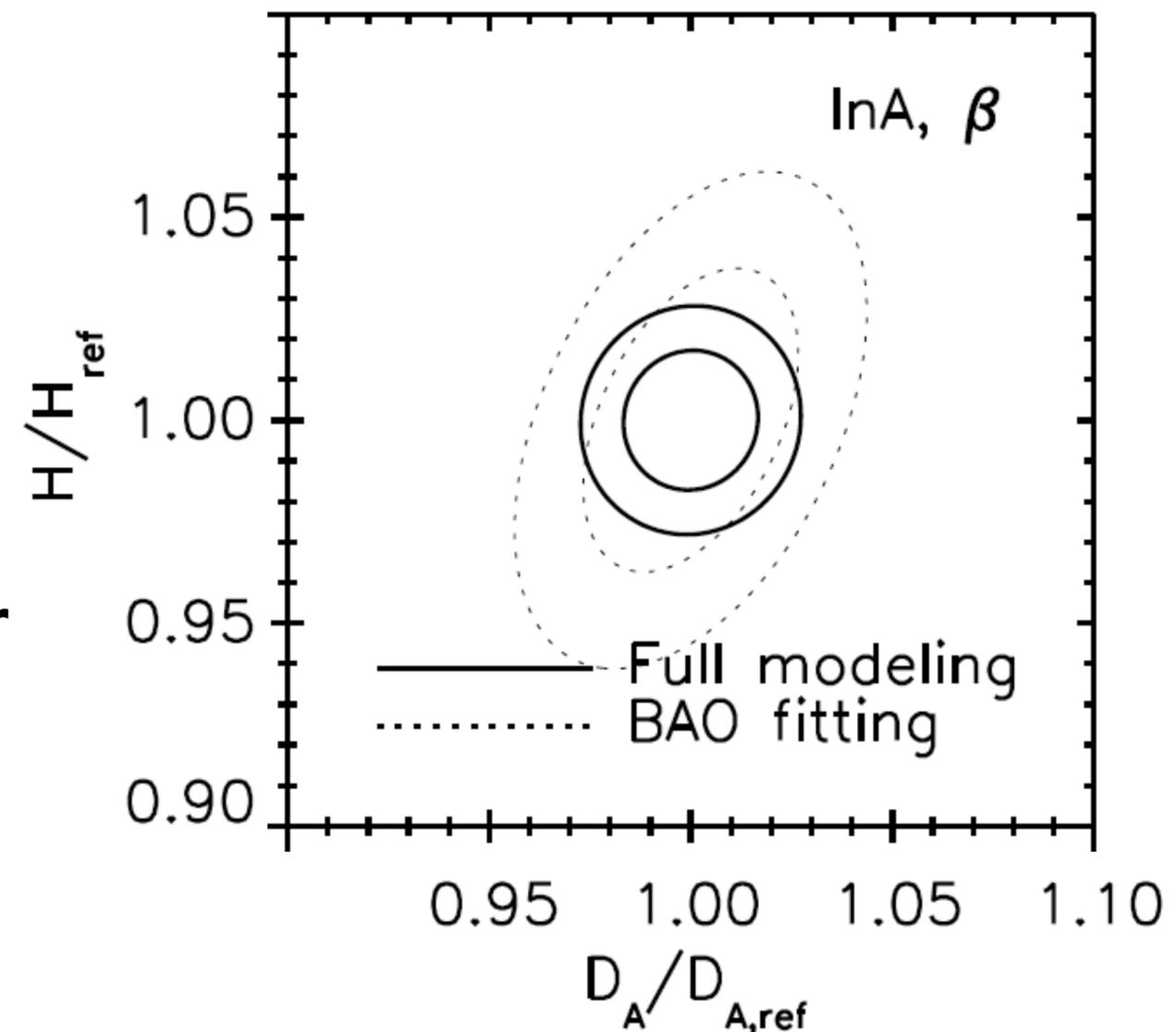
BAO vs Full Modeling

- Full modeling improves upon the determinations of D_A & H by more than a factor of two.
- On the D_A - H plane, the size of the ellipse shrinks by more than a factor of four.



Alcock-Paczynski: The Most Important Thing For HETDEX

- **Where does the improvement come from?**
- The Alcock-Paczynski test is the key. *This is the most important component for the success of the HETDEX survey.*



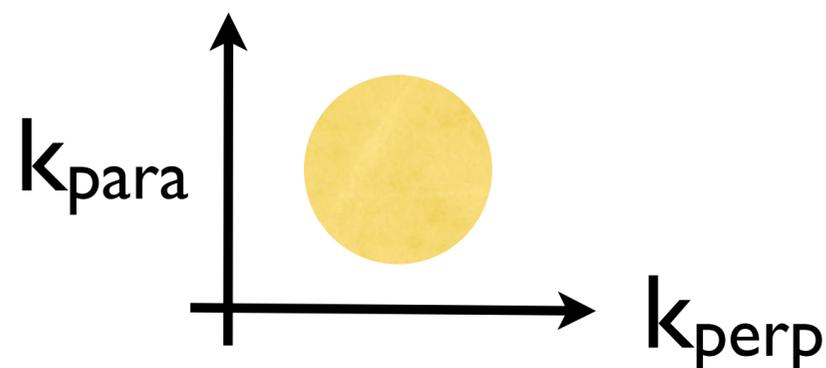
The AP Test: How That Works

- The key idea: (*in the absence of the redshift-space distortion - we will include this for the full analysis; we ignore it here for simplicity*), the distribution of the power should be **isotropic** in Fourier space.

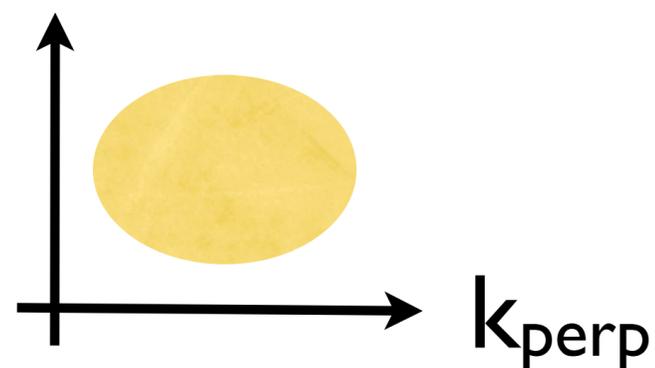
The AP Test: How That Works

- **D_A** : (RA, Dec) to the transverse separation, r_{perp} , to the transverse wavenumber
 - $k_{\text{perp}} = (2\pi)/r_{\text{perp}} = (2\pi)[\text{Angle on the sky}]/\mathbf{D_A}$
- **H** : redshifts to the parallel separation, r_{para} , to the parallel wavenumber
 - $k_{\text{para}} = (2\pi)/r_{\text{para}} = (2\pi)\mathbf{H}/(c\Delta z)$

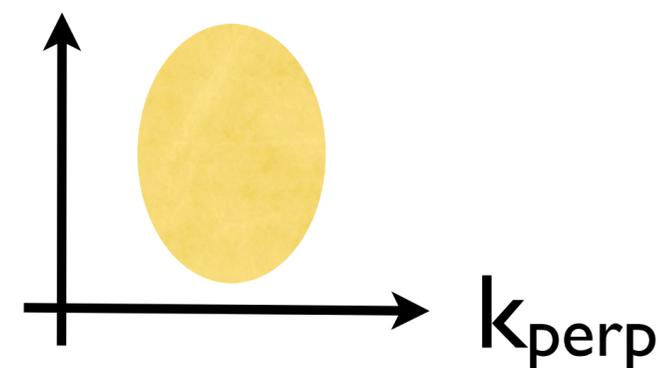
If D_A and H are correct:



If D_A is wrong:



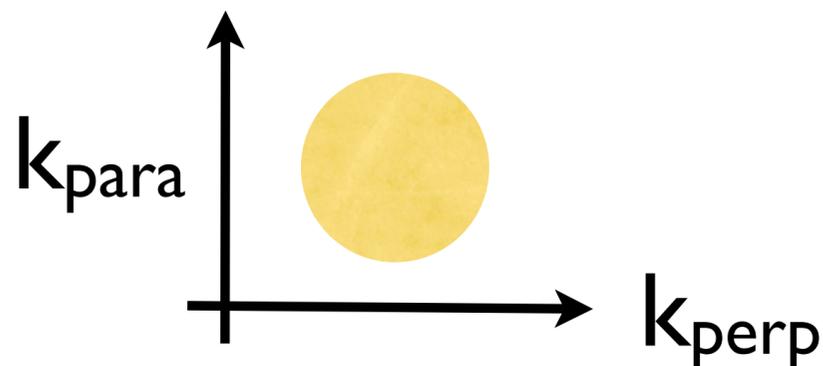
If H is wrong:



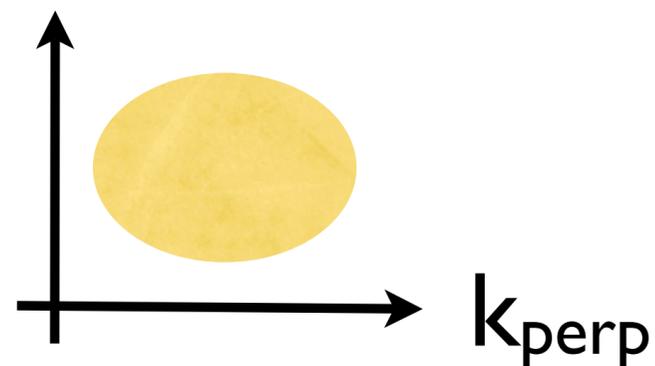
The AP Test: How That Works

- **D_A** : (RA, Dec) to the transverse separation, r_{perp} , to the transverse wavenumber
- $k_{\text{perp}} = (2\pi)/r_{\text{perp}} = (2\pi)[\text{Angle on the sky}]/\mathbf{D_A}$
- **H** : redshifts to the parallel separation, r_{para} , to the parallel wavenumber
- $k_{\text{para}} = (2\pi)/r_{\text{para}} = (2\pi)\mathbf{H}/(c\Delta z)$

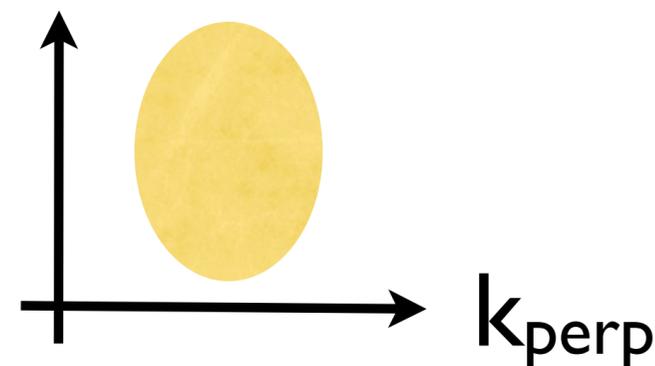
If D_A and H are correct:



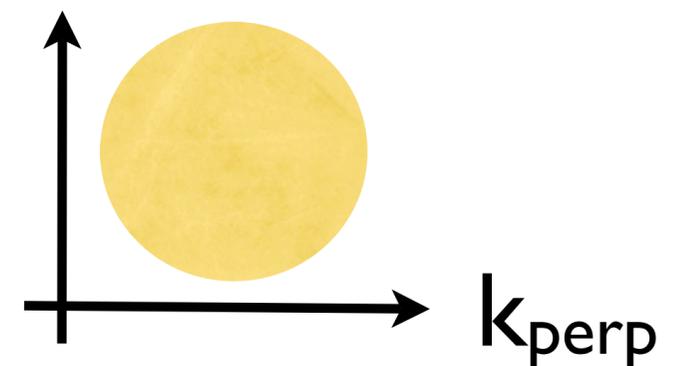
If D_A is wrong:



If H is wrong:

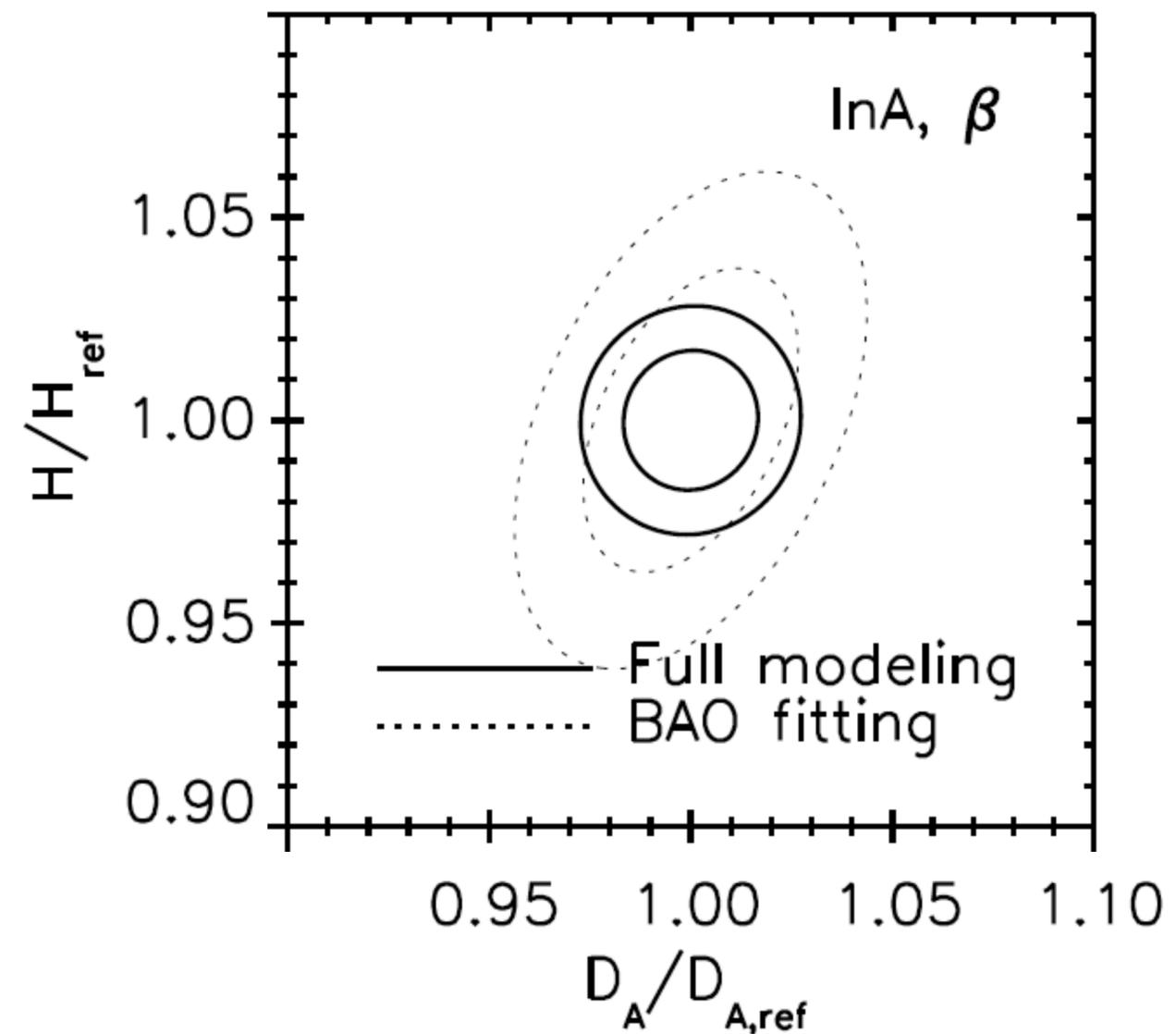


If D_A and H are wrong:



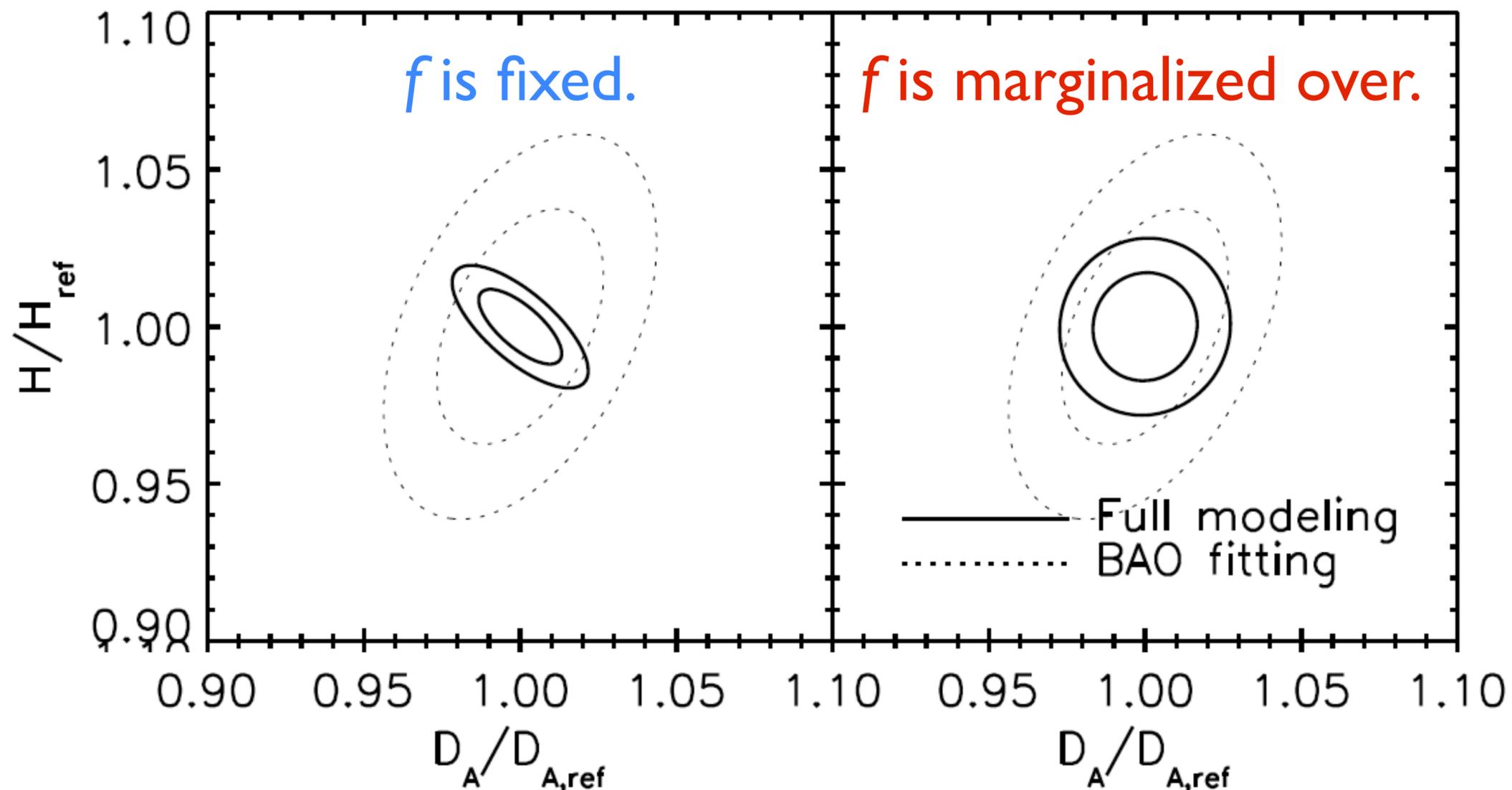
$D_A H$ from the AP test

- So, the AP test can't be used to determine D_A and H separately; however, it gives a measurement of **$D_A H$** .
- Combining this with the BAO information, and marginalizing over the redshift space distortion, we get the solid contours in the figure.



Redshift Space Distortion

- Both the AP test and the redshift space distortion make the distribution of the power anisotropic. Would it spoil the utility of this method?
- Some, but not all!



WMAP Amplitude Prior

- WMAP measures the amplitude of curvature perturbations at $z \sim 1090$. Let's call that R_k . The relation to the density fluctuation is

$$\delta_{m,\mathbf{k}}(z) = \frac{2k^3}{5H_0^2\Omega_m} \mathcal{R}_k T(k) D(k, z)$$

- Variance of R_k has been constrained as:

$$\Delta_{\mathcal{R}}^2(k_{WMAP}) = (2.208 \pm 0.078) \times 10^{-9} (68\% \text{ CL})$$

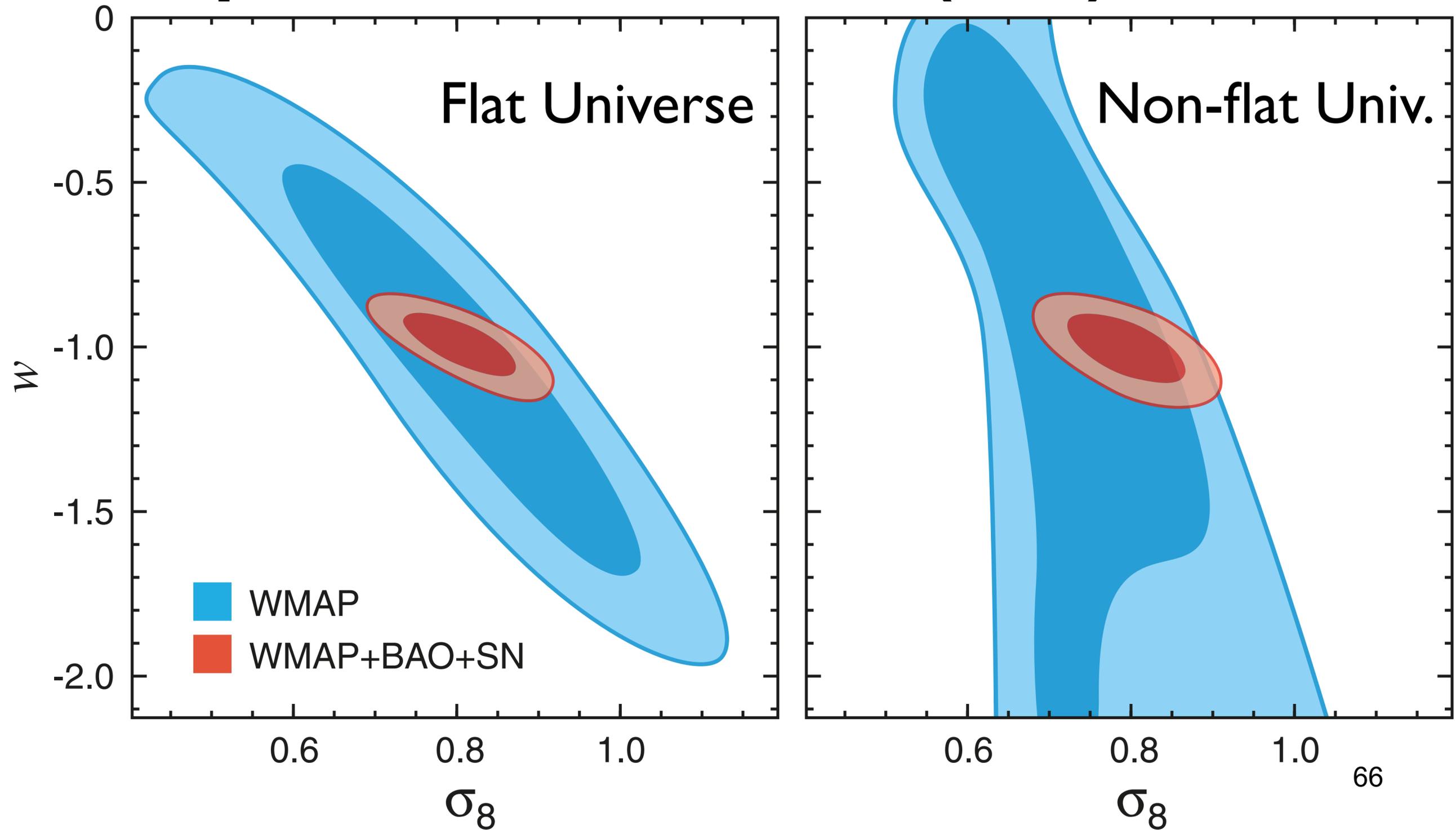
where $k_{WMAP} = 0.027 \text{ Mpc}^{-1}$

Then Solve This Diff. Equation...

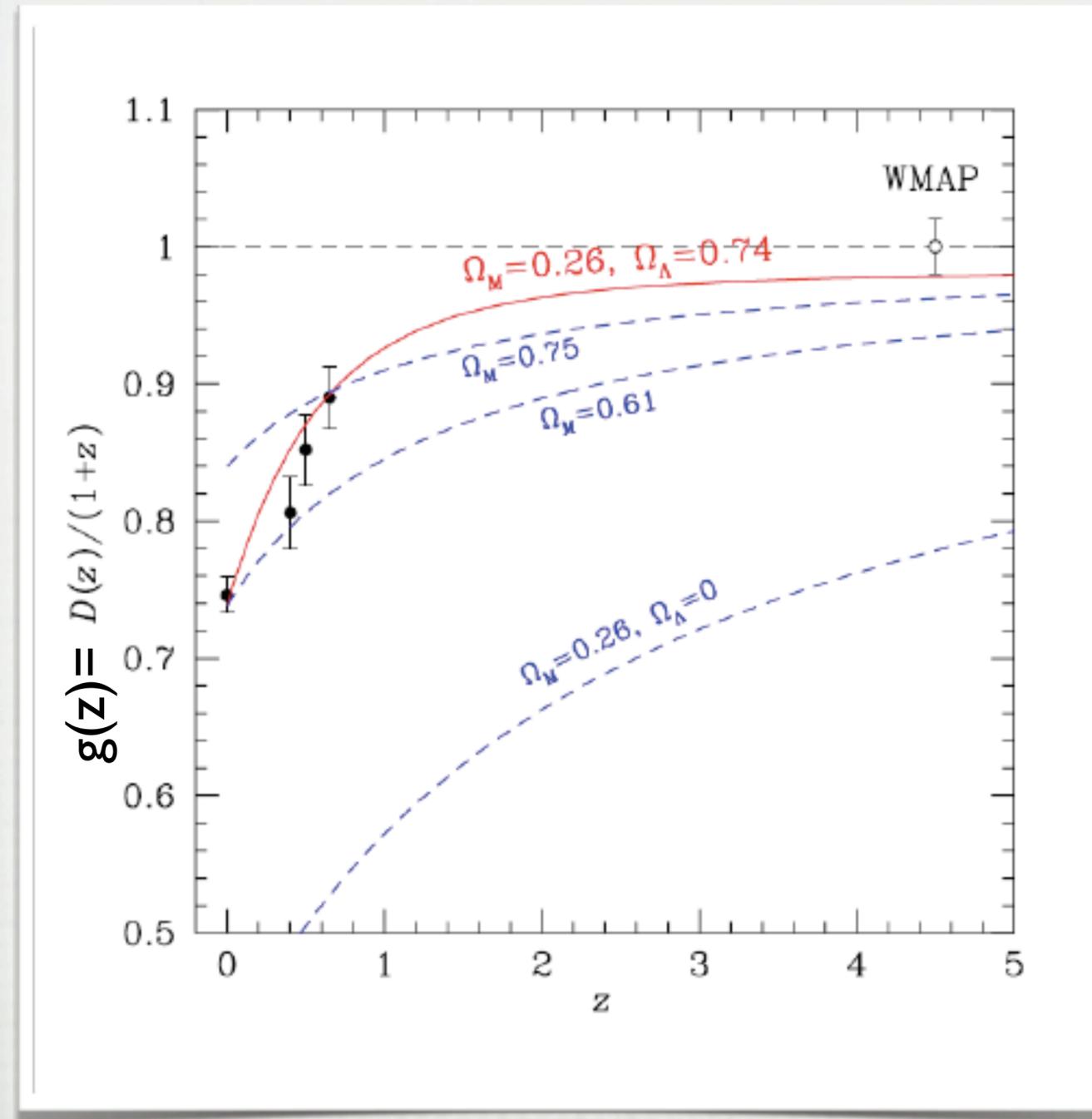
Ignoring the mass of neutrinos and modifications to gravity, one can obtain the growth rate by solving the following differential equation (Wang & Steinhardt 1998; Linder & Jenkins 2003): $\mathbf{g}(\mathbf{z})=(\mathbf{1}+\mathbf{z})\mathbf{D}(\mathbf{z})$

$$\frac{d^2 g}{d \ln a^2} + \left[\frac{5}{2} + \frac{1}{2} (\Omega_k(a) - 3 \mathbf{W}(a) \Omega_{de}(a)) \right] \frac{dg}{d \ln a} + \left[2\Omega_k(a) + \frac{3}{2} (1 - \mathbf{W}(a)) \Omega_{de}(a) \right] g(a) = 0, \quad (76)$$

Degeneracy Between Amplitude at $z=0$ (σ_8) and w



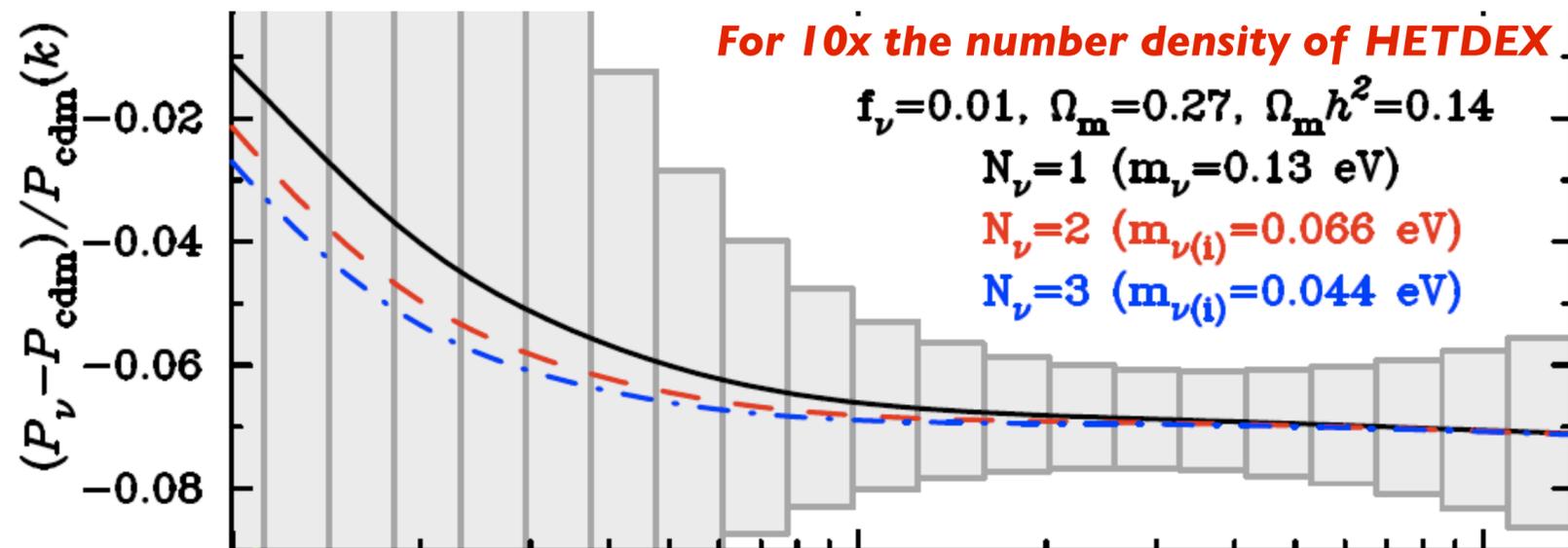
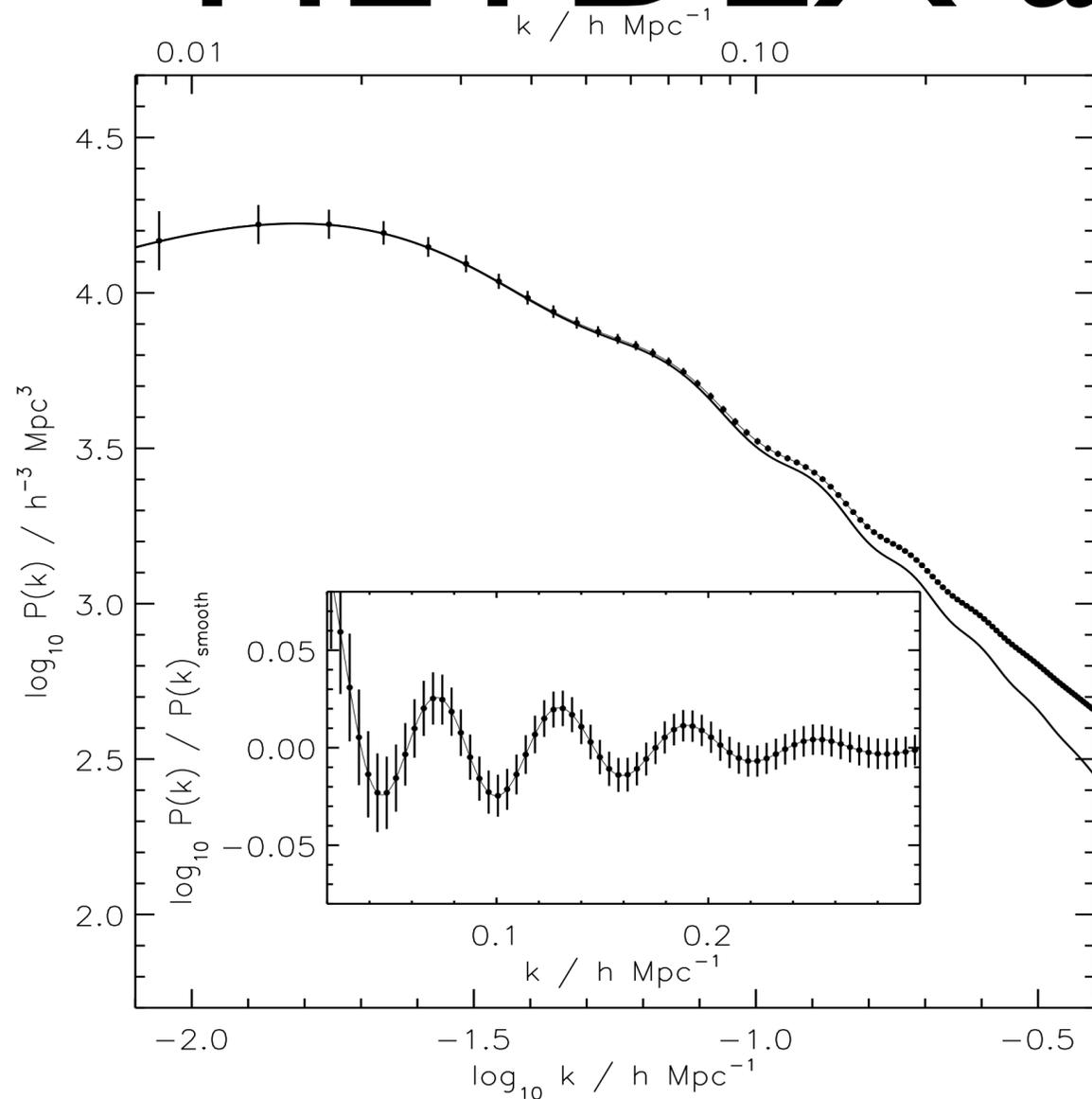
DETECTION OF Λ BY GROWTH HISTORY



Alexey Vikhlinin,
from a slide
presented at the
IPMU Dark Energy
Conference in
Japan, June 2009

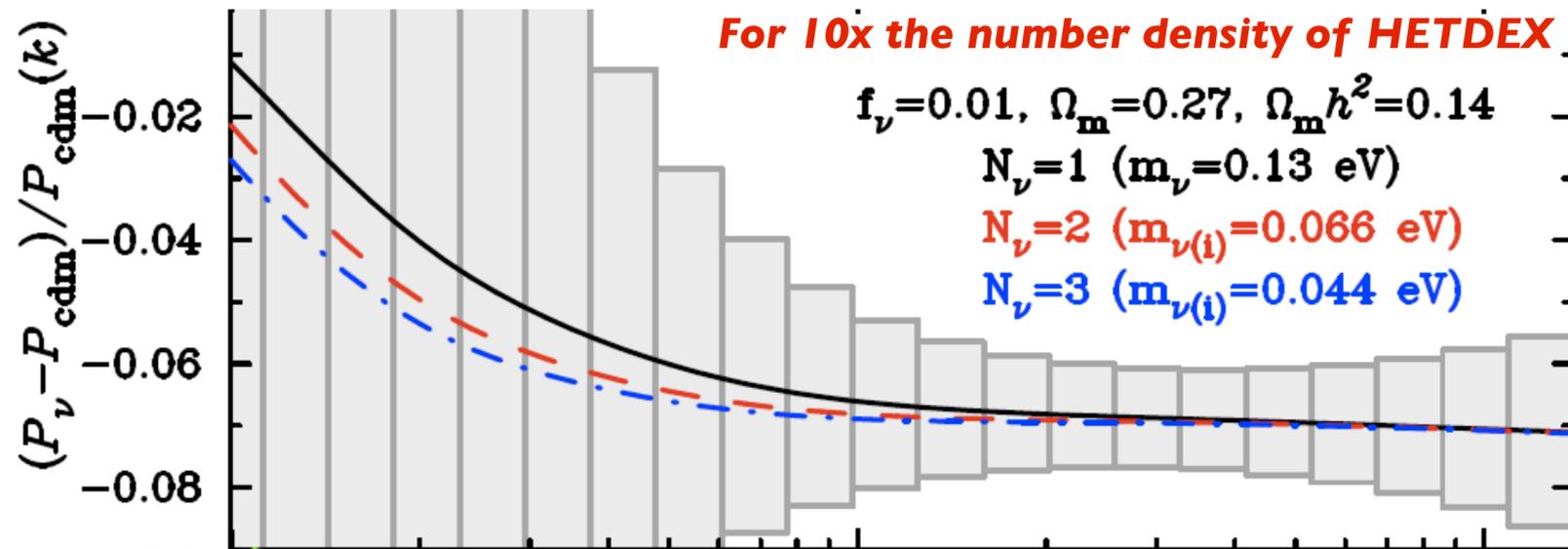
$\Lambda > 0$ at $\sim 3.5\sigma$ from perturbations growth only
at $\sim 5\sigma$ from growth + geometry

HETDEX and Neutrino Mass



- Neutrinos suppress the matter power spectrum on small scales ($k > 0.1 \text{ h Mpc}^{-1}$).
- A useful number to remember:
- For $\sum m_\nu = 0.1 \text{ eV}$, the power spectrum at $k > 0.1 \text{ h Mpc}^{-1}$ is suppressed by **$\sim 7\%$** .
- We can measure this easily!

Neutrino Mass and P(k)



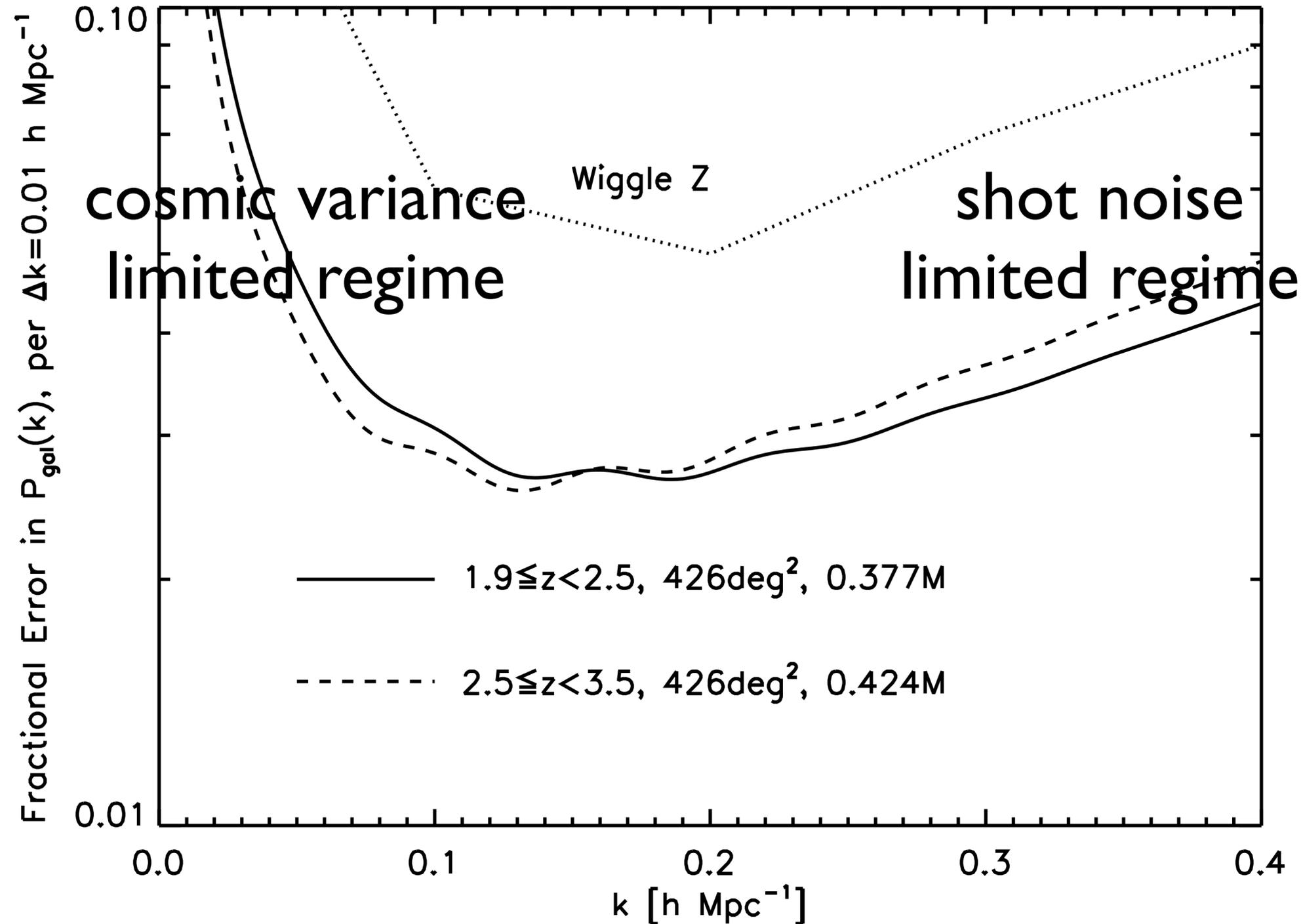
- Total neutrino mass: coming from the small scale

- $\Delta P/P \sim -8\Omega_\nu/\Omega_m = -[8/(\Omega_m h^2)]\sum m_\nu/(\dots)$

- Where the suppression begins depends on individual masses!

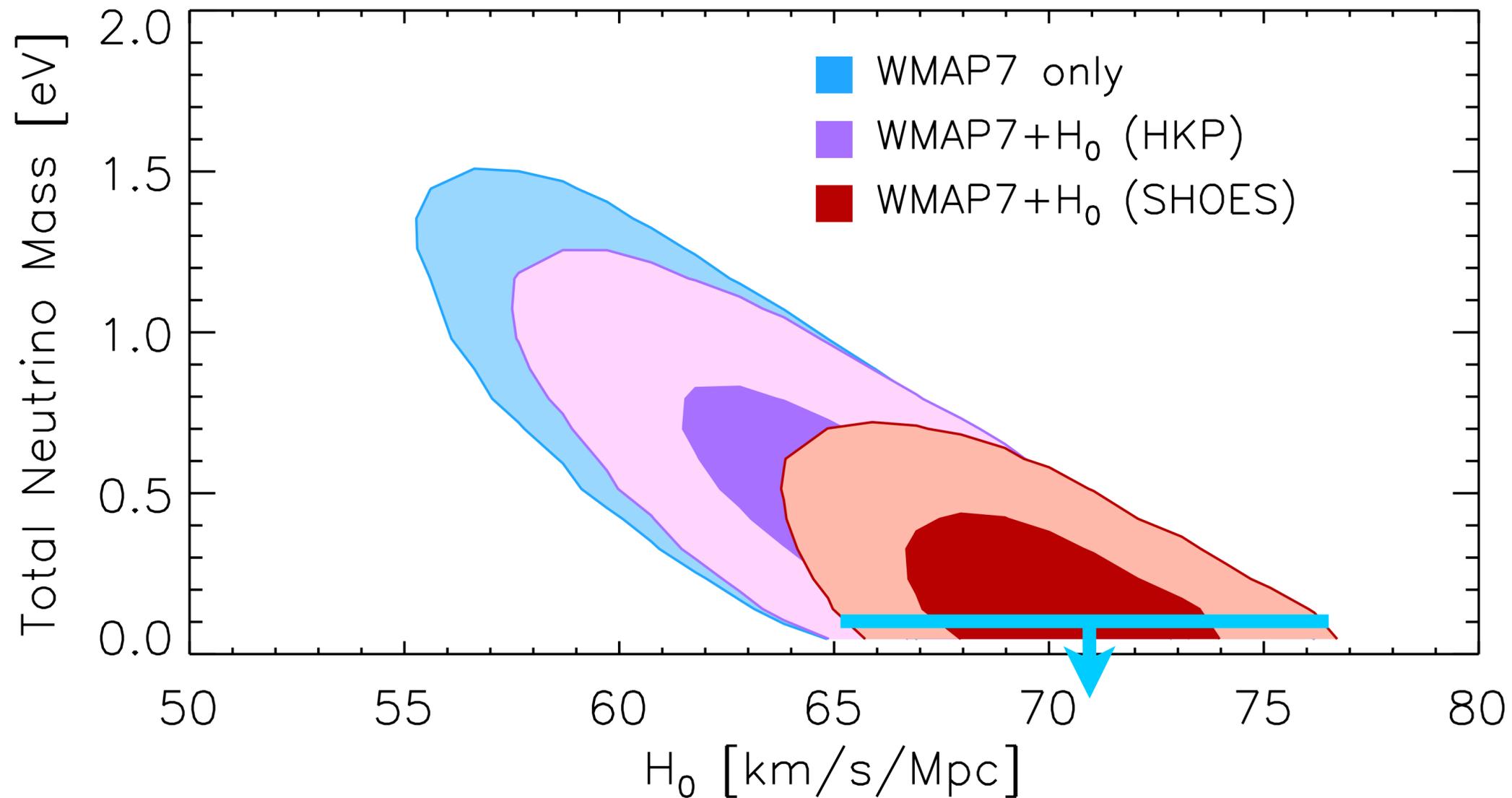
- $k_{\text{fs},i}(z) \equiv \sqrt{\frac{3}{2}} \frac{H(z)}{(1+z)\sigma_{v,i}(z)} \simeq \frac{0.677}{(1+z)^{1/2}} \left(\frac{m_{\nu,i}}{1 \text{ eV}}\right) \Omega_m^{1/2} h \text{ Mpc}^{-1}$

Expectation for HETDEX



- CV limited: error goes as $1/\sqrt{\text{volume}}$
- SN limited: error goes as $1/(\text{number density})/\sqrt{\text{volume}}$

Expected HETDEX Limit



- ~6x better than WMAP 7-year+ H_0

Summary

- Three (out of four) questions:
 - What is the physics of inflation?
 - $P(k)$ shape (esp, $dn/d\ln k$) and non-Gaussianity
 - What is the nature of dark energy?
 - $D_A(z)$, $H(z)$, growth of structure
 - What is the mass of neutrinos?
 - $P(k)$ shape
- CMB and large-scale structure observations can lead to major breakthroughs in any of the above questions.