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Cosmic Birefringence A New Probe of Dark Matter and Dark Energy based on

- Minami & EK, PRL, 125, 221301 (2020)
- EK, Nature Reviews Physics, 4 (2022)
- Eskilt & EK, PRD, 106, 063503 (2022)
- Diego-Palazuelos, et al., arXiv:2210.07644

Eiichiro Komatsu MPA Institute Seminar, November 14, 2022

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- Diego-Palazuelos, Eskilt, Minami, et al., PRL, 128, 091302 (2022)

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Highlights 2021

Highlights 2020

Highlights 2019

Highlights 2018

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Highlights 2015

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Author



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Director

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Original publications ay a menu

New analysis strengthens the hint of new physics in polarized radiation from the early Universe

JUNE 01, 2022

In 2020, a tantalizing hint of new physics violating "parity symmetry" was found in polarization data of the cosmic microwave background obtained with the Planck satellite at high frequencies. Based on the Planck data and a simplified assumption about the impact of the polarized dust emission in the Milky Way, the scientists reported a violation of the symmetry of the laws of physics under inversion of spatial coordinates with 99.2% confidence level. An international team led by MPA director Eiichiro Komatsu has now improved the analysis method. By considering the dust emission explicitly and using more data from not only Planck but also from WMAP the astrophysicists measured the parity-violating signal with 99.987% confidence level. If this should be confirmed in the future as a genuine cosmological signal, it would have profound implications for the fundamental physics behind dark matter, dark energy, and quantum gravity.



Photons of the cosmic microwave background (CMB), the afterglow of the primordial fireball Universe, are linearly polarized (Figure 1). This pattern can be used to search for new physics violating "parity" symmetry" - the symmetry of the laws of physics under an inversion of spatial coordinates. For example, electromagnetism works the same way whether one is in the original system or in a mirrored system with all spatial coordinates flipped. A violation of parity symmetry has only been observed in the weak interaction of the standard model of elementary particles and fields - so far. Can the Universe also violate parity symmetry?









ESA's Planck

Foreground-cleaned Temperature (smoothed)

Credit: ESA



Emitted 13.8 billions years ago





Foreground-cleaned Temperature (smoothed) + Polarisation

Credit: ESA

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Emitted 13.8 billions years ago



Standard Cosmological Model (ACDM) Requires New Physics Physics beyond Standard Model of elementary particles and fields

Dark Sector: What is dark matter (CDM)? What is dark energy (Λ)?

behind cosmic inflation?

Polarisation of the CMB may hold the key to the answers.

Early Universe: What powered the Big Bang? What is the fundamental physics



Standard Cosmological Model (ACDM) Requires New Physics Physics beyond Standard Model of elementary particles and fields

- **Dark Sector**: What is dark matter (CDM)? What is dark energy (Λ)?
 - Cosmic birefringence in CMB polarisation
- Early Universe: What powered the Big Bang? What is the fundamental physics behind cosmic inflation?
 - Imprint of primordial gravitational waves in CMB polarisation
- *Polarisation* of the CMB may hold the key to the answers.



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Review Article Published: 18 May 2022

New physics from the polarized light of the cosmic microwave background Key Words:

Eiichiro Komatsu

Nature Reviews Physics (2022) Cite this article

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Available also at arXiv:2202.13919

Cosmic Microwave Background (CMB) Polarization **Parity Symmetry**







Credit: WMAP Science Team The surface of "last scattering" by electrons (Scattering generates *polarisation*!)

What powered the Big Bang?

What is dark matter/energy?



How does the electromagnetic wave of the CMB propagate?

Today's topic

What is dark matter/energy?



How does the electromagnetic wave of the CMB propagate?

Today's topic

What is dark matter/energy?



Cosmic Birefringence The Universe filled with a "birefringent material"

- If the Universe is filled with a pseudoscalar field (e.g., an axion field) coupled to the electromagnetic tensor via a Chern-Simons coupling:
- Ni (1977); Turner & Widrow (1988) the effective Lagrangian for axion electrodynamics is $\mathcal{L} = -\frac{1}{2}\partial_{\mu}\theta\partial^{\mu}\theta - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + g_{a}\theta F_{\mu\nu}\tilde{F}^{\mu\nu}, \qquad (3.7)$

where g_a is a coupling constant of the order α , and the vacuum angle $\theta = \phi_a / f_a$ ($\phi_a = axion$ field). The equations

$$F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu} \qquad \sum_{\mu\nu}F_{\mu\nu}F^{\mu\nu} =$$

Carroll, Field & Jackiw (1990); Carroll & Field (1991); Harari & Sikivie (1992)





 $= 2(\mathbf{B} \cdot \mathbf{B} - \mathbf{E} \cdot \mathbf{E}) \qquad \sum F_{\mu\nu} \tilde{F}^{\mu\nu} = -4\mathbf{B} \cdot \mathbf{E}$ 11 Parity Even $\mu
u$ **Parity Odd**







Carroll, Field & Jackiw (1990); Carroll & Field (1991); Harari & Sikivie (1992)

Cosmic Birefringence This "axion" field can be The Universe filled with a "birefringent material" dark matter or dark energy!

- to the electromagnetic tensor via a Chern-Simon's coupling:
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• If the Universe is filled with a pseudoscalar field (e.g., an axion field) coupled



This term makes the phase velocities of right- and left-handed polarisation states of photons different, leading to rotation of the linear polarisation direction.





Standard Maxwell Theory Warm up (1)

• To isolate a transverse wave, we require $A_0=0$ and $div(A_i)=0$. Then, in vacuum,

$$\left(\frac{\partial^2}{\partial\eta^2} - \nabla^2\right) A_i(\eta, \mathbf{x}) = 0 \qquad ds^2 = a^2(-d\eta^2 + d\mathbf{x}^2)$$

 Go to Fourier space, choose the propagation direction of A_i to be in z-axis, and define right- and left-handed polarisation states as



$$A_1 \mp i A_2$$
 • A+: Right-handed states $\sqrt{2}$ • A-: Left-handed states







Standard Maxwell Theory Warm up (2)

• To isolate a transverse wave, we require $A_0=0$ and div $(A_i)=0$. Then, in vacuum,

and define right- and left-handed polarisation states as



Go to Fourier space, choose the propagation direction of A_i to be in z-axis,

$$A_1 \mp i A_2$$
 • A₊: Right-handed stated on the set of the set o

• A_: Left-handed state







Cosmic Birefringence Derivation (1)

Now, include the Chern-Simons term!

the effective Lagrangian for axion electrodynamics is $\mathcal{L} = -\frac{1}{2}\partial_{\mu}\theta\partial^{\mu}\theta - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + g_{a}\theta F_{\mu\nu}\tilde{F}^{\mu\nu}, \qquad (3.7)$ where g_a is a coupling constant of the order α , and the

 The equation of motion is modified to $\left(-\omega_{\pm}^2 + k^2\right)A_{\pm}(\eta) = 0$

Carroll, Field & Jackiw (1990); Carroll & Field (1991); Harari & Sikivie (1992)



$$\left(-\omega_{\pm}^{2}+k^{2}\pm4g_{a}k\theta'\right)A_{\pm}(\eta) = \frac{\omega_{\pm}^{2}}{k^{2}} = 1\pm\frac{4g_{a}\theta'}{k} \quad (\theta'=\partial\theta/\partial\eta)$$





Cosmic Birefringence Derivation (1)

Now, include the Chern-Simons term!

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$$\left(-\omega_{\pm}^{2}+k^{2}\pm4g_{a}k\theta'\right)A_{\pm}(\eta)=$$

$$\frac{\omega_{\pm}^{2}}{k^{2}}=1\pm\frac{4g_{a}\theta'}{k}=\left(1\pm\frac{2g_{a}\theta'}{k}\right)^{2}-\frac{4g_{a}\theta'}{k}$$

Cosmic Birefringence Derivation (1)

Now, include the Chern-Simons term!

the effective Lagrangian for axion electrodynamics is $\mathcal{L} = -\frac{1}{2}\partial_{\mu}\theta\partial^{\mu}\theta - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + g_{a}\theta F_{\mu\nu}\tilde{F}^{\mu\nu}, \qquad (3.7)$ where g_a is a coupling constant of the order α , and the

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$$\left(-\omega_{\pm}^{2}+k^{2}\pm4g_{a}k heta'
ight)A_{\pm}(\eta)=rac{\omega_{\pm}}{k}\simeq1\pmrac{2g_{a} heta'}{k}$$
 Phase velocities of and left-handed stars are slightly different of an elightly different of an elightly different of an elightly difference of an elightly differe

Cosmic Birefringence Derivation (2)

• With

$$rac{\omega_\pm}{k}\simeq 1\pm rac{2g_a heta'}{k}$$
 Phase velocities a and left-handed are slightly differentiated of the structure of

$$\label{eq:alpha} -\beta = \int d\eta \; \frac{\omega_+ - \omega_-}{2} = 2g_a \int d\eta \; \theta' = 2g_a \int dt$$
 The effect accumulates over the distance => CMB polarisation is sensitive to this effect.

Carroll, Field & Jackiw (1990); Carroll & Field (1991); Harari & Sikivie (1992)

of rightstates erent!

• The plane of linear polarisation rotates clockwise on the sky by an angle β :

Cosmic Birefringence Recap

- to the electromagnetic tensor via a Chern-Simons coupling:
- Ni (1977); Turner & Widrow (1988) the effective Lagrangian for axion electrodynamics is $\mathcal{L} = -\frac{1}{2}\partial_{\mu}\theta\partial^{\mu}\theta - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + g_{a}\theta F_{\mu\nu}\tilde{F}^{\mu\nu}, \qquad (3.7)$ where g_a is a coupling constant of the order α , and the vacuum angle $\theta = \phi_a / f_a$ ($\phi_a = axion$ field). The equations

Carroll, Field & Jackiw (1990); Carroll & Field (1991); Harari & Sikivie (1992)

This "axion" field can be dark matter or dark energy!

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Cosmic Birefringence Recap

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Carroll, Field & Jackiw (1990); Carroll & Field (1991); Harari & Sikivie (1992)

This "axion" field can be dark matter or dark energy!

• If the Universe is filled with a pseudoscalar field (e.g., an axion field) coupled

Motivation Why study the cosmic birefringence?

- The Universe's energy budget is dominated by two dark components:
 - Dark Matter
 - Dark Energy
- Either or both of these can be an axion-like field!
 - See Marsh (2016) and Ferreira (2020) for reviews.
- Thus, detection of parity-violating physics in polarisation of the cosmic microwave background can transform our understanding of Dark Matter/ Energy.

(Simpler) Motivation Why study the cosmic birefringence?

- 1957).

Let's look!

• We know that the weak interaction violates parity (Lee & Yang 1956; Wu et al.

Why should the laws of physics governing the Universe conserve parity?

B-mode : Polarisation directions are 45 degrees tilted w.r.t the wavenumber direction

B-mode : Polarisation directions are **45 degrees tilted** w.r.t the wavenumber direction

IMPORTANT": These "E and B modes" are jargons in the CMB community, and completely unrelated to the electric and magnetic fields of the electromagnetism!!

B-mode : Polarisation directions are **45 degrees tilted** w.r.t the wavenumber direction

Parity Flip E-mode remains the same, whereas B-mode changes the sign

Seljak & Zaldarriaga (1997); Kamionkowski, Kosowsky & Stebbins (1997)

 Two-point correlation functions invariant under the parity flip are

$$\langle E_{\boldsymbol{\ell}} E_{\boldsymbol{\ell}'}^* \rangle = (2\pi)^2 \delta_D^{(2)} (\boldsymbol{\ell} - \boldsymbol{\ell}') C$$
$$\langle B_{\boldsymbol{\ell}} B_{\boldsymbol{\ell}'}^* \rangle = (2\pi)^2 \delta_D^{(2)} (\boldsymbol{\ell} - \boldsymbol{\ell}') C$$
$$\langle T_{\boldsymbol{\ell}} E_{\boldsymbol{\ell}'}^* \rangle = \langle T_{\boldsymbol{\ell}}^* E_{\boldsymbol{\ell}'} \rangle = (2\pi)^2 \delta_D^{(2)} (\boldsymbol{\ell} - \boldsymbol{\ell}')$$

- The other combinations <TB> and <EB> are not invariant under the parity flip.
 - We can use these combinations to probe parity-violating physics (e.g., axions)

CMB Power Spectra Progress over the last 3 decades 10²

- This is the typical figure that you find in talks and lectures on CMB.
 - The temperature power spectrum and the E- and B-mode polarisation power spectra have been measured well.
- Our focus is the EB spectrum, which is not shown here.

Lue, Wang & Kamionkowski (1999); Feng et al. (2005, 2006); Liu, Lee & Ng (2006)

E-B mixing by rotation of the plane of linear polarisation

- Observed E- and B-mode polarisation, E^o and B_l^o, are related to those before rotation as
- $E^{o}_{\ell} \pm i B^{o}_{\ell} = (E_{\ell} \pm i B_{\ell}) e^{\pm 2i\beta}$
- which gives
- $E_{\ell}^{o} = E_{\ell} \cos(2\beta) B_{\ell} \sin(2\beta)$
- $B_{\ell}^{o} = E_{\ell} \sin(2\beta) + B_{\ell} \cos(2\beta)$

Lue, Wang & Kamionkowski (1999); Feng et al. (2005, 2006); Liu, Lee & Ng (2006) Zhao et al. (2015) Searching for the birefringence

- Computing observed difference between EE and BB spectra,
 - $C_{\ell}^{EE,\text{obs}} = C_{\ell}^{EE} \cos^2(2\beta) + C_{\ell}^{BB} \sin^2(2\beta) C_{\ell}^{EB} \sin(4\beta)$ $C_{\ell}^{BB,\text{obs}} = C_{\ell}^{EE} \sin^2(2\beta) + C_{\ell}^{BB} \cos^2(2\beta) + C_{\ell}^{EB} \sin(4\beta)$ $C_{\ell}^{EE,\text{obs}} - C_{\ell}^{BB,\text{obs}} = (C_{\ell}^{EE} - C_{\ell}^{BB})\cos(4\beta) - 2C_{\ell}^{EB}\sin(4\beta)$
- We find

$$C_{\ell}^{EB,\text{obs}} = \frac{1}{2} (C_{\ell}^{EE} - C_{\ell}^{BB}) \sin \left(\frac{1}{2} (C_{\ell}^{EE,\text{obs}} - C_{\ell}^{BB,\text{o}}) - \frac{1}{2} (C_{\ell}^{EE,\text{obs}} - C_{\ell}^{BB,\text{o}}) \right)$$

 $(4\beta) + C_{\ell}^{EB} \cos(4\beta)$ $\frac{C_{\ell}^{EB}}{\cos(4\beta)}$ $^{\text{bs}}) \tan(4\beta)$

EB is generated by the *difference* between EE and BB spectra.

CMB Power Spectra EE >> **BB**!

トン

- In our Universe, CMB EE is much greater than BB. This makes CMB sensitive to birefringence.
- This is the typical figure that you find in talks and lectures on CMB.
 - The temperature power spectrum and the E- and B-mode polarisation power spectra have been measured well.
- Our focus is the EB spectrum, which is not shown here.

This is the EB power spectrum (WMAP+Planck) Nearly full-sky data (92% of the sky)

Eskilt & EK, arXiv:2205.13962

This is the EB power spectrum (WMAP+Planck) Galactic plane removed (62% of the sky)

Eskilt & EK, arXiv:2205.13962

Cosmic Birefringence fits well(?) Nearly full-sky data (92% of the sky)

Eskilt & EK, arXiv:2205.13962

Cosmic Birefringence fits well(?) Galactic plane removed (62% of the sky)

Eskilt & EK, arXiv:2205.13962

Signal is robust with respect

The Biggest Problem: Miscalibration of detectors

Wu et al. (2009); Komatsu et al. (2011); Keating, Shimon & Yadav (2012) Impact of miscalibration of polarisation angles **Cosmic or Instrumental?**

- to the sky coordinates (and we did not know it)?

(but we do not know it)

 Is the plane of linear polarisation rotated by the genuine cosmic birefringence effect, or simply because the polarisation-sensitive directions of detectors are rotated with respect

If the detectors are rotated by α , it seems that we can measure only the SUM $\alpha+\beta$.

The past measurements The quoted uncertainties are all statistical only (68%CL)

- $\alpha + \beta = -6.0 \pm 4.0 \text{ deg}$ (Feng et al. 2006)
- $\alpha+\beta = -1.1 \pm 1.4$ deg (WMAP Collaboration, Komatsu et al. 2009; 2011)
- $\alpha+\beta = 0.55 \pm 0.82$ deg (QUaD Collaboration, Wu et al. 2009)
- $\alpha + \beta = 0.31 \pm 0.05$ deg (Planck Collaboration 2016)
- $\alpha + \beta = -0.61 \pm 0.22$ deg (POLARBEAR Collaboration 2020)
- $\alpha+\beta = 0.63 \pm 0.04 \text{ deg}$ (SPT Collaboration, Bianchini et al. 2020)
- $\alpha+\beta = 0.12 \pm 0.06$ deg (ACT Collaboration, Namikawa et al. 2020)
- $\alpha+\beta = 0.07 \pm 0.09$ deg (ACT Collaboration, Choi et al. 2020)

first measurement

The past measurements Now including the estimated systematic errors on a • $\beta = -6.0 \pm 4.0 \pm ??$ deg (Feng et al. 2006)

- $\beta = -1.1 \pm 1.4 \pm 1.5$ deg (WMAP Collaboration, Komatsu et al. 2009; 2011)
- $\beta = 0.55 \pm 0.82 \pm 0.55$ deg (QUaD Collaboration, Wu et al. 2009)
- •
- $\beta = 0.31 \pm 0.05 \pm 0.28$ deg (Planck Collaboration 2016)
- $\beta = -0.61 \pm 0.22 \pm ??$ deg (POLARBEAR Collaboration 2020)
- $\beta = 0.63 \pm 0.04 \pm ??$ deg (SPT Collaboration, Bianchini et al. 2020)
- $\beta = 0.12 \pm 0.06 \pm ??$ deg (ACT Collaboration, Namikawa et al. 2020)
- $\beta = 0.07 \pm 0.09 \pm ??$ deg (ACT Collaboration, Choi et al. 2020)

Uncertainty in the calibration of a has been the major limitation

$f_{sky} = 0.90$ CO+PS+5% (1deg apodization)

CO+PS+20% (1deg apodization)

CO+PS (1deg apodization)

$f_{sky} = 0.93^{-1}$ = nearly full sky

CO+PS+10% (1deg apodization) $f_{sky} = 0.85$

. .

 $f_{sky} = 0.63$

CO+PS+30% (1deg apodization)

....

The Key Idea: The polarised Galactic foreground emission as a calibrator

ESA's Planck

Directions of the magnetic field inferred from polarisation of the thermal dust emission in the Milky Way

Credit: ESA

Polarised dust emission within our Milky Way!

Emitted "right there" - it would not be affected by the cosmic birefringence.

Foreground-cleaned Temperature (smoothed) + Polarisation

Credit: ESA

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Emitted 13.8 billions years ago

Searching for the birefringence Including the miscalibration angle

• Idea: Miscalibration of the polarization angle a rotates both the foreground and CMB, but β affects only the CMB. $E_{\ell,m}^{o} = E_{\ell,m}^{fg} \cos(2\alpha) - B_{\ell,m}^{fg} \sin(2\alpha) + E_{\ell,m}^{CMB} \cos(2\alpha + 2\beta) - B_{\ell,m}^{CMB} \sin(2\alpha + 2\beta) + E_{\ell,m}^{N}$ $B_{\ell,m}^{o} = E_{\ell,m}^{fg} \sin(2\alpha) + B_{\ell,m}^{fg} \cos(2\alpha) + E_{\ell,m}^{CMB} \sin(2\alpha + 2\beta) + B_{\ell,m}^{CMB} \cos(2\alpha + 2\beta) + B_{\ell,m}^{N}$ • Thus,

$$\begin{split} \langle C_{\ell}^{EB,\mathrm{o}} \rangle = & \frac{\tan(4\alpha)}{2} \left(\frac{\langle C_{\ell}^{EE,\mathrm{o}} \rangle - \langle C_{\ell}^{BB,\mathrm{o}} \rangle}{2} \right) + \frac{\sin(4\beta)}{2\cos(4\alpha)} \left(\frac{\langle C_{\ell}^{EE,\mathrm{CMB}} \rangle - \langle C_{\ell}^{BB,\mathrm{CM}} \rangle}{\mathrm{known\ accurately}} \right) \\ & + \frac{1}{\cos(4\alpha)} \langle C_{\ell}^{EB,\mathrm{fg}} \rangle + \frac{\cos(4\beta)}{\cos(4\alpha)} \langle C_{\ell}^{EB,\mathrm{CMB}} \rangle \,. \end{split}$$

How does it work?

Minami et al. (2019)

- When the data are dominated by CMB, the sum of two angles, $\alpha + \beta$, is determined precisely.
 - This is the diagonal line.
- The foreground determines α with some uncertainty, breaking the degeneracy. Then $\sigma(\beta) \sim \sigma(\alpha)$ because $\sigma(\alpha+\beta) << \sigma(\alpha).$
- When the data are dominated by the foreground, it can determine a but not β due to the lack of sensitivity to the CMB.

From Planck HFI Data (Minami & EK 2020)

 $\langle C_{\ell}^{EB,o} \rangle = \frac{\tan(4\alpha)}{2}$ $\langle C_{\ell}^{EE, \text{CMB}_{\chi}} \rangle$ $sin(4\beta)$ $\langle \langle C_{\ell}^{EE,\mathrm{o}} \rangle - \langle C_{\ell}^{BB,\mathrm{o}} \rangle \rangle$ $2\cos(4\alpha)$

- Can we see β by eyes?
- First, take a look at the observed EE-BB spectra.
 - Red: Total
 - Blue: The best-fitting CMB model
 - The difference is due to the FG (and maybe unknown systematics)

From Planck HFI Data (Minami & EK 2020)

 $\langle C_\ell^{EB, \mathrm{o}_\lambda} \rangle$ $\frac{\tan(4\alpha)}{2}\left(\langle C_{\ell}^{EE,o}\rangle - \langle C_{\ell}^{BB,o}\rangle\right)$ $(\langle C_{\ell}^{EE, \text{CMB}} \rangle)$ $\sin(4\beta)$

- Can we see β by eyes?
 - Red: The signal attributed to the miscalibration angle, α_v
 - Blue: The signal attributed to the cosmic birefringence, β
- Red + Blue is the best-fitting model for explaining the data points

Angles	Results (deg)
β	0.35 ± 0.14
$lpha_{100}$	-0.28 ± 0.13
$lpha_{143}$	0.07 ± 0.12
$lpha_{217}$	-0.07 ± 0.11
$lpha_{353}$	-0.09 ± 0.11

Assumption for the baseline result What about the intrinsic EB correlation of the foreground emission?

$$\langle C_{\ell}^{EB,o} \rangle = \frac{\tan(4\alpha)}{2} \left(\langle C_{\ell}^{EE,o} \rangle - \langle C_{\ell}^{BB,o} \rangle \right) + \frac{\sin(4\beta)}{2\cos(4\alpha)} \left(\langle C_{\ell}^{EE,CMB} \rangle - \langle C_{\ell}^{BB,CMI} + \frac{1}{\cos(4\alpha)} \left(\langle C_{\ell}^{EB,fg} \rangle \right) + \frac{\cos(4\beta)}{\cos(4\alpha)} \left(\langle C_{\ell}^{EB,CMB} \rangle \right).$$

- but we take into account the foreground (dust) EB, $\langle C_{\ell}^{EB,fg} \rangle$
 - TB/TE measured from the data. Not really a modelling.

• For the baseline result, we ignore the intrinsic EB correlation of the CMB, $\langle C_{\ell}^{EB,CMB} \rangle$

• We account for the dust EB by assuming that EB/EE is proportional to

Including the foreground EB Introducing a new angle, "y"

• When we do not ignore the intrinsic foreground EB, the observed foreground EB (including the miscalibration angle contribution, α) is given by

$$C_{\ell}^{EB,\mathrm{FG,o}} = \frac{1}{2}\sin(4\alpha)\left(C_{\ell}^{EE,\mathrm{FG}} - C_{\ell}^{BB,\mathrm{FG}}\right) + \frac{C_{\ell}^{EB,\mathrm{FG}}\cos(4\alpha)}{\mathrm{new \ term}}$$

- Using a formula for trigonometric functions, $A\sin\varphi + B\cos\varphi = \sqrt{A^2 + B}$
- We obtain

 $C_{\ell}^{EB,\mathrm{FG,o}} = \sqrt{J_{\ell}^2 + (C_{\ell}^{EB,\mathrm{FG}})}$

Minami et al. (2019); Minami & Komatsu (2020b); Diego-Palazuelos, Eskilt, et al. (2022)

$$B^{2}\sin(\varphi + \theta), \quad \tan \theta = B/A$$

$$\underline{\alpha} \rightarrow \alpha + \gamma \qquad J_{\ell} \equiv \frac{1}{2} \left(C_{\ell}^{EE, FG} - C_{\ell}^{B} + \Omega \right)$$

 $(1)^{2}\sin(4\alpha + 4\gamma_{\ell}) \quad \tan(4\gamma_{\ell}) \equiv C_{\ell}^{EB, FG}$

Relating EB to TB

• A generic approach:

$C_{\rho}^{EB,dust}$

Minami et al. (2019); Minami & Komatsu (2020b); Diego-Palazuelos, Eskilt, et al. (2022)

$C_{n}TB, dust$

$\pi TE, dust$

Relating EB to TB

- How do we model the new angle

 - Then $\gamma EE, dust$

Clark et al. (2021); Diego-Palazuelos, Eskilt, et al. (2022)

$$\begin{array}{l} & \frac{C_{\ell}^{EB,{\rm dust}}}{C_{\ell}^{EE,{\rm dust}}} \propto \frac{C_{\ell}^{TB,{\rm d}}}{C_{\ell}^{TE,{\rm d}}} \end{array} \end{array} \\ \\ \textbf{v?} & \frac{C_{\ell}^{EE,{\rm dust}}}{C_{\ell}^{EE,{\rm dust}}} \end{array} \end{array} \end{array} \end{array}$$

• Our ansatz, motivated by a physical consideration of Clark et al. (2021):

 $C_{\ell}^{EB,\text{dust}} = A_{\ell} C_{\ell}^{EE,\text{dust}} \sin(4\psi_{\ell}^{\text{dust}})$

Free I-dependent amplitude parameters $\psi_{\ell}^{\text{dust}} = \frac{1}{2} \arctan(C_{\ell}^{TB,\text{dust}}/C_{\ell}^{TE,\text{dust}})$

 $\gamma^{T'B}, \mathrm{dust}$ $\tilde{C}^{EE,\text{dust}}_{\ell} - C^{BB,\text{dust}}_{\ell} \overline{C^{TE,\text{dust}}_{\ell}}$

for small angles.

Miscalibration angles (WMAP and Planck) Nearly full-sky data (92% of the sky)

Eskilt & EK, arXiv:2205.13962

LFI HFI **WMAP** β

- The angles are all over the place, and are well within the quoted calibration uncertainty of instruments.
 - 1.5 deg for WMAP
 - 1 deg for Planck
- They cancel!
 - The power of adding independent datasets.

1.5

No frequency dependence is found **Consistent with the expectation from cosmic birefringence**

 $0.33^{\circ} \pm 0.10^{\circ}$

Eskilt (2022); Eskilt & EK, arXiv:2205.13962

- No evidence for frequency dependence:
 - For β~(v/150GHz)ⁿ $n = -0.20^{+0.41} - 0.39 (68\% CL)$
 - Faraday rotation (n=-2) is disfavoured.

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Conclusion $\beta = 0.34 \pm 0.09 \text{ deg } (68\% \text{CL}; \text{ nearly full sky})$

- Robust against the sky fraction used for the analysis.
- No evidence for frequency dependence of β.
 - Consistent with a cosmological signal.
- Good news: The impact of the known instrumental systematics is negligible.
 - We found this using the Planck simulations (*Diego-Palazuelos et al.*, arXiv:2210.07655).
- If the measured β is confirmed as cosmological, it would have profound implications for the fundamental physics behind dark matter and energy.
- It is fair to say that there is something in the Planck data. No evidence for significant impacts of the Galactic foreground or the known systematics.
 - But we keep investigating!

Relating EB to TB A physical justification for this assumption $ho^{EE,{ m dust}}$

- How do we model the new angle y physically?
 - We don't really know for sure yet. This is a fascinating opportunity for Galactic science!
- Nonetheless, there may be a clue from the dust TB correlation.
 - **Discovery of a non-zero (positive) dust TB correlation** by the Planck collaboration was a surprise.
 - We still do not know its origin (see Huffenberger et al. 2020 and Clark et al. 2021 for the first attempts to explain it).
 - However, it seems reasonable to relate the possible dust EB correlation to the dust TB correlation.

TE, TB, and EB correlation from a filament Huffenberger, Rotti & Colins (2020)

Misalignment of filaments and magnetic fields creates TE>0, TB>0 and EB>0

