

# On the origin of primordial gravitational waves

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# (Spectator) axion-U(1) inflation

$$S = \int d^4x \sqrt{-g} \left[ \frac{M_{pl}^2 R}{2} - \underbrace{\frac{1}{2}(\partial\phi)^2 - V(\phi)}_{\text{Standard inflaton sector}} - \underbrace{\frac{1}{2}(\partial\chi)^2 - U(\chi)}_{\text{Axion}} - \underbrace{\frac{1}{4}F_{\mu\nu}F^{\mu\nu}}_{\text{Gauge field}} - \underbrace{\lambda \frac{\chi}{f} F_{\mu\nu} \tilde{F}^{\mu\nu}}_{\text{Chern-Simons coupling}} \right]$$

Barnaby + '12,  
Namba + '15,  
Mukohyama + '14

- ▶ Axion and gauge field are spectators
- ▶ Inflation realised by standard inflaton
- ▶ Time-dependent axion + Chern-Simons coupling → breaks conformal invariance of gauge field → amplification of only one helicity → **parity-violating GWs!**
- ▶ Amplitude of gauge field fluctuations controlled by axion's velocity:

$$A_- \propto e^{\pi\xi} = e^{\frac{\pi\lambda|\dot{\chi}|}{2Hf}}$$

# Vacuum vs sourced fluctuations

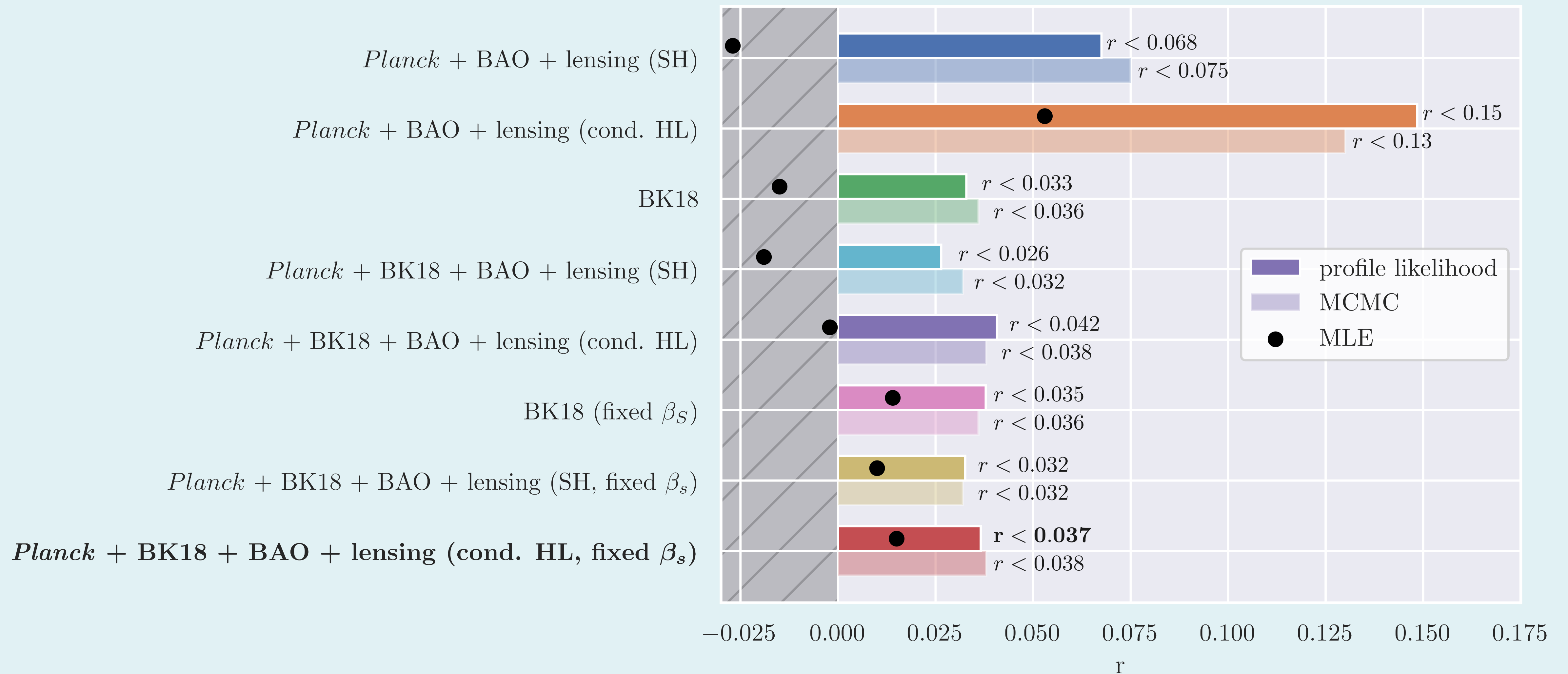
MODES	VACUUM	SOURCED	OBSERVATIONS
SCALARS	Slightly red-tilted, ~ Gaussian	Large and non-Gaussian or negligible depending on model	Remarkably consistent with vacuum
TENSORS	~ Scale-invariant, ~ Gaussian, non-chiral	Scale-dependent, non-Gaussian, chiral	? (To be checked)

# Constraint on vacuum fluctuations

PC & E. Komatsu 2022 [arXiv:2205.05617](https://arxiv.org/abs/2205.05617)

# Current CMB constraints on vacuum fluctuations

The CMB is the most promising way to detect vacuum fluctuations



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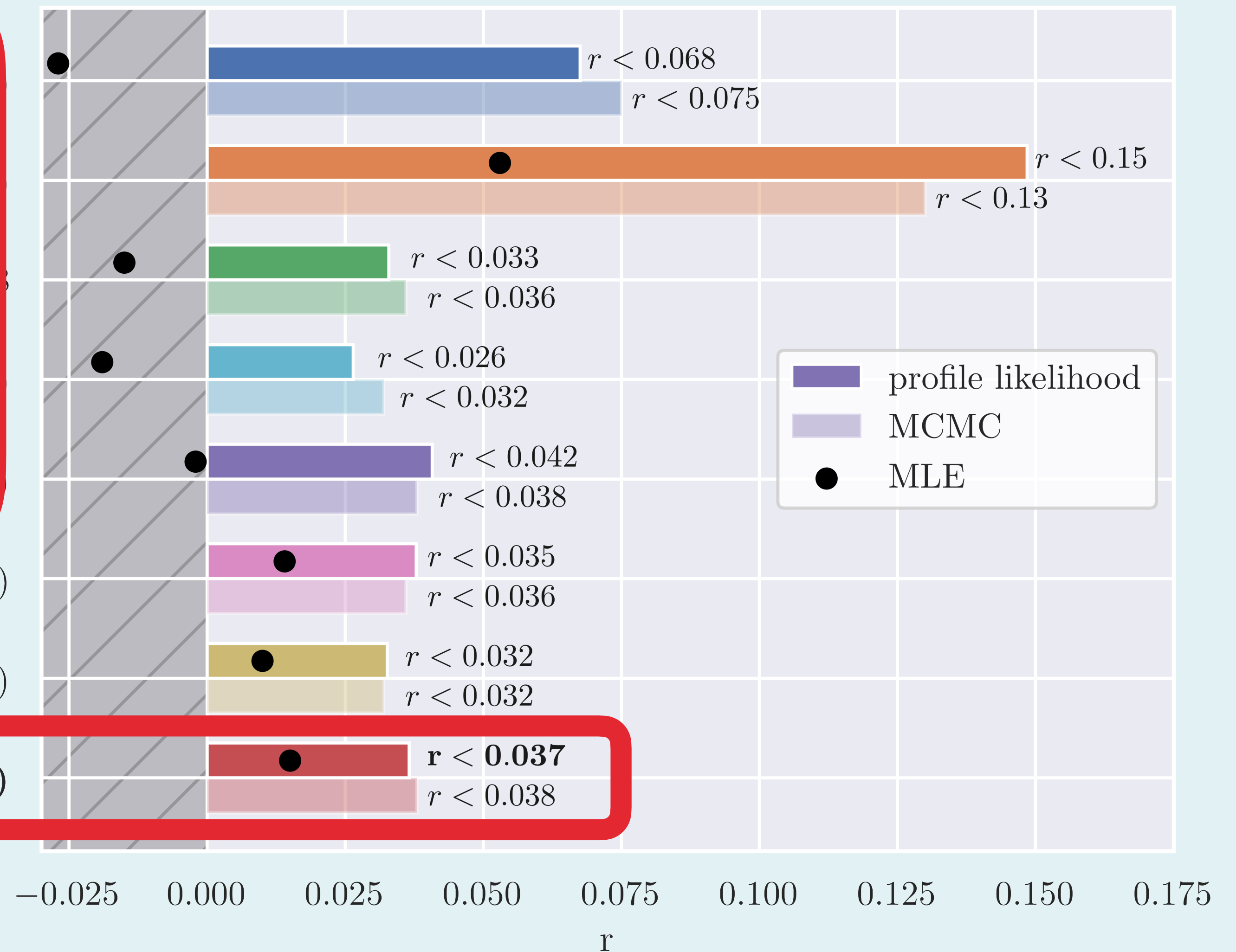
We discussed robustness against prior choices and volume effects using the profile likelihood

PC & E. Komatsu 2022 [arXiv:2205.05617](https://arxiv.org/abs/2205.05617)

BK18 (fixed  $\beta_S$ )

*Planck* + BK18 + BAO + lensing (SH, fixed  $\beta_S$ )

*Planck* + BK18 + BAO + lensing (cond. HL, fixed  $\beta_S$ )



# Constraint on gravitational waves sourced by matter fields

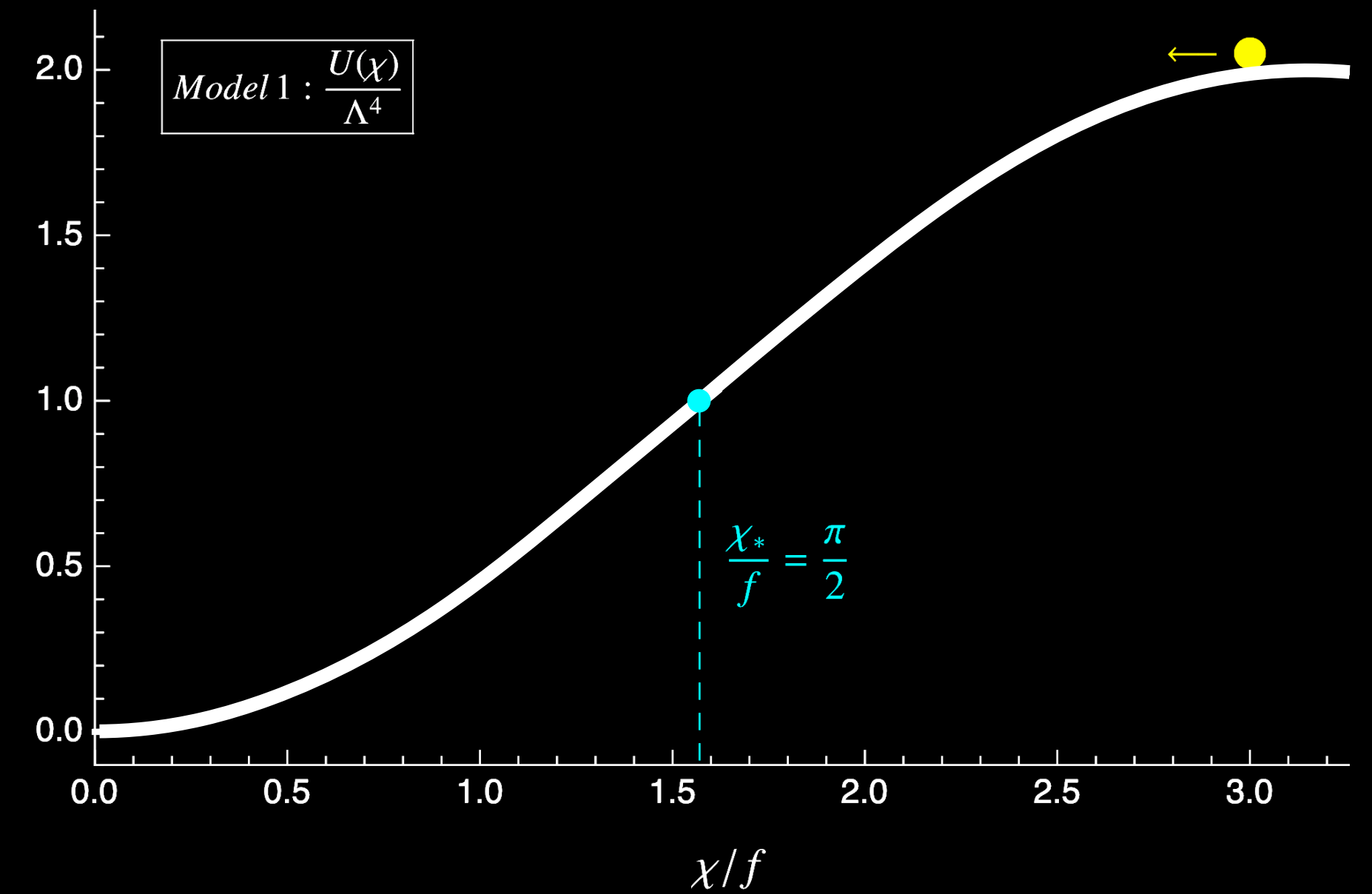
PC & O. Özsoy, I. Obata, M. Shiraishi 2022 [arXiv:2203.03401](https://arxiv.org/abs/2203.03401), accepted in JCAP

# Axion potential and sourced scalar modes

- ▶ Sourced **non-Gaussian scalar modes** from inverse decay of gauge fields  $A_i + A_i \rightarrow \delta\chi \rightarrow \delta\phi \propto \mathcal{R}$
- ▶ **Localized profile for axion velocity**  $\rightarrow$  amplify only large scales modes where CMB constraints are weaker ( $\ell \lesssim 100$ )
- ▶ 2 choices of axion potential (**M1, M2**):

$$\left\{ \begin{array}{l} U_{M1}(\chi) = \Lambda^4 \left[ 1 - \cos\left(\frac{\chi}{f}\right) \right] \quad \text{Freese + '90, Namba + '15} \\ U_{M2}(\chi) = \mu^3 \chi + \Lambda^4 \left[ 1 - \cos\left(\frac{\chi}{f}\right) \right] \quad \& \quad \Lambda^4 \lesssim \mu^3 f \quad \text{Özsoy + '20} \end{array} \right.$$

- ▶ Axion's velocity peaks at time  $\tau_*$
- ▶ Effective coupling strength  $\xi = -\frac{\lambda\dot{\chi}}{2Hf}$  peaks at  $\xi_* = \xi(\tau_*)$





# Sourced modes

- ▶ Strongly **scale-dependent** sourced tensor spectrum
- ▶ Gaussian bump feature:

$$\mathcal{P}_j^{(s)}(k) = \left[ \epsilon_\phi \mathcal{P}_{\mathcal{R}}^{(v)}(k) \right]^2 f_{2,j}^c[\xi_*, \delta] \exp \left[ -\frac{1}{2\sigma_{2,j}^2[\xi_*, \delta]} \ln^2 \left( \frac{k}{k_* x_{2,j}^c[\xi_*, \delta]} \right) \right]$$

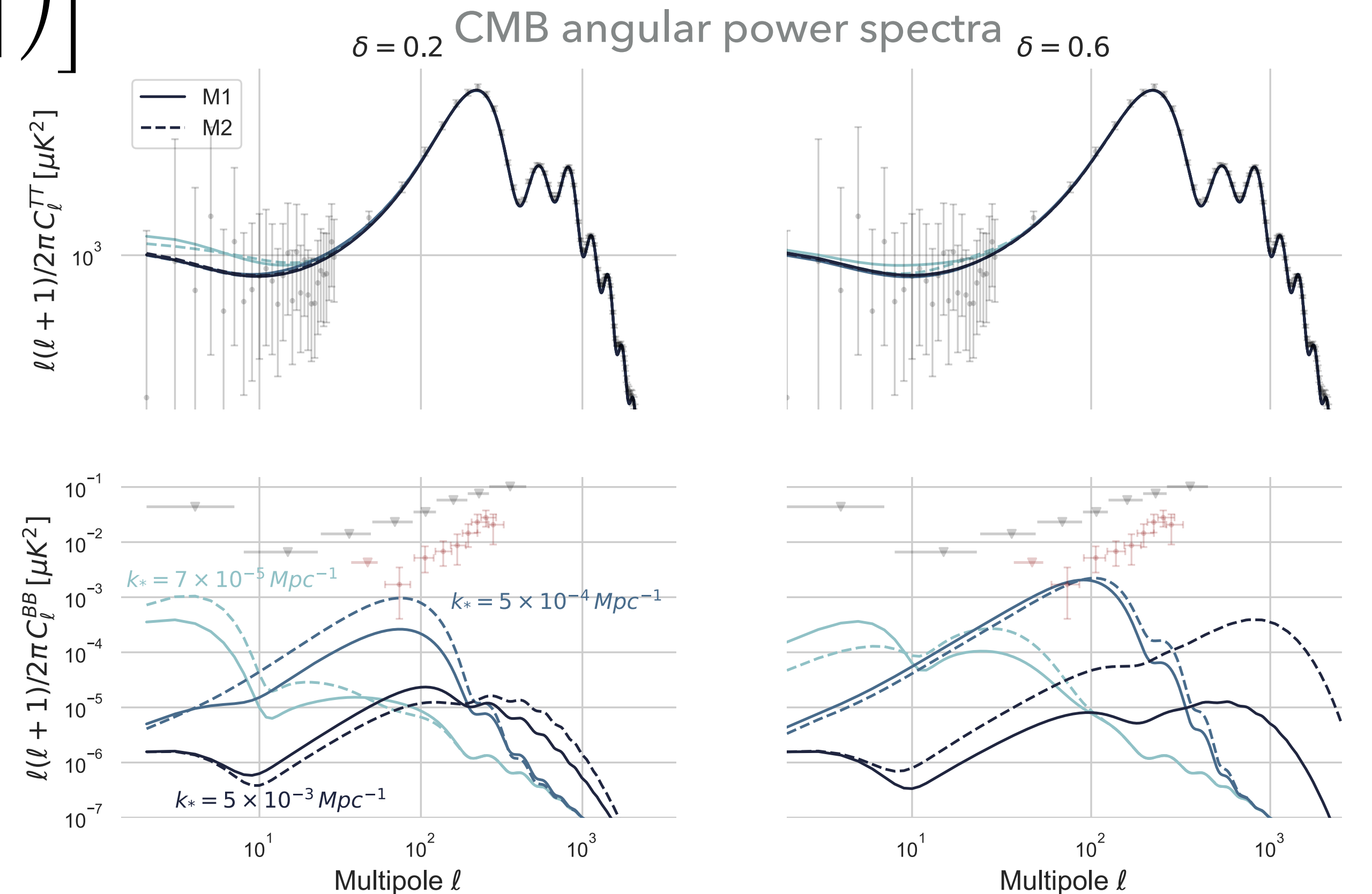
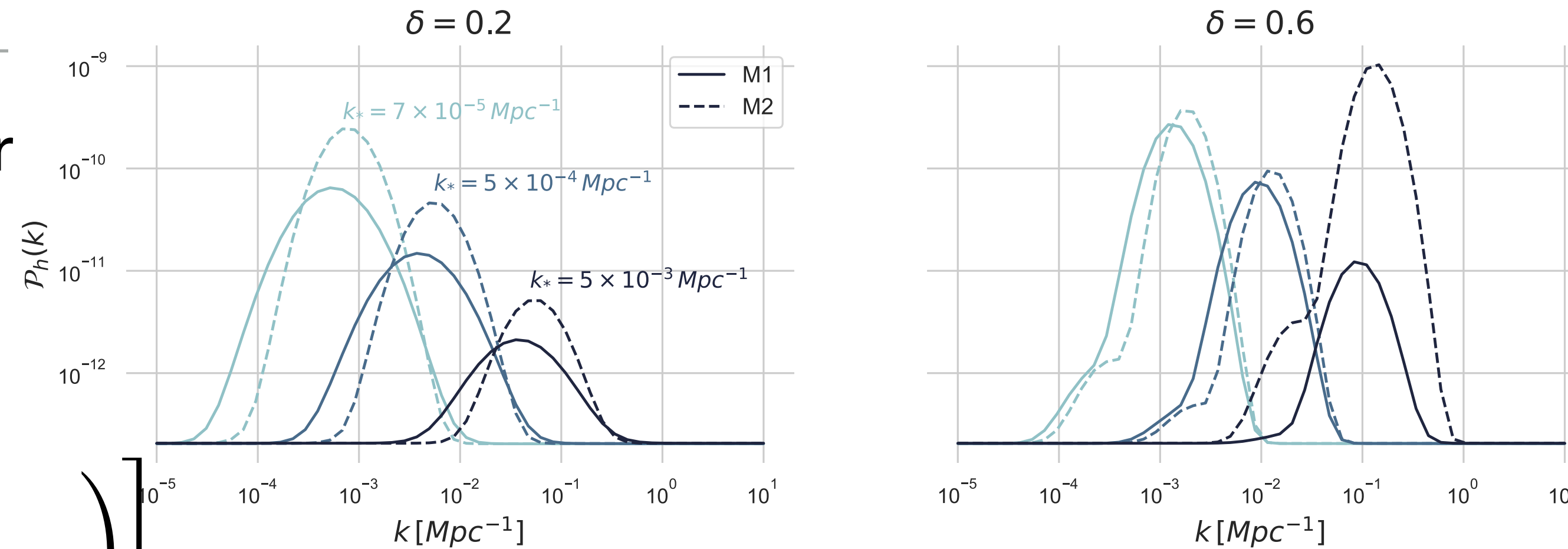
$$j = \{ \mathcal{R}, \pm \}$$

$\xi_*$  → amplitude of the peak

$\delta$  → ( $\propto$  axion's acceleration) width of the peak

$k_*$  → position of the peak (scale exiting the horizon when axion's velocity is max)

## Tensor power spectrum



## Method: data and likelihood

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### GOAL:

Determine the extent to which axion-U(1) can amplify tensor modes while staying consistent with CMB constraints

- ▶ We use **Planck + BICEP/Keck** latest data → state-of-the-art for constraining large cosmological scales
- ▶ Parity-violating correlations TB, EB give only very weak constraints (*Gerbino + '16*)
- ▶ Tensor bispectrum is complementary but expected to have lower SNR w.r.t. power spectrum (*Shiraishi + '19*)
- ▶ We exploit the **scale-dependence of the 2-point function** (most constraining at the current state)

# Method: profile likelihood

- ▶ Frequentist
- ▶ Profile likelihood for a parameter of interest  $\theta$ :
  1. fix  $\theta$  to multiple values
  2. Minimize  $\chi^2(\theta) = -2 \log \mathcal{L}(\theta)$
  3. By construction minimum  $\chi_{min}^2$  coincides with global MLE ("best-fit")

METHOD	MCMC	PROFILE LIKELIHOOD
PRIORS	Yes	No
INVARIANT UNDER REPARAMETRIZATION	Yes with Jeffreys prior	Yes
VOLUME EFFECTS	Maybe	No

# Method: profile likelihood

- ▶ Frequentist
- ▶ Profile likelihood for a parameter of interest
  1. fix  $\theta$  to multiple values
  2. Minimize  $\chi^2(\theta) = -\ln L(\theta)$
  3. By construction minimum coincides with global ("best-fit")

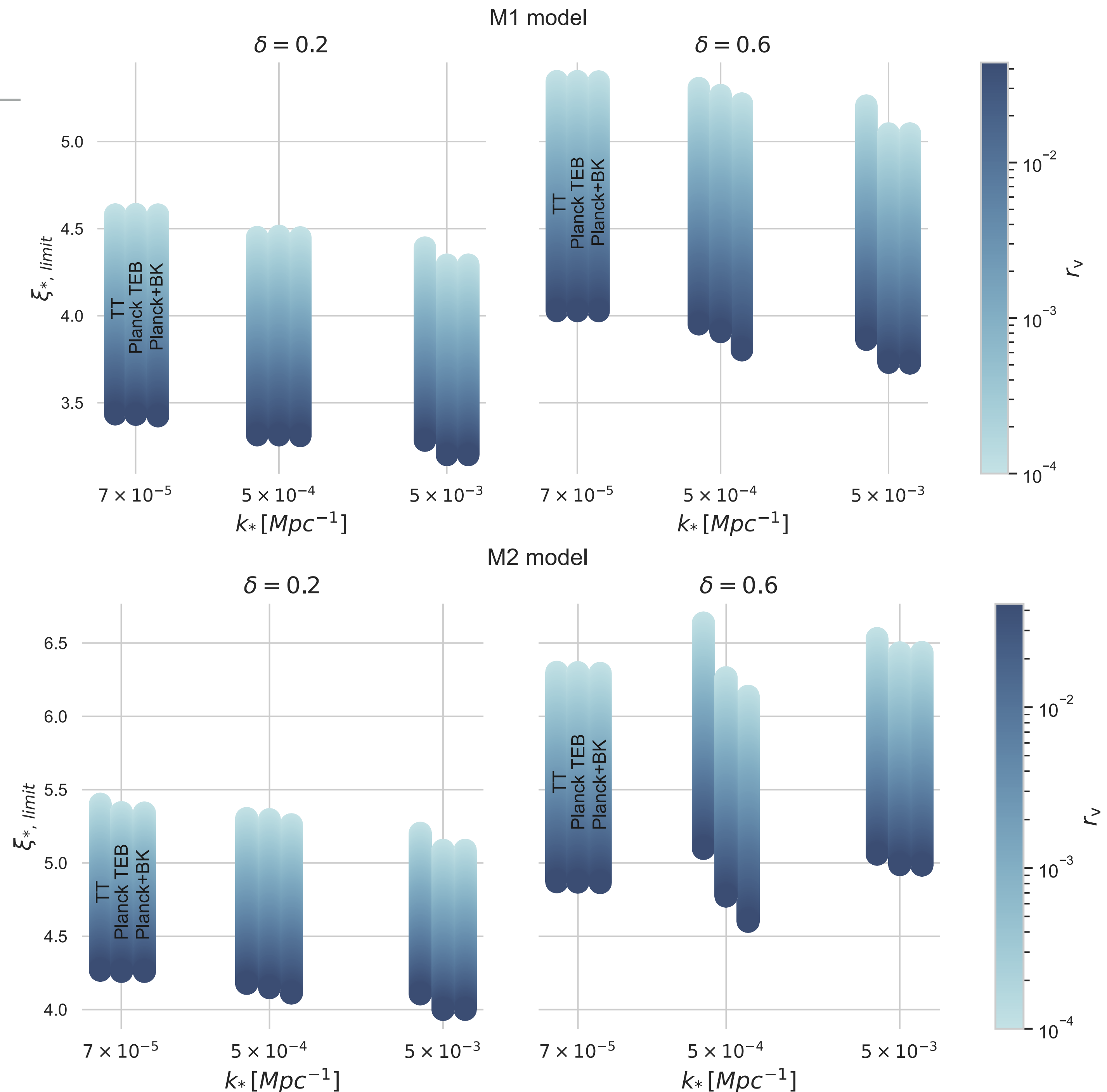
**We applied the profile likelihood to vacuum Planck +BICEP constraints and compared to MCMC**

**PC & E.Komatsu 2022**  
**arXiv:2205.05617**

MCMC	PROFILE LIKELIHOOD
	No
	Yes
	No

# Temperature vs polarization

- ▶ Observational constraints driven by *Planck* temperature (i.e. by sourced scalars!)
- ▶ *B*-modes data weakly constrained → minor effect on the bounds
- ▶ Large scale temperature data are already **cosmic-variance limited** in *Planck* data → improve polarization!

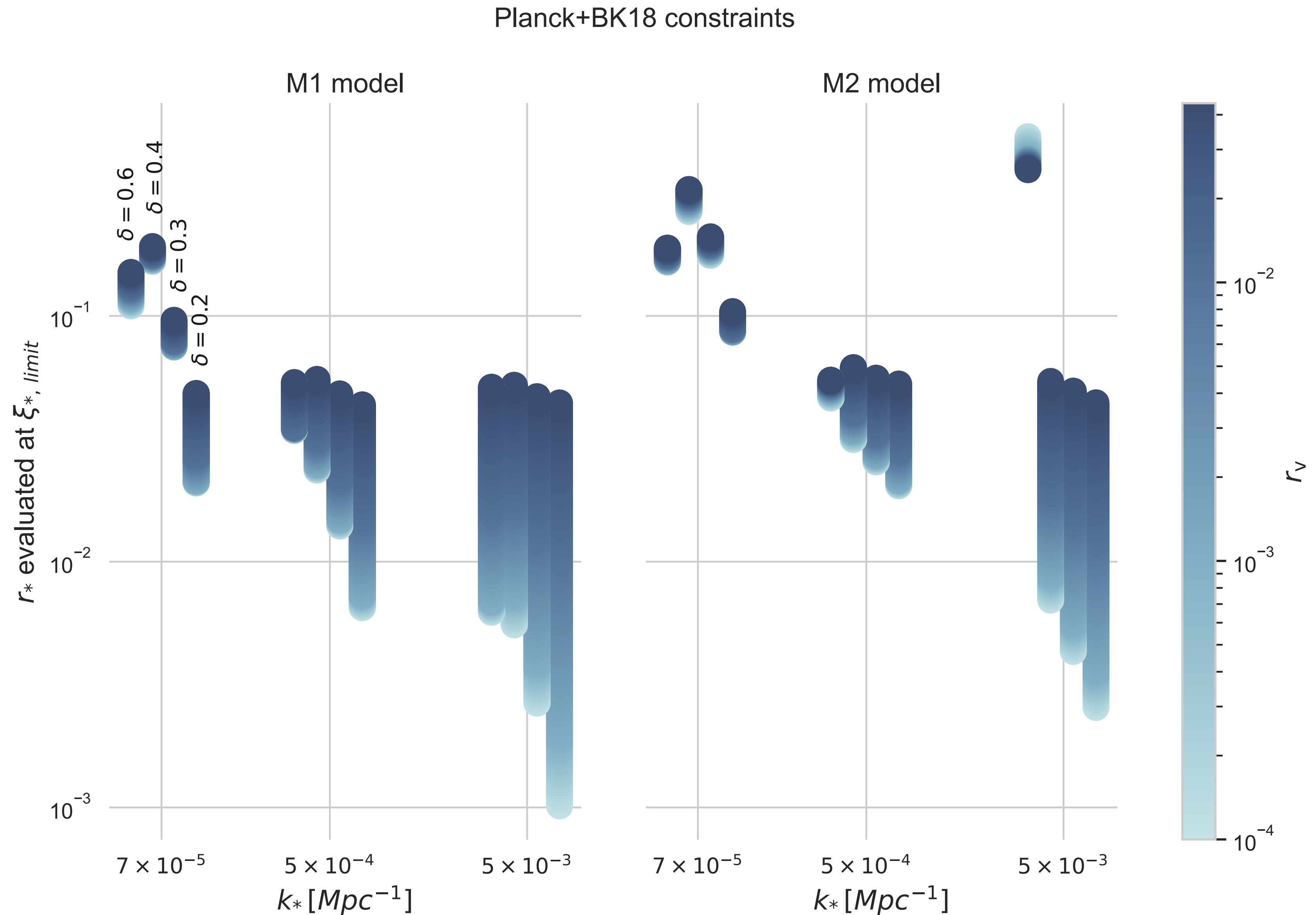


# Total tensor-to-scalar ratio

- ▶ Total (vacuum + sourced) tensor-to-scalar ratio:

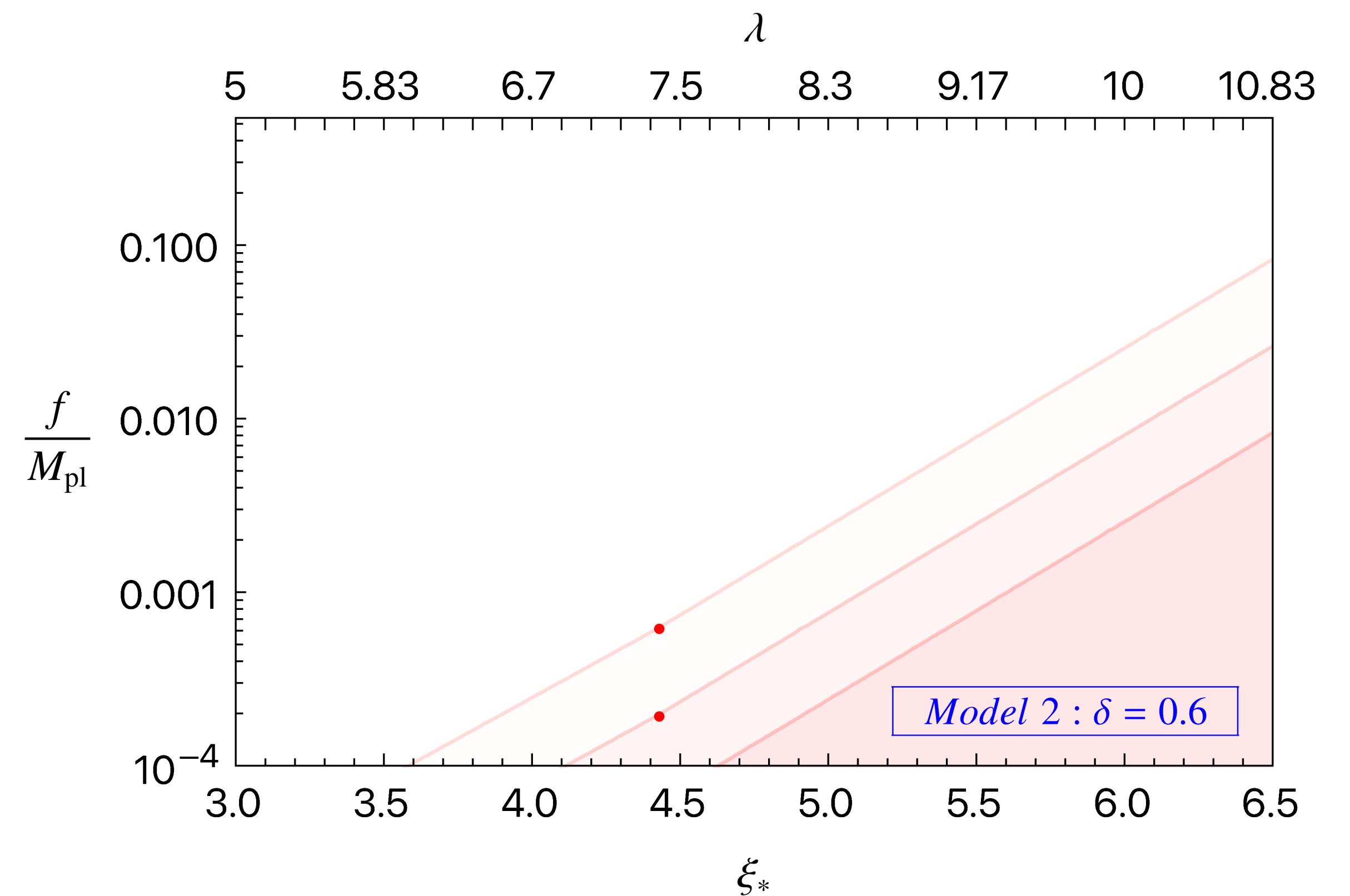
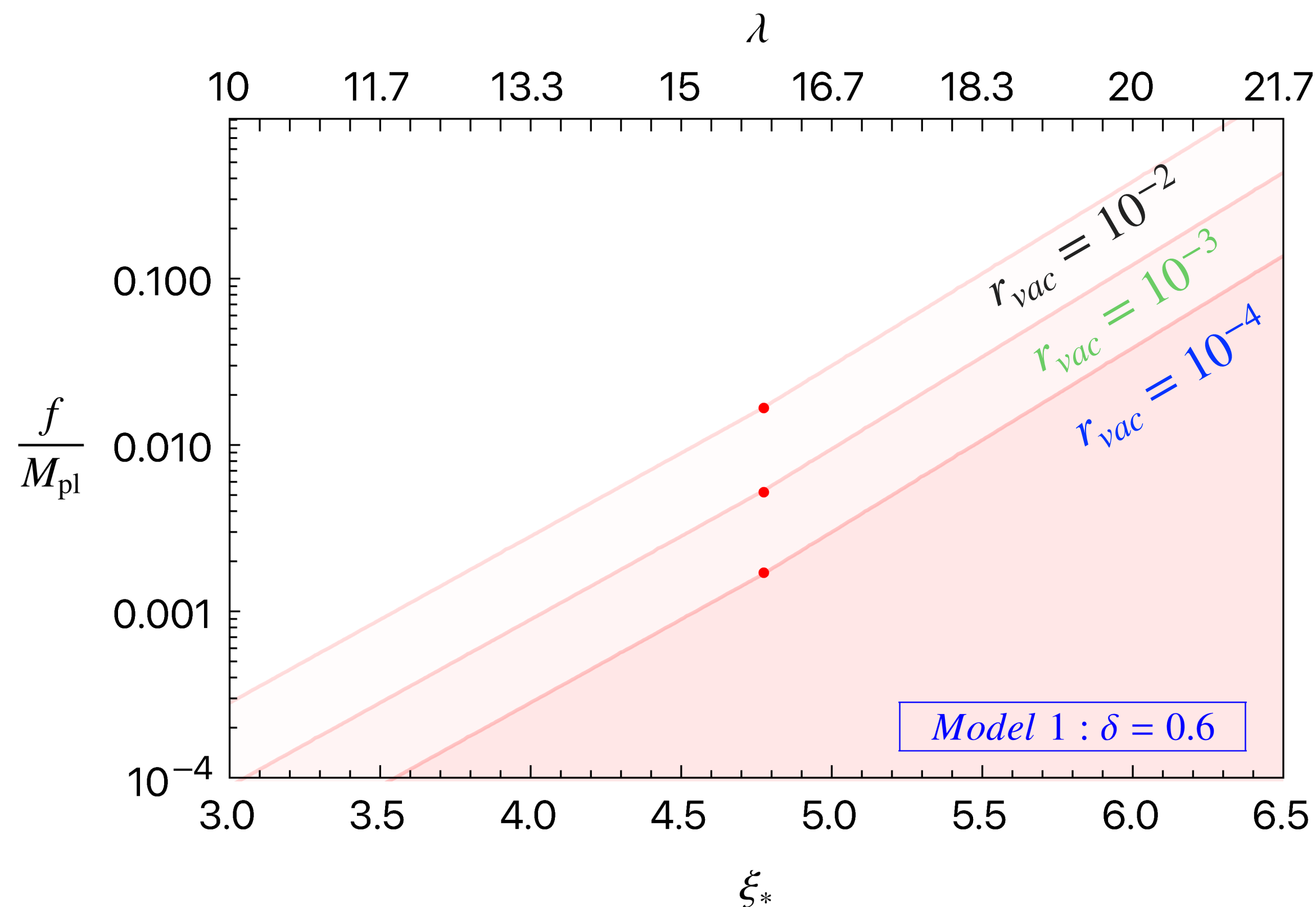
$$r_*(k) = \frac{\sum_{\lambda=\pm} \left[ \mathcal{P}_\lambda^{(\text{vac})}(k) + \mathcal{P}_\lambda^{(\text{sourced})}(k) \right]}{\mathcal{P}_\mathcal{R}^{(\text{vac})}(k) + \mathcal{P}_\mathcal{R}^{(\text{sourced})}(k)}$$

- ▶ Larger sourced signal allowed at larger scales
- ▶ Still possible to get significant sourced contribution to  $r_*$  even with small  $r_{vac}$



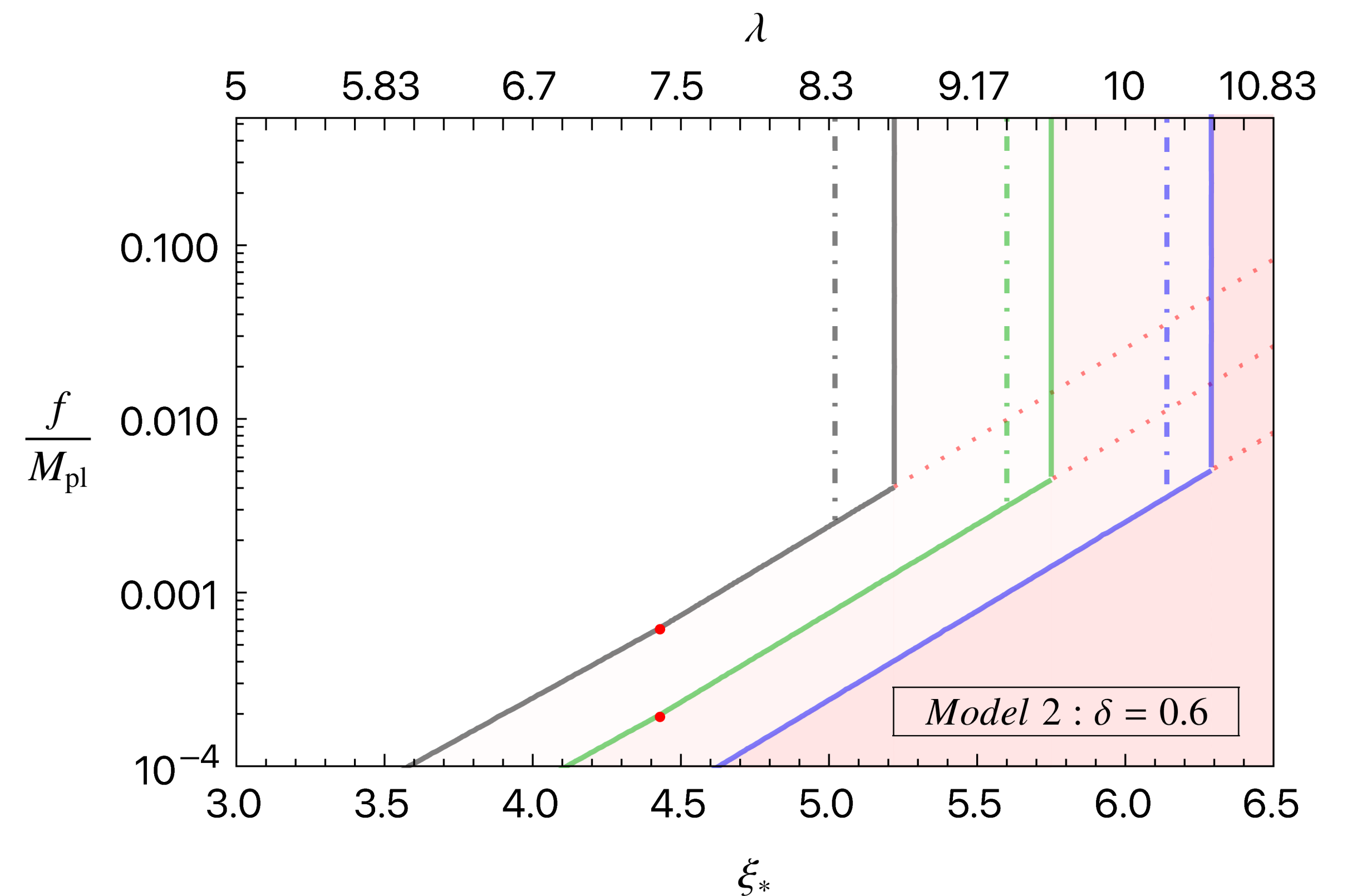
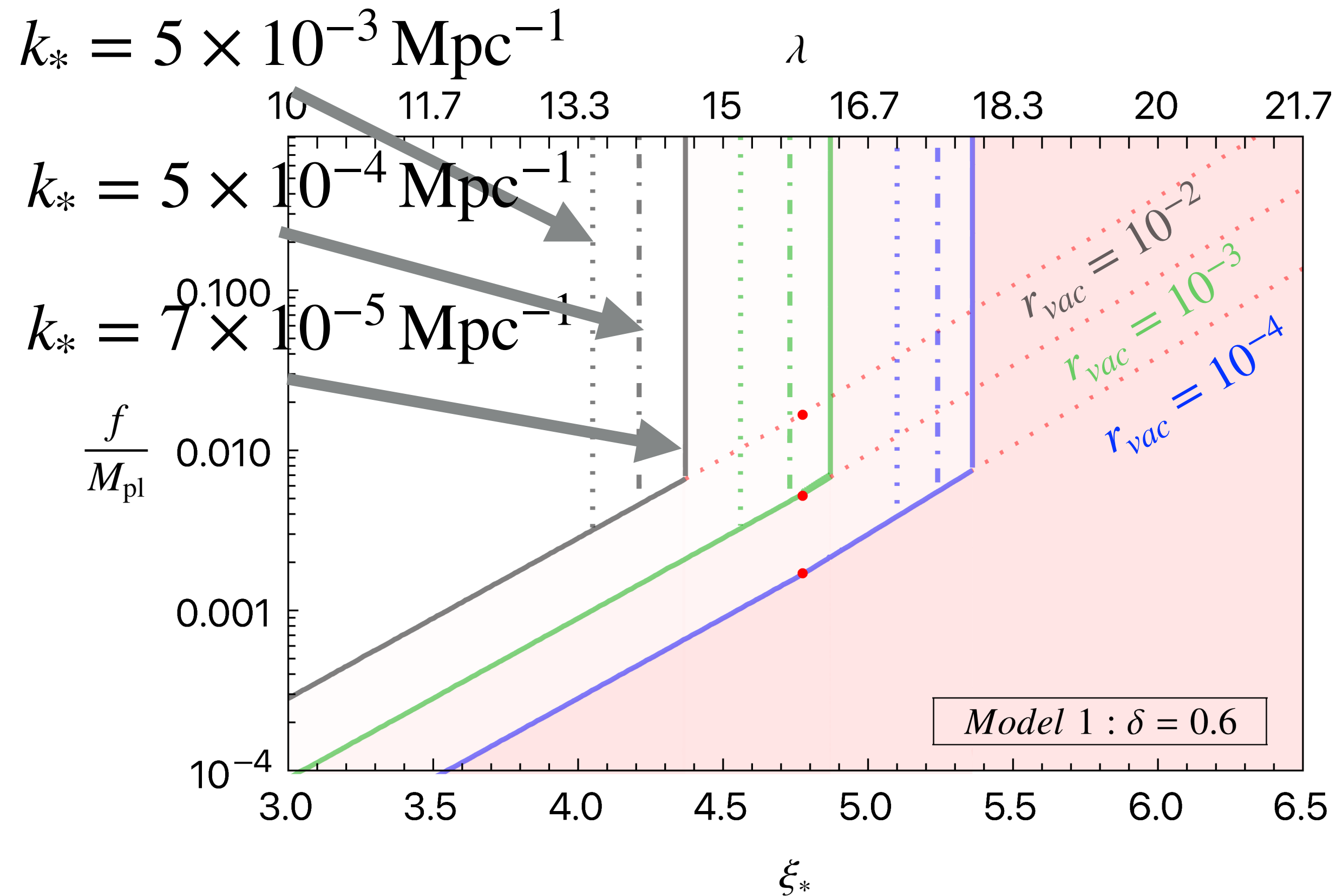
# Theoretical self-consistency bounds

- ▶ **Backreaction** upper/lower bounds on  $f/M_{pl}$  (axion must be a spectator / gauge field production should not influence background evolution of axion)
- ▶ **Perturbativity** (exponentially large gauge field amplitudes can drive the system out of perturbative regime) → lower bound on  $f/M_{pl}$



# Theoretical self-consistency + observations bounds

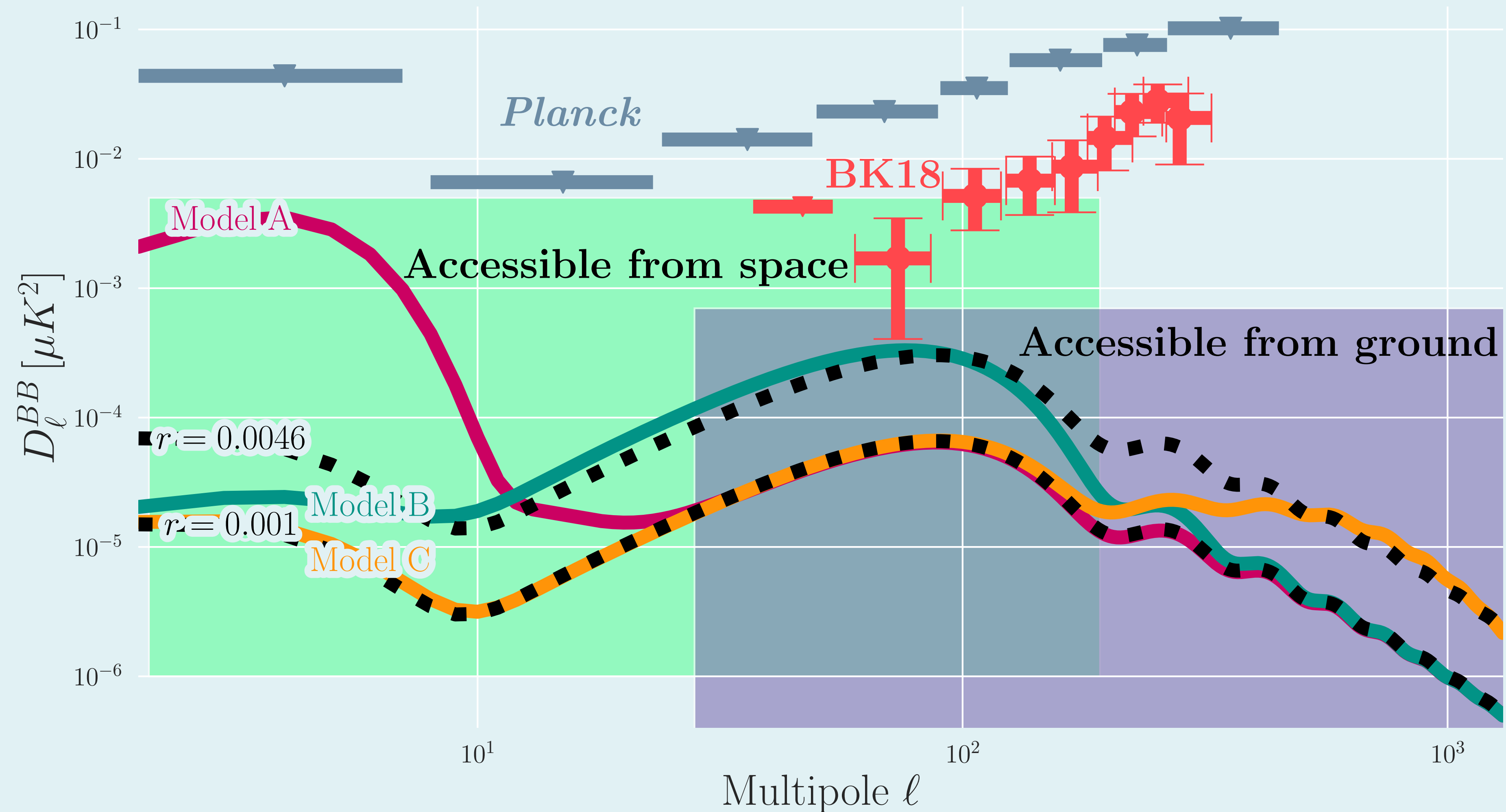
- ▶ Observational constraints are competitive with theoretical bounds
- ▶ Parameter space shrinks but still remains large and interesting for future B-mode experiments





# Conclusions: axion-U(1) is still interesting for future experiments

- ▶ Future **space** and **ground-based** B-mode experiments will be necessary to **distinguish vacuum from sourced** (Models A, B, C) primordial gravitational waves
- ▶ Full-sky space mission will improve also EB, TB measurements and detect  $\mathcal{O}(1)$  tensor non-Gaussianity



# (Spectator) axion-SU(2) inflation

$$\mathcal{L} = \underbrace{\mathcal{L}_{\text{inflaton}}}_{\text{Standard inflaton sector}} + \underbrace{\frac{1}{2} (\partial_\mu \chi)^2 - \mu^4 \left[ 1 + \cos\left(\frac{\chi}{f}\right) \right]}_{\text{Axion with cosine-type potential}} - \underbrace{\frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu}}_{\text{Gauge field}} + \underbrace{\frac{\lambda}{4f} \chi F_{\mu\nu}^a \tilde{F}^{a\mu\nu}}_{\text{Chern-Simons coupling}}$$

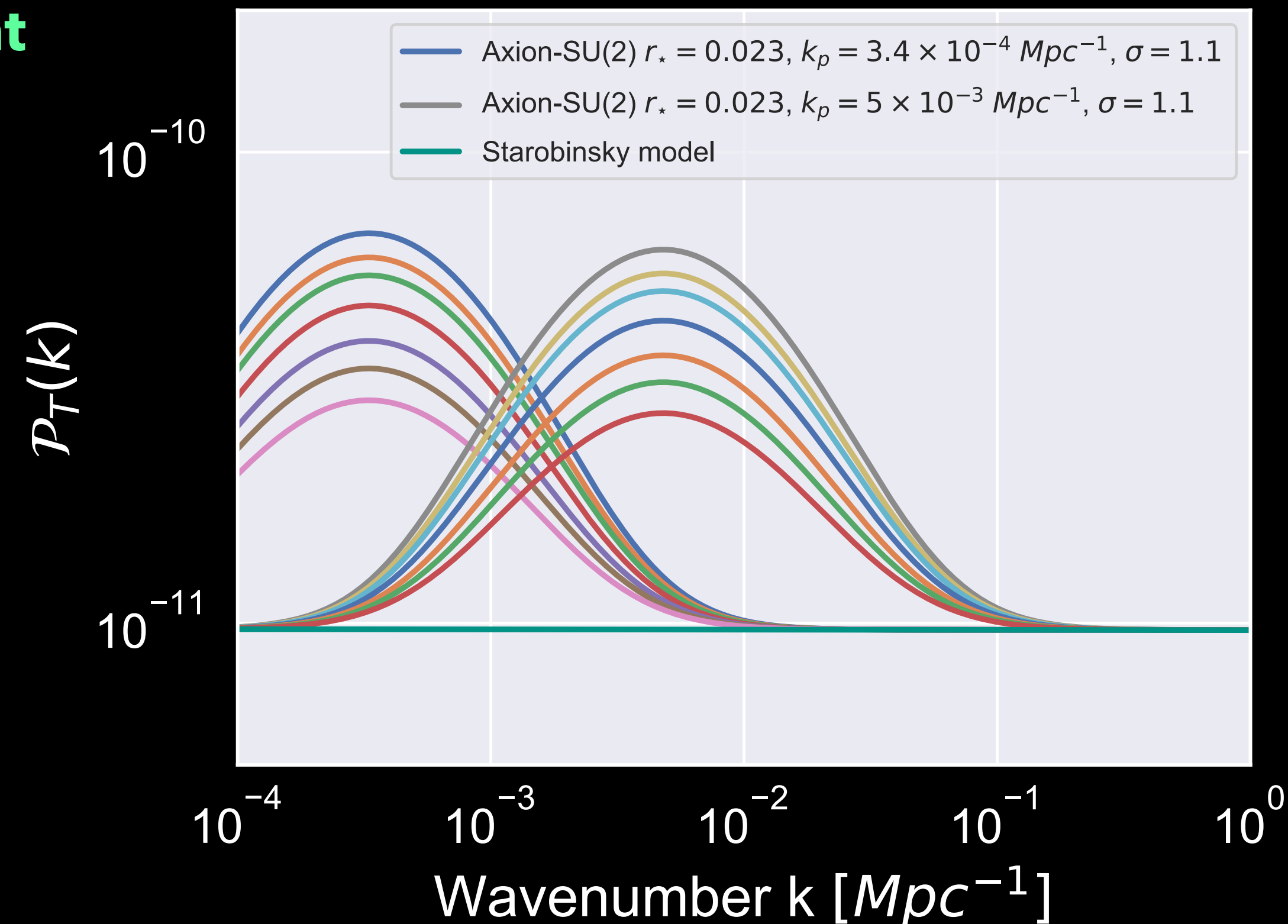
*Dimastrogiovanni + '17*

- ▶ **Chiral, non-Gaussian** sourced tensors with **scale-dependent** spectrum

$$P_T^{L, \text{Sourced}}(k) = r_* P_{\mathcal{R}}(k) \exp \left[ -\frac{1}{2\sigma^2} \ln^2 \left( \frac{k}{k_p} \right) \right]$$

*Thorne + '17*

- ▶ **U(1)** source tensors at  $2^{\text{nd}}$  order → source scalars and scalar non-Gaussianity
- ▶ **SU(2)** source tensors at linear order! → negligible sourced scalars and scalar non-Gaussianity!



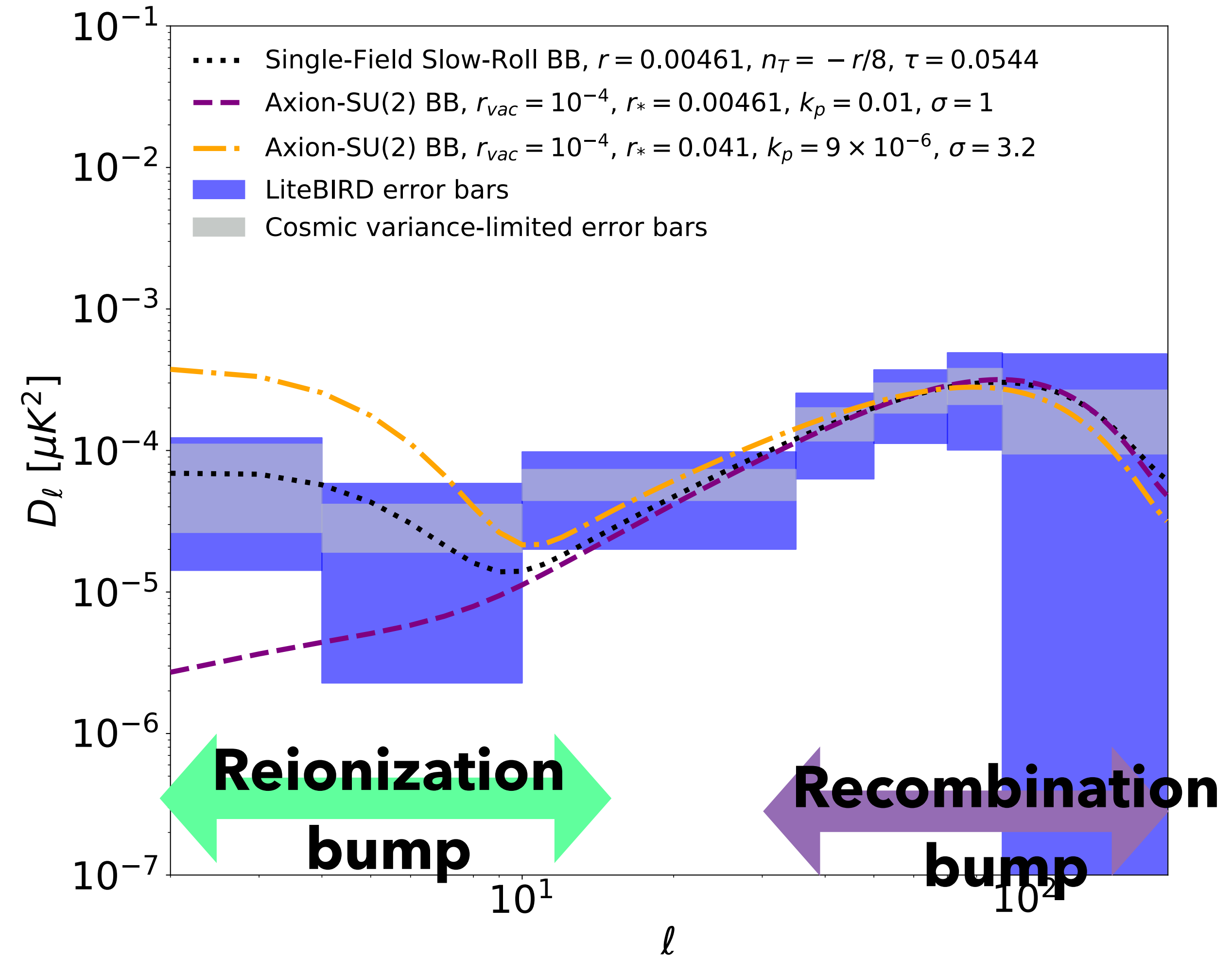
# LiteBIRD constraints on axion-SU(2) inflation

PC & E. Komatsu and the LiteBIRD Collaboration in preparation

# Constraining axion-SU(2) with LiteBIRD

PC, E. Komatsu + LiteBIRD collaboration (in preparation)

- ▶ SU(2) can source GWs that exceed vacuum contribution at reionization bump scales by factor  $\sim 5$  (Ishiwata + '21)
- ▶ **Goal:** show that full-sky survey with access to **reionization bump** is necessary to understand origin of primordial GWs
- ▶ **Method:** Realistic simulations, profile likelihood



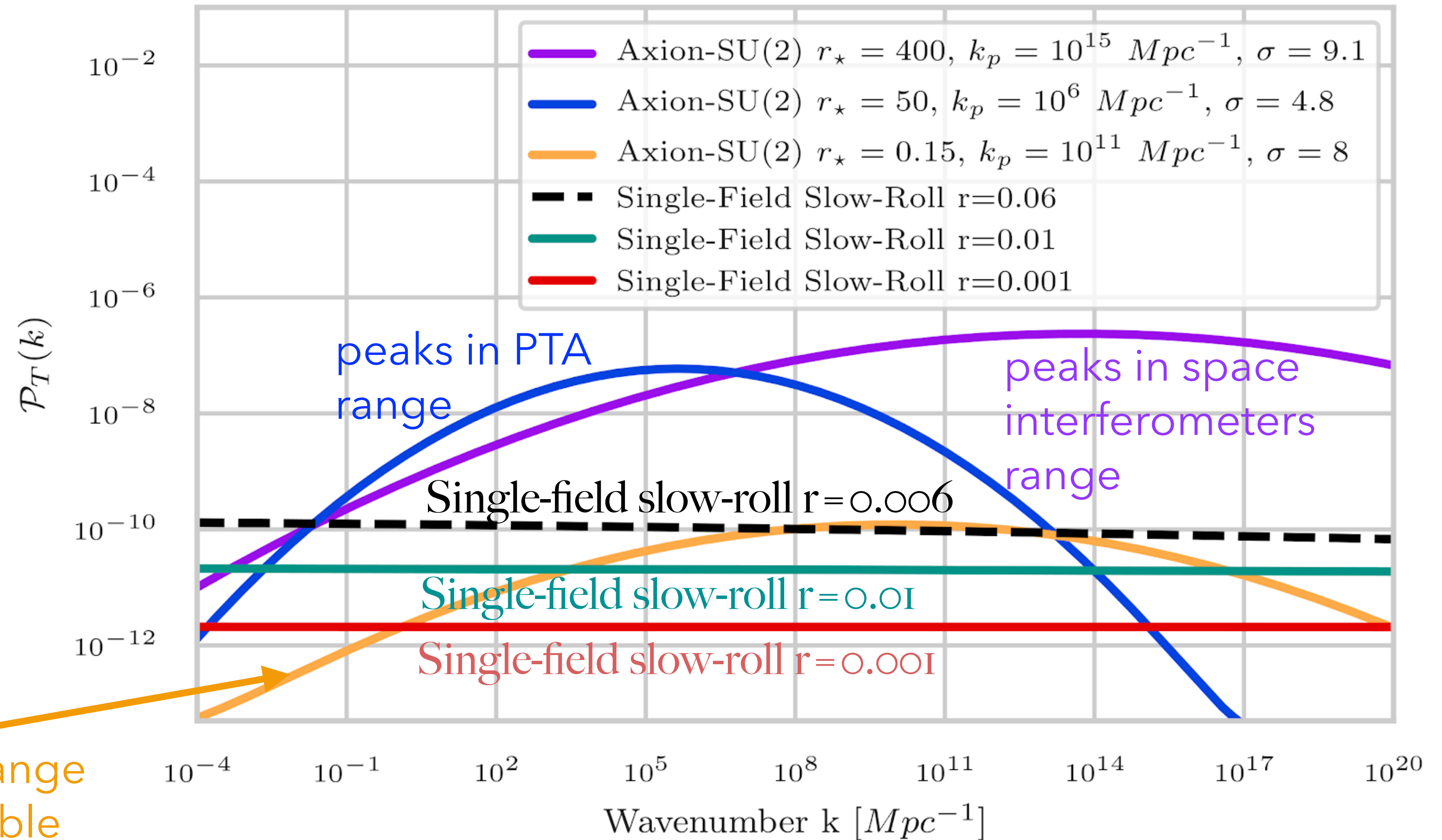
# Constraints on axion-SU(2) from CMB, PTA and interferometers

PC & E. Komatsu, D. Poletti, C. Baccigalupi 2021 [arXiv:2007.04241](https://arxiv.org/abs/2007.04241), JCAP 2021, 01,012

# Testing SU(2) with interferometers

- ▶ Gauge fields (e.g. SU(2)) can generate a signal detectable in many decades in frequency while staying consistent with CMB bounds
- ▶ Testable at PTA and space interferometers!

peaks in space interferometers range but is not detectable at CMB scales



# Astrophysical foregrounds for direct detection experiments

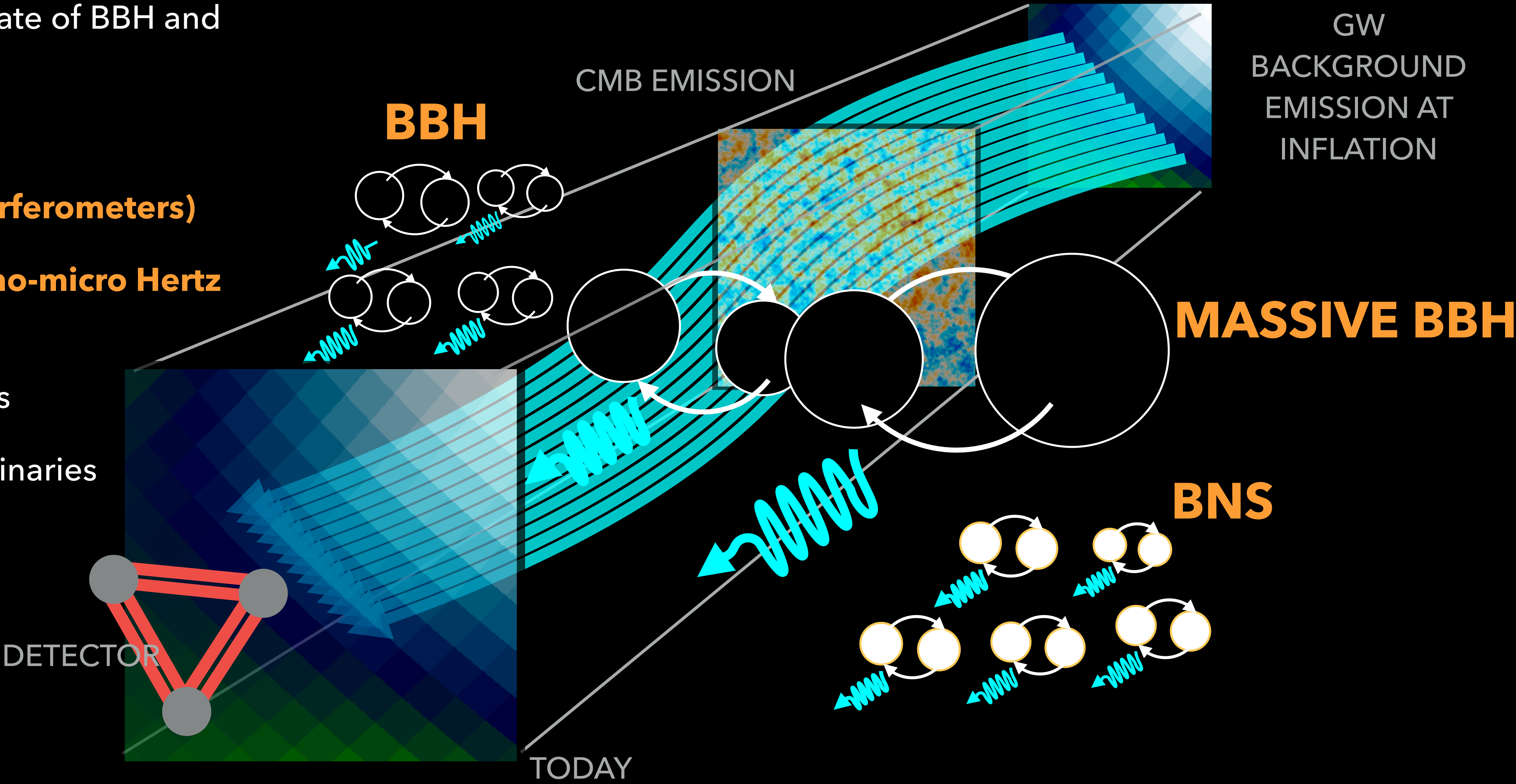
- ▶ Superposition of many astrophysical sources integrated over time
- ▶ LIGO/Virgo measured rate of BBH and BNS mergers
- ▶ Main sources:

- ▶ **BBH + BNS (all interferometers)**

- ▶ **Massive BBH in nano-micro Hertz range**

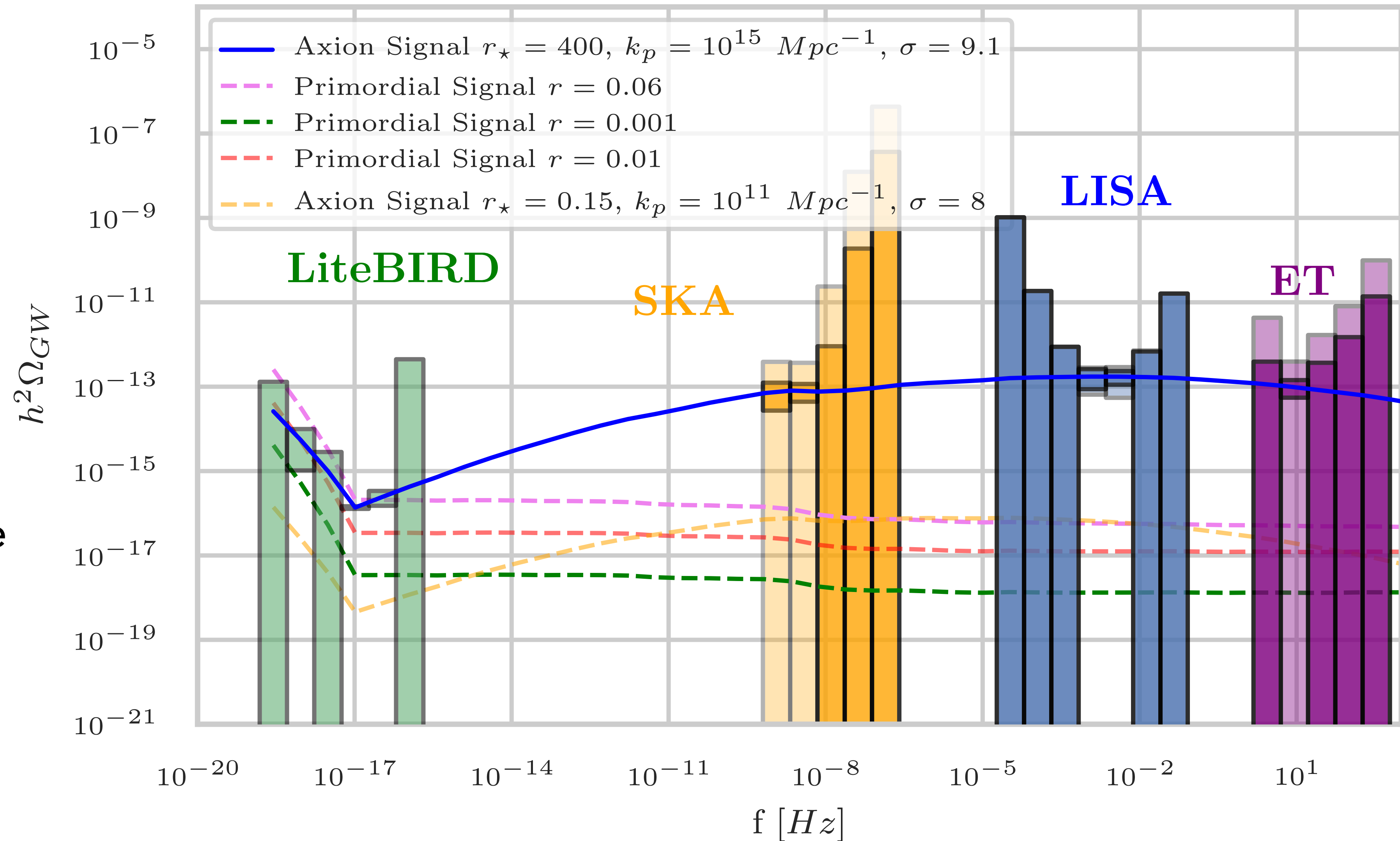
- ▶ Galactic WD binaries

- ▶ Extra-Galactic WD binaries



# axion-SU(2) at interferometers scales

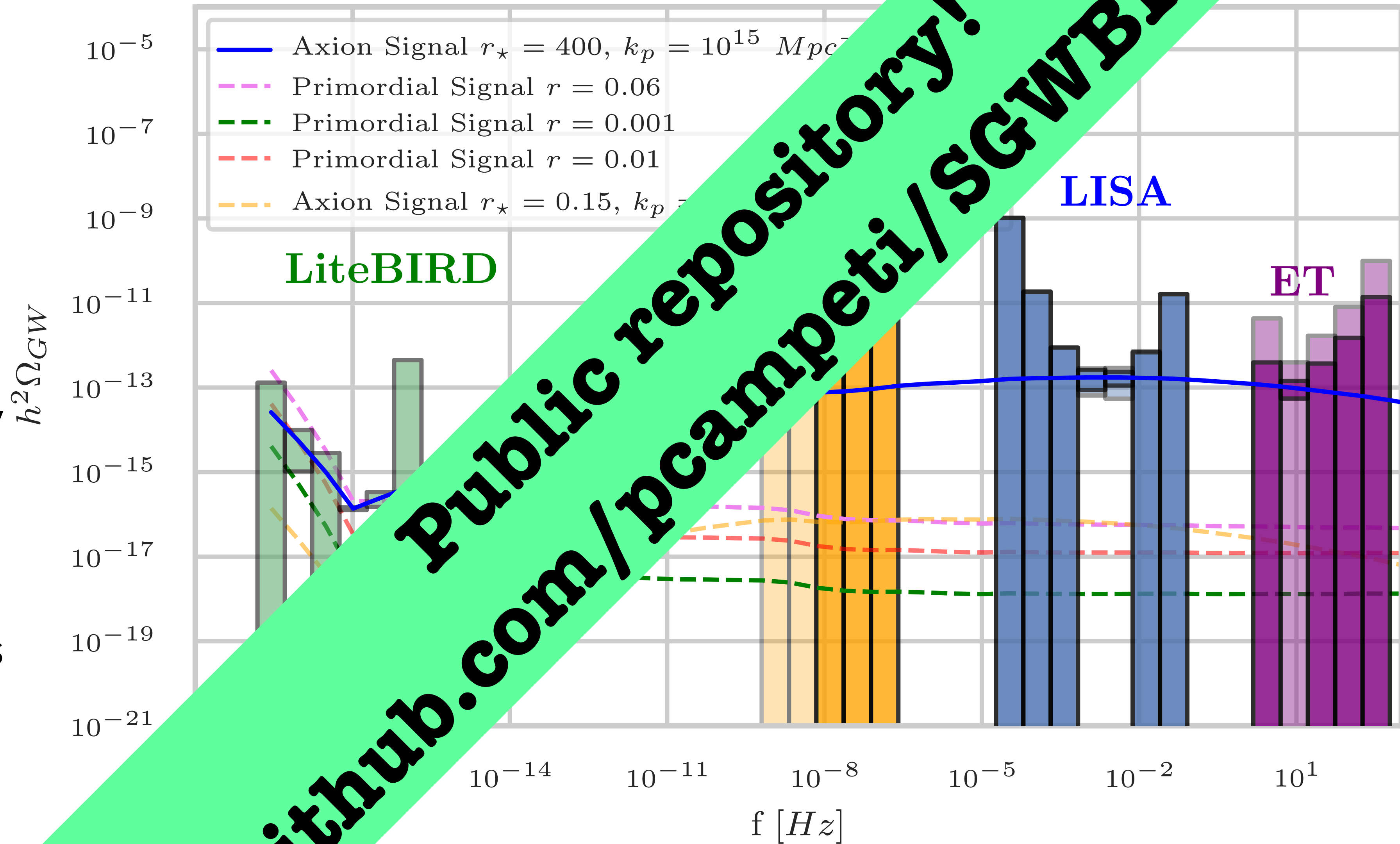
- ▶ We derived **new filter** for cross-correlation and foreground marginalisation for interferometers (including multi-band foreground cleaning)
- ▶ We take into account **foregrounds** for every experiment (lighter shade error bars).
- ▶ Coherent assumptions and realism for every experiment





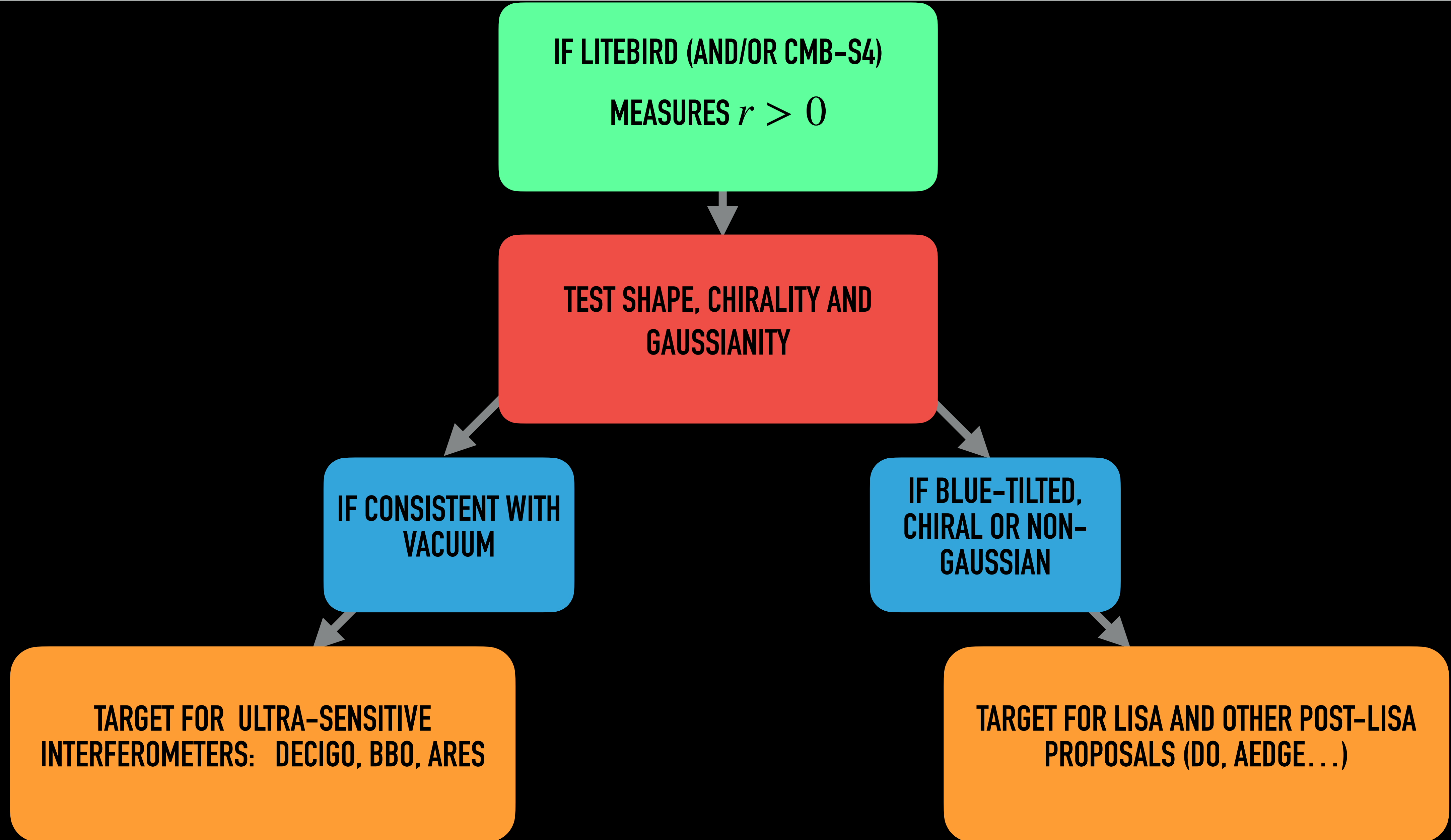
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# Future roadmap

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# Conclusions

1.



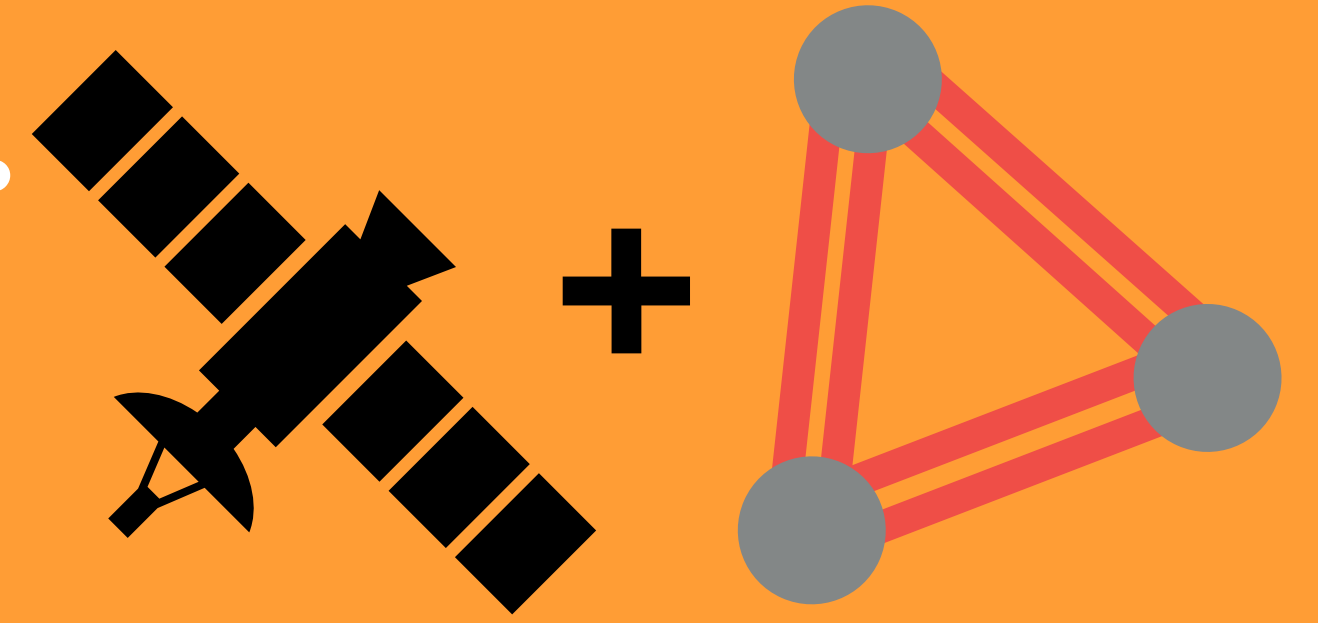
In light of **Planck** and **BICEP/Keck** data the parameter space of axion-U(1) models remains large and interesting for future B-mode experiments.

2.



Axion-SU(2) can produce sourced GWs exceeding vacuum contribution especially at reionization bump scales, which are accessible only with a full-sky space mission like **LiteBIRD**.

3.



- ▶ Measuring the **shape** of the GW spectrum along many decades in frequency is needed to understand its **origin**.
- ▶ Control of **foregrounds** is fundamental for all probes.
- ▶ **B-mode** experiments are the most sensitive and closest in time.
- ▶ Results suggest a future roadmap.