# Isotropic cosmic birefringence from early dark energy 

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Mini-workshop at MPA
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- Rotation of the CMB polarization (= CMB Birefringence)
T. Fujita, Y. Minami, KM, H. Nakatsuka

Phys. Rev. D 105 103519, arXiv: 2110.03228
T. Fujita, KM, H. Nakatsuka, S. Tsujikawa

Phys. Rev. D 105 103518, arXiv: 2203.03977

## This talk

"Isotropic cosmic birefringence from early dark energy"
E. Komatsu, KM, T. Namikawa

In preparation

- Big Bang Nucleosynthesis
- Primordial black holes
- Axion-gauge dynamics in inflation
I. Cosmic birefringence
II. Early dark energy
III. CB from EDE
IV. Summary


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## Cosmic birefringence

- Cosmic birefringence

Cosmic birefringence is a rotation of the plane on linear polarization in CMB.

CMB polarization is decomposed into Parity-even $E$ mode
Parity-odd $B$ mode


Cosmic birefringence mixes $E$ and $B$ modes:


## Cosmic birefringence

- Parity violating signal in CMB

New analysis of the Planck data reported "cosmic birefringence".
[Minami \& Komatsu(2020)]
The measured cosmic birefringence is considered to be

## (isotropic

(independent of the photon frequency
The isotropic rotation angle is estimated as

$$
\beta=0.342_{-0.091^{\circ}}^{+0.094^{\circ}} \text { at } 68 \% \text { C.L. [Eskilt \& Komatsu (2022)] }
$$

This signal indicates the existence of new physics!

## Cosmic birefringence

■ Cosmic birefringence from axion
Axion is a candidate of the origin of $\beta$.

$$
\mathscr{L}=-\frac{1}{2}\left(\partial_{\mu} \phi\right)^{2}-V(\phi)-\frac{1}{4} F_{\mu \nu} F^{\mu \nu}-\frac{1}{4} g \phi F_{\mu \nu} \tilde{F}^{\mu \nu}
$$

Dispersion relations for left/right circular polarization photons:

$$
\omega_{ \pm} \simeq k \mp \frac{g}{2}\left(\frac{\partial \phi}{\partial t}+\frac{\vec{k}}{k} \cdot \vec{\nabla} \phi\right)=k \mp \frac{g}{2} \frac{\mathrm{~d} \phi}{\mathrm{~d} t}
$$

$\rightarrow \beta$ is represented by the difference of the axion field value:

$$
\beta(t)=-\frac{1}{2} \int_{t}^{t_{0}} \mathrm{~d} t\left(\omega_{+}-\omega_{-}\right)=\frac{g}{2}\left[\phi\left(t_{0}\right)-\phi(t)\right]
$$

## Cosmic birefringence

## ■ Tomographic approach

CMB polarization mainly comes from

- Recombination ( $z \sim 1090$ )
- Reionization $(z \sim 7)$

Recombination



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## Early dark energy

- Hubble tension

Hubble tension is a discrepancy of the value of $H_{0}$ between the local measurement and early-universe predictions.


Hubble Space Telescope $74.03 \pm 1.42 \mathrm{~km} / \mathrm{s} / \mathrm{Mpc}$

Planck 2018
$67.4 \pm 0.5 \mathrm{~km} / \mathrm{s} / \mathrm{Mpc}$

## Early dark energy

■ Idea of Early dark energy [Karwal \& Kamionkowski (2016)]


Increase $H(z)$ at $z \gtrsim z_{\star}$
$\rightarrow$ Decrease $r_{\mathrm{s}}^{\star}$ and then $D_{A}^{\star}$
$\rightarrow$ The inferred value of $H_{0}$ increases.

## Early dark energy

■ Early dark energy model
EDE must satisfy the followings:

- Relevant at matter-radiation equality
- Behaves like dark energy at early times
- Dilutes faster than the matter after $z_{\star}$

An EDE model includes a field with a potential:

$$
V_{\mathrm{cos}}^{(n)} \equiv m^{2} f^{2}\left[1-\cos \left(\frac{\phi}{f}\right)\right]^{n}, \quad n \geq 2
$$

[Poulin, Smith, Karwal, Kamionkowski (2019)]
Around $\phi=0, V(\phi)$ is approximated $\propto \phi^{2 n}$.

$$
\rho_{\phi} \propto a^{-6 n /(n+1)}
$$

[Turner (1983)]

## Early dark energy

## ■ Best-fit for EDE

We consider a model with $V=m^{2} f^{2}\left[1-\cos \left(\frac{\phi}{f}\right)\right]^{3}$.

Using a profile likelihood, the favored value of $f_{\text {EDE }}$ is estimated:

$$
\begin{array}{r}
f_{\mathrm{EDE}}=0.072_{-0.060}^{+0.071} \\
\\
\text { at } 95 \% \text { C.L. }
\end{array}
$$

[Herold, Ferreira, Komatsu (2022)]


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## Isotropic CB from EDE

## ■ Implementation

In the previous work, the effect of $\beta$ has been implemented
with an input data of $\phi(\eta)$. [Nakatsuka, Namikawa, Komatsu (2022)]

$$
\begin{aligned}
& \frac{{ }_{ \pm 2} \Delta_{P, l}\left(\eta_{0}, q\right)}{}=-\frac{3}{4} \sqrt{\frac{(l+2)!}{(l-2)!}} \int_{0}^{\eta_{0}} \mathrm{~d} \eta \tau^{\prime} e^{-\tau(\eta)} \Pi(\eta, q) \frac{j_{l}(x)}{x^{2}} e^{ \pm 2 i \beta(\eta)} \\
& \text { Fourier of } Q, U \\
& \Delta_{E, l}(q) \pm i \Delta_{B, l}(q) \equiv-{ }_{ \pm 2} \Delta_{P, l}\left(\eta_{0}, q\right) \quad \beta(\eta)=\frac{g}{2}\left[\phi\left(\eta_{0}\right)\right.
\end{aligned}
$$

To deal with EDE, $\phi(\eta)$ should be solved consistently with the background cosmology.

We extend the code to solve following CLASS_EDE code.

## Isotropic CB from EDE

## ■ EB angular power spectrum

We consider the EDE model with the best-fit parameters for
$f_{\mathrm{EDE}}=0.01,0.07$, and 0.14 .
Here, we use $g=M_{\mathrm{Pl}}^{-1}$ and $C_{l}^{E B}$ scales as $\propto g$.



## Isotropic CB from EDE

## - EB angular power spectrum

$C_{l}^{E B}$ is not proportional to $C_{l}^{E E}$.
The peak is shifted by $\Delta l \gtrsim 10$.


## Isotropic CB from EDE

## ■ EB angular power spectrum

For $f_{\mathrm{EDE}}=0.14$, the sign of $C_{l}^{E B}$ flips in mid $l$.


## Isotropic CB from EDE

## - EB angular power spectrum



## Isotropic CB from EDE

- EB angular power spectrum

Rough translation into $\beta: \frac{\operatorname{Max}\left[C_{l}^{E B}\right]}{\operatorname{Max}\left[C_{l}^{E E}\right]}=\frac{1}{2} \sin \left(4 \beta_{\mathrm{eff}}\right)$

We obtain $\beta_{\text {eff }}=\beta_{\text {obs }}$ with $g M_{\mathrm{Pl}}=\{1.2,0.42,0.47\}$.
$\rightarrow g=\mathcal{O}\left(M_{\mathrm{Pl}}^{-1}\right)$ is favored.

$f_{\text {EDE }}$
$-0.01$
$-0.07$

- 0.14


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- I extended the modified CLASS codes and calculated the EB power spectrum when an EDE field induces cosmic birefringence.
- For all $f_{\text {EDE }}$ I considered,
the EB spectrum is not proportional to the EE spectrum.
- Especially, $f_{\mathrm{EDE}}=0.14$ has a distinct shape.
- Future direction:

Anisotropic CB
Including lensing


