

Max-Planck-Institut für Astrophysik



## Analytical model for non-thermal pressure in galaxy clusters & its application to mass estimation

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ICM physics and modeling, June 16th, 2015

# Non-thermal pressure in galaxy clusters

### Hydro simulations find kinetic pressure in the ICM

- increasing fraction with radius
- of order 20% at r<sub>500</sub>



Evrard90, Rasia+04, I2, Dolag+05, Nagai+07, Lau+09, Battaglia+I2...

# **P**<sub>rand</sub>: the major known physical contributor to the HSE mass bias

- B field contribution unclear but not dominating
- CR upper limits already tight



#### P<sub>rand</sub> hard to observe (esp. at large radii where cluster masses are estimated)

see Zhuraveleva and Vacca's talks, though

# Analytical model for Prand

## Is this possible ...?



## Pressure ~ energy density

turbulent ICM

# Our model

## Injection & dissipation of random kinetic energy at Eulerian positions

$$\frac{d\sigma_{nth}^2}{dt} = -\frac{\sigma_{nth}^2}{t_d} + \eta \frac{d\sigma_{tot}^2}{dt} + [Diffusion] + [Advection]$$
  
$$\sigma_{nth}^2 = P_{rand} / \rho_{gas} \propto E_{rand} \text{ per unit mass}$$

dissipation injection



### But WHERE and HOW ?

Original source of energy: gravitational energy of infalling material



A previous idea: Cavaliere +11:  $E_{kin} \rightarrow E_{th} + E_{rand}$  at accretion shock

## Injection

### WHERE and HOW ?

### Our idea: trace the bulk of energy flow

#### Low Mach number internal (merger) shocks process more kinetic energy



$$\frac{\mathrm{d}\sigma_{\mathrm{nth}}^2}{\mathrm{d}t} = -\frac{\sigma_{\mathrm{nth}}^2}{t_{\mathrm{d}}} + \eta \, \frac{\mathrm{d}\sigma_{\mathrm{tot}}^2}{\mathrm{d}t}$$

### injection

same source responsible for the heating of ICM, and synchronized with growth of gravitational potential

$$\sigma_{tot}^2 = \sigma_{th}^2 + \sigma_{nth}^2 \quad \sim T \sim \varphi$$

efficiency  $\eta \leq 1$  (characteristic of weak shocks)

## Dissipation

Time scale determined by the turnover time of the largest eddies - doesn't depend on how viscosity works on small scales



"Big whorls have little whorls That feed on their velocity; And little whorls have lesser whorls And so on to viscosity"

 $t_d = \beta t_{dyn} / 2$ 

-- Lewis F. Richardson Weather prediction by numerical processes (1922)

$$\frac{\mathrm{d}\sigma_{\mathrm{nth}}^2}{\mathrm{d}t} = -\frac{\sigma_{\mathrm{nth}}^2}{t_{\mathrm{d}}} + \eta \, \frac{\mathrm{d}\sigma_{\mathrm{tot}}^2}{\mathrm{d}t}$$

## Properties of non-thermal fraction f<sub>nth</sub>

$$t_{growth} = \frac{\sigma_{tot}^2}{d\sigma_{tot}^2/dt} \approx \frac{M}{Mdot}$$

$$\frac{d\sigma_{nth}^2}{dt} = \frac{\sigma_{nth}^2}{t_d} + \eta \frac{d\sigma_{tot}^2}{dt} \qquad growth rate dependence$$

$$\int_{0}^{0} \frac{d\sigma_{nth}^2}{d\sigma_{tot}^2/dt} = \frac{\sigma_{nth}^2}{t_d} + \eta \frac{d\sigma_{tot}^2}{dt} \qquad growth rate dependence$$

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$$\int_{0}^{0} \frac{d\sigma_{nth}^2}{d\sigma_{tot}^2/dt} = \frac{\sigma_{nth}^2}{d\tau} + \frac{\sigma_{nth}^2}{d\tau} + \frac{\sigma_{nth}^2}{d\tau} = \frac{\sigma_{nth}^2}{d\tau} = \frac{\sigma_{nth}^2}{d\tau} + \frac{\sigma_{nth}^2}{d\tau} = \frac{\sigma_{nth}^2}{d\tau} + \frac{\sigma_{nth}^2}{d\tau} = \frac{\sigma_{nth}^2}{d\tau} = \frac{\sigma_{nth}^2}{d\tau} + \frac{\sigma_{nth}^2}{d\tau} = \frac{\sigma_{n$$

## Predicted non-thermal fraction vs simulations

A mass-limited sample of 65 simulated clusters at z=0

Omega500 simulation (Nelson+14)

0.6 0.7  $M_{200m} = 6.4e + 14 \ h^{-1} M_{\odot}$  $\Gamma < 1.8$ , simulated  $\rm M_{200m}$  = 6.8e+14  $\rm h^{-1}\,M_{\odot}$  $\Gamma$  < 1.8, modeled 0.6 0.5  $M_{200m} = 7.2e + 14 h^{-1} M_{\odot}$  $\Gamma$  > 2.7, simulated  $M_{200m} =$  7.6e+14  $h^{-1} M_{\odot}$  $\Gamma$  > 2.7, modeled 0.5  $M_{200m} = 1.8e + 15 \ h^{-1} M_{\odot}$ 0.4  $M_{200m} = 2.3e + 15 h^{-1} M_{\odot}$  $f_{nth}$ 0.4 **ر** 1014 0.3 0.3 0.2 0.2 0.1 sample average 0.1 a few clusters faster/slower growing samples 0.0 0.0 0.2 0.4 0.6 0.8 1.0 0.2 0.4 0.5 0.6 0.7 0.8 0.9 1.0 0.1 0.3 r / r<sub>200m</sub>  $r/r_{200m}$ 

use  $\sigma_{tot}(r,t)$  from simulation as input

reproduce the variation among clusters

both mean & scatter match; confirms the relation between f<sub>nth</sub> & growth rate

note: not all relaxed

Shi, Komatsu, Nelson, Nagai, 2015

## From non-thermal pressure to mass bias

- How well can we correct for HSE mass bias using predicted Prand?
- If we know the accretion histories, can we correct for individual clusters ?

# depends (~10% for relaxed clusters) on the particular pipeline used for estimating the mass (since ICM has structures)

### mass bias

Prand

$$M_{\rm HSE}(< r) \equiv -\frac{r^2}{G\rho_{\rm gas}(r)} \frac{\partial P_{\rm th}}{\partial r}$$

$$M_{\rm corr}(< r) \equiv -\frac{r^2}{G\rho_{\rm gas}(r)} \frac{\partial P_{\rm tot}}{\partial r}$$

the curse of derivative and division - fitting / smoothing necessary

## HSE mass of rather relaxed clusters: 5-10% scatter among different methods



top relaxed 14/65 from X-ray mock of Omega500 clusters

### Corrected mass using <u>simulated</u> P<sub>rand</sub>: much less biased on average



top relaxed 14/65 from X-ray mock of Omega500 clusters

What causes the residues? probably density structure and accelerations

### Corrected mass using <u>predicted</u> P<sub>rand</sub>: much less biased on average, a bit more scatter



top relaxed 14/65 from X-ray mock of Omega500 clusters

### 5-10% scatter between different fitting limits (r500 or 1.5 r500)



On average: larger bias when *fitting to larger radii* (non-thermal pressure more prominent)

### Correction works well for the sample mean, irrespective of methods, fitting range, or even dynamical state of the sample



Using predicted Prand

## Conclusions

A physical motivated 1d model for non-thermal pressure without free parameters,

Key elements:

- Infall kinetic energy converts to turbulence  $(\eta)$  + thermal energy (1- $\eta$ ), mostly by weak internal shocks

Captures behaviors in hydro simulations,

Improves cluster mass estimation

- correcting for individual cluster seems hard due to real life complications
- good for the sample mean in all cases