simulation turbulence in galaxy clusters

insights on stochastic acceleration and the impact of microphysics

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"The Matryoshka Run: Eulerian Refinement Strategy to Model Turbulence in Cosmic Structure" fm, ApJ 782 21 (2014)

"The Matryoshka Run (II): Time Dependent Turbulence Statistics, Stochastic Particle Acceleration and Microphysics Impact in a Massive Galaxy Cluster" fm, ApJ 800 60 (2015)

motivations for turbulence acceleration

- diffuse radio sources appear only in massive systems
- they appear to be triggered by mergers / bi-modality (but see Enßlin+ 2011)
- spectral curvature
- source statistics suggests a lifetime of ca 1 Gyr
- we don't see gamma-rays that would suggest a secondary origin
- other more sophisticated tests but like radial profile requiring some assumption about B or the like

outline

- (computational modeling)
- some properties of turbulence in galaxy clusters
- particle acceleration by turbulence, impact of microphysics of weak shocks

Eulerian Refinement Strategy: Zoom-in + Matryoshka of grids

	l	L (h ⁻¹ Mpc)	Nℓ	n _l	Δx_{ℓ} (h ⁻¹ kpc)	
	0	240	512	2	470	
	1	120	512	2	235	\mathbf{N}
Σ	2	60	512	2	117	
	3	30	512	4	58.6	
	4	15	1024	2	14.6	10
	5	7.5	1024		7.3	

fm, ApJ 782, 21 (2014)



5756.57 -

208.99

7.59

0.28

0.01



Statistics

Hodge-Helmholtz decomposition



Evolution of Turbulence



fm, ApJ 800, 60 (2015)

Particle Acceleration

transport eq.
$$\frac{df}{dt} - \nabla D_{xx} \nabla f - \frac{1}{p^2} \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial f}{\partial p} = 0$$

Fokker-Planck coeff.

$$D_{pp} \equiv \left\langle \frac{\Delta p^2}{2\Delta t} \right\rangle, \qquad D_{xx} \equiv \left\langle \frac{\Delta x^2}{2\Delta t} \right\rangle$$

advection rate in p-space

$$\Gamma_{p} \equiv \frac{\dot{p}}{p} = \frac{1}{p^{3}} \frac{\partial}{\partial p} \left(p^{2} D_{pp} \right) = 4 \frac{D_{pp}}{p^{2}}$$

Particle Acceleration

$$\frac{\partial f}{\partial t} - \nabla D_{xx} \nabla f - \frac{1}{p^2} \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial f}{\partial p} = 0$$

 Transit-Time-Damping (Fisk 1976, BL07):
 2nd order Fermi process, particles resonate with fast MHD waves and get reflected by the mirror force

$$\frac{dp_{\parallel}}{dt} = -\frac{p_{\perp}}{2B^2} \mathbf{v}_{\perp} \left(\vec{B} \cdot \nabla \right) \vec{B}$$

$$D_{pp} = p^2 \frac{\pi I_{\vartheta}}{16c} \langle k \rangle_E \left\langle \left(\delta u_c \right)^2 \right\rangle$$

$$\left\langle k\right\rangle_{E} \equiv \frac{1}{\left\langle \left(\delta u_{c}\right)^{2}\right\rangle} \int^{k_{c}} dk k W(k) \approx \frac{\zeta_{2}}{1-\zeta_{2}} k_{L} \left[\left(\frac{k_{c}}{k_{L}}\right)^{1-\zeta_{2}} - 1 \right]$$

• Non Resonant Mech. (Ptuskin 1988, BL07): stochastic acceleration due to velocity divergence according to the adiabatic process

$$\frac{dp}{dt} = -\frac{p}{3} \nabla \cdot \vec{v}$$

$$\frac{\text{micro}}{\sqrt{2}} \int_{pp}^{\sqrt{2}} \frac{2}{9} \zeta_2 \frac{I_{\xi}}{D} \left(\frac{k_L D}{c_s}\right)^{\zeta_2} \left(\left(\delta u_c\right)^2\right)$$

$$I_{\xi} \equiv \int_{k_L D/c_s}^{k_c D/c_s} d\xi \frac{\xi^{1-\zeta_2}}{1+\xi^2}$$

Spectra of Turbulent Cascade



Spectra of Turbulent Cascade



Brunetti & Jones 2014

Key Cascade Physics

- the cascade of the compressional modes, Alfven vs Burgers: how much dissipation occurs during the cascade ? If enough to steepen the structure functions then the mechanisms become very inefficient
- the slope of the cascade of compressional modes affects
 - i) the value of the energy-averaged wave-vector which tends to k_L as $\zeta \rightarrow 1$
 - ii) the cascade cutoff which can become much larger
- the collisionality of the plasma; if thermal particles have their *mfp* reduced by micro instabilities (mirror, firehose...), they won't resonate with and damp the MHD waves anymore, only CRs do, so their acceleration efficiency increases dramatically

$$D_{pp} = p^2 \frac{C_D C_W}{x_{CR}} \xi k_L \frac{\left\langle \left(\delta u_c\right)^2 \right\rangle}{c_s^3}$$



Time evolution of spectral indexes



kraicknan's cascade for cutfoff + burgers' (simulation's) slope for <k>E no micro instabilities

kraicknan's cascade for cutfoff and <k>E + micro instabilities set *mfp*

takeaway result

- the simulation model of the turbulence pins down an important ingredient entering the acceleration rate, which is the amount of compressional turbulent energy available for TTD or NR mechanisms
- the microphysics of the ICM plasma, however, also enters the acceleration rates, and because we have fixed the above unknown, we can now expose its impact
- the acceleration rates depend on at least the following microphysics (but possibly other as well) of:
- the cascade of the compressional modes, Alfven vs Burgers: how much dissipation occurs during the cascade ? If enough to steepen the structure functions then the mechanisms become very inefficient
- the collisionality of the plasma; if micro instabilities (mirror, firehose...) reduce the thermal particles *mfp* then the acceleration efficiency is very high
- we also need to understand the properties of magnetic fields