Multi-scale, multi-physics structure formation simulations Volker Springel



- The Next Generation Illustris Simulations (IllustrisTNG)
- The role of black holes in galaxy formation simulations
- Modelling directions: Magnetic fields, cosmic rays, radiation hydro, ...





Heidelberg Institute for Theoretical Studies



Stars, Planets, and Galaxies Harnack Haus Berlin, April 2018 The multifaceted need for **feedback**

Star formation in the ISM is surprisingly inefficient

THE GAS CONSUMPTION TIMESCALE OF STAR FORMATION

depletion time:

$$t_{\rm dep} \equiv M_{\rm gas}/\dot{M}_*$$

gravitational free-fall time:

$$t_{\rm ff} = \sqrt{\frac{3\pi}{32G\rho}}$$

dimensionless "efficiency" of star formation:

$$\epsilon_{\rm ff} \equiv \frac{t_{\rm ff}}{t_{\rm dep}}$$

observed is:
 $\dot{\Sigma}_{\star} \simeq \epsilon_{\rm ff} \frac{\Sigma_{\rm H_2}}{t_{\rm ff}}$
 $\epsilon_{\rm ff} \sim 0.01$



Krumholz et al. (2014)

Abundance matching gives the expected halo mass – stellar mass relation in ΛCDM

MODULATION OF GLOBAL STAR FORMATION EFFICIENCY AS A FUNCTION OF HALO MASS



Quasars release plenty of energy

$$L_Q \sim 10^{12} L_{\odot}$$
 $t_Q \sim 10^7 - 10^8 \,\mathrm{yr}$

Black hole energetics suggests that they

could influence the evolution of galaxies

COMPARING SUPERNOVA AND BLACK HOLE ENERGIES

$$E_Q \sim 10^{60} - 10^{61} \,\mathrm{erg}$$

a billion supernovae !

But how does AGN energy couple to halo gas?

quasars / AGN





galaxies

Total available feedback energy from BHs is comparable to that of supernovae

$$\rho_{\rm BH} \simeq 0.001 \,\rho_{\star} \qquad E_{\rm BH}/V \simeq 0.1 \,\rho_{\rm BH} \,c^2 \qquad \qquad \frac{E_{\rm BH}}{E_{\rm SN}} \simeq 1.8$$
$$E_{\rm SN}/V \simeq \frac{10^{51} \,\rm erg}{100 \,\rm M_{\odot}} \,\rho_{\star} \qquad \qquad \frac{E_{\rm BH}}{E_{\rm SN}} \simeq 1.8$$

The dynamic range challenge



x [kpc]

x [kpc]

x [kpc]

x [kpc]

The Illustris and IllustrisTNG cosmological simulations

Galaxy formation physics in the Illustris simulations

 Cooling and metal enrichment Nine elements followed independently 	Star formation and windsVariant of Springel & Hernquist (2003)
 Mass and metal loss of stars treated continuously over time based on stellar population synthesis models (similar to Wirsma et al. 2009) 	 Cold dense gas stabilized by an ISM equation of state Winds are phenomenologically introduced, with an energy given as a fixed fraction of the
 Ionization balance and cooling from H and He followed with direct chemical network (Katz et al. 1996) 	 supernova energy The wind velocity is variable, the mass flux follows for energy-driven winds
 Metal line cooling added through CLOUDY lookup tables in density, temperature and redshift 	 Fiducial model scales wind with local dark matter velocity dispersion Winds are launched outside of star-forming
 Simple self-shielding correction (Rahmati et al. 2013) 	gas, and metal-loading can be reduced if desired

Black hole accretion and feedback

- Black hole seeding and accretion model (Springel et al. 2005)
- Quasar-mode feedback for high accretion rates
- Radio-mode feedback for low accretion rates based on bubble-heating model (Sijacki et al. 2006)
- Radiative AGN feedback (change in heating/cooling due to variation of UVB) in proximity to an active black hole
- Reduction of accretion rate in low-pressure/low-density regimes to avoid large hot bubbles around black holes in quiescent state
- Black holes tied to potential minimum of halos

Illustris Simulation

Vogelsberger, Genel, Springel, Torrey, Sijacki, Xu, Snyder, Bird, Nelson, Hernquist

The Illustris simulation reproduces the morphological mix of galaxies SIMULATED HUBBLE TUNING FORK DIAGRAM











ellipticals











disk galaxies



















The stellar mass functions match observations at high redshift well

STELLAR MASS FUNCTIONS OF ILLUSTRIS COMPARED TO OBSERVATIONS AT DIFFERENT EPOCHS



Genel et al. (2014)

People in the "Next Generation Illustris Simulations" team

A COLLABORATION BETWEEN HEIDELBERG, HARVARD, AND THE MIT



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Massachusetts Institute of Technology Jet injection by BHs can be modelled at high-resolution in individual "zoom" galaxy cluster simulations HEATING CLUSTERS THROUGH JET-ICM INTERACTIONS

Weinberger et al. (2017)

Jet after 42 Myr



ICM after lobe passage (168 Myr)



Intense winds from supermassive black holes may keep 'red geyser' galaxies turned off

EXAMPLE OF OBSERVATIONAL EVIDENCE FOR BLACK HOLE WINDS



(Genzel et al. 2014, Förster Schreiber et al. 2014)

We have adopted a simple model for kinetic AGN winds in cosmological simulations of galaxy formation

ILLUSTRIS-TNG AGN MODEL AND ITS IMPACT

 $\Delta \dot{E}_{\rm high} = \epsilon_{\rm f, high} \epsilon_{\rm r} \dot{M}_{\rm BH} c^2$ $\dot{M}_{\rm Bondi} = \frac{4\pi G^2 M_{\rm BH}^2 \rho}{c^3}$ We distinguish a "high" and a "low" accretion flow state: Pure kinetic feedback in $\frac{\dot{M}_{\rm Bondi}}{\dot{M}_{\rm Edd}} \ge \chi = \min \left| \chi_0 \left(\frac{M_{\rm BH}}{10^8 \,{\rm M}_\odot} \right)^\beta, 0.1 \right|$ "low" radio mode: $\dot{M}_{\rm Edd} = \frac{4\pi G M_{\rm BH} m_p}{\epsilon_r \sigma_{\rm T} c}$ $\Delta \dot{E}_{\rm low} = \epsilon_{\rm f,kin} \dot{M}_{\rm BH} c^2$ $M_{
m star,10}$ SHMR $/M_{200,
m c}/(\Omega_b/\Omega_m)$ [%] high res 10^{10} 25Behroozi+ (2013), including ICL Kormendy + Ho (2013) -220 10^{9} $M_{
m BH} \, [{
m M}_{\odot}]$ $\log_{10}(\dot{m})$ 15 10^{8} 10 10^{7} -55 10^{6} -6 10^{11} 10^{14} 10^{12} 10^{13} 10^{11} 10^{10} 10^{12} 10^{9} 10^{8} 10^{7} $M_{200,c} \, [M_{\odot}]$ Weinberger et al. (2017) $M_{\rm bulge} \, [{\rm M}_{\odot}]$

"high" quasar mode:

We obtain sudden quenching, setting in at around $M_* \sim 2x10^{10} M_{sun}$

The gas fractions in galaxy groups and poor clusters provide a sensitive constraint on viable AGN feedback models

GAS FRACTIONS WITH THE NEW ILLUSTRIS-TNG AGN MODEL



Weinberger et al. (2017)

The "Next Generation Illustris Simulations" (IllustrisTNG) are our novel, significantly improved models for cosmic structure formation

IllustrisTNG Collaboration (2017)



Overview over the simulation set of IllustrisTNG

RUNS WERE DONE ON HAZEL-HEN AT HLRS WITH CPU-TIME GRANTED BY GCS

Series	Run	Boxsize		$N_{ m gas}$	N _{dm}	N _{tracer}	$m_{\rm b}$	<i>m</i> _{dm}	ϵ
		$[h^{-1}Mpc]$	[Mpc]				$[h^{-1}\mathrm{M}_{\odot}]$	$[h^{-1}\mathrm{M}_{\odot}]$	$[h^{-1}kpc]$
TNG300	TNG300(-1)	205	302.6	2500 ³	2500 ³	2500 ³	7.44×10^{6}	3.98×10^{7}	1.0
up to 24000 cores, 35 Mio core-h	TNG300-2	205	302.6	1250^{3}	1250^{3}	1250^{3}	5.95×10^{7}	3.19×10^{8}	2.0
	TNG300-3	205	302.6	625^{3}	625^{3}	625^{3}	4.76×10^{8}	2.55×10^{9}	4.0
	TNG300-DM(-1)	205	302.6		2500^{3}			4.73×10^{7}	1.0
	TNG300-DM-2	205	302.6		1250^{3}			3.78×10^{8}	2.0
	TNG300-DM-3	205	302.6		625 ³			3.03×10^{9}	4.0
TNG100	TNG100(-1)	75	110.7	1820 ³	1820 ³	2×1820^{3}	9.44×10^{5}	5.06×10^{6}	0.5
up to 10720 cores, 18 Mio core-h	TNG100-2	75	110.7	910 ³	910 ³	2×910^{3}	7.55×10^{6}	4.04×10^{7}	1.0
	TNG100-3	75	110.7	455 ³	455 ³	2×455^{3}	6.04×10^{7}	3.24×10^{8}	2.0
	TNG100-DM(-1)	75	110.7		1820^{3}			6.00×10^{6}	0.5
	TNG100-DM-2	75	110.7		910 ³			4.80×10^{7}	1.0
	TNG100-DM-3	75	110.7		455 ³			3.84×10^{8}	2.0
TNG50	TNG50-1	35	51.7	2160 ³	2160 ³	2160 ³	5.74×10^{4}	3.07×10^{5}	0.2
up to 24000 cores, 90 Mio core-h	TNG50-2	35	51.7	1080^{3}	1080^{3}	1080^{3}	4.59×10^{5}	2.46×10^{6}	0.4
	TNG50-3	35	51.7	540 ³	540^{3}	540^{3}	3.67×10^{6}	1.97×10^{7}	0.8
	TNG50-4	35	51.7	270^{3}	270^{3}	270^{3}	2.94×10^{7}	1.57×10^{8}	1.6
	TNG50-DM-1	35	51.7		2160^{3}			3.64×10^{5}	0.2
	TNG50-DM-2	35	51.7		1080^{3}			2.92×10^{6}	0.4
	TNG50-DM-3	35	51.7		540^{3}			2.33×10^{7}	0.8
	TNG50-DM-4	35	51.7		270^{3}			1.87×10^{8}	1.6

IllustrisTNG was executed on Hazel Hen in Stuttgart, Germany





IllustrisTNG gives detailed predictions for the spatial distribution of stars, and its relation to the underlying dark matter backbone

DARK MATTER AND STELLAR DENSITY IN ILLUSTRIS-TNG



Pillepich et al. (2018)

Dark Matter Column Density [log M_{sun} kpc⁻²]

IllustrisTNG reproduces the observed color-bimodality of galaxies thanks to AGN feedback

COLOR DISTRIBUTION OF GALAXIES OF DIFFERENT MASS COMPARED TO SDSS



Nelson et al. (2017)

The new feedback model in TNG produces a sharp characteristic scale in the halo – stellar mass relationship

ILLUSTRIS COMPARED TO ABUNDANCE MATCHING MODELS FOR THE SMHM RELATION

Pillepich et al. (2018)



Kinetic AGN feedback takes over at late times in large galaxies

FEEDBACK ENERGY AS A FUNCTION OF TIME AND FINAL STELLAR MASS

Weinberger et al. (2018)



IllustrisTNG predicts galaxy correlation functions in good agreement with the most accurate galaxy surveys

PROJECTED TWO-POINT FUNCTIONS IN DIFFERENT MASS BINS



IllustrisTNG predicts pronounced differences in the clustering of red and blue galaxies in good agreement with data

CLUSTERING IN DIFFERENT MASS AND COLOR BINS COMPARED TO SDSS



IllustrisTNG predicts that the correlation length of galaxy clustering depends both on stellar mass and redshift

CORRELATION LENGTH AS A FUNCTION OF REDSHIFT AND STELLAR MASS





TNG allows us to measure the scale-dependence of the power spectrum of different tracers relative to the non-linear clustering of all the matter

SCALE-DEPENDENT BIAS IN THE POWER SPECTRUM OF DIFFERENT TRACERS





Black hole growth influences disk sizes

BLACK HOLE GROWTH BETWEEN Z=1 AND Z=0 CORRELATES WITH DISK SCALE LENGTHS IN HALOS OF MILKY WAY MASS



Grand et al. (2016)

The sizes of different galaxies types reproduce observed trends with stellar mass well

ILLUSTRIS-TNG GALAXY SIZES AS A FUNCTION OF STELLAR MASS



Genel et al. (2018)

Observed metallicity profiles of galaxy clusters are reproduced by IllustrisTNG

METALLICITY MAPS AND PROFILES FOR RICH GALAXY CLUSTERS



Ζ Fe Ettori + 2015- Mernier + 2017(scaled) Leccardi + 2008Leccardi + 2008 . 10⁰ ${\rm Fe/Fe}_{\odot}$ Z/Z ⊙, z = 0.210⁻¹ TNG100-1 10^{-2} 10^{-1} 10^{0} $r/r_{\rm 500,crit}$ 1.0 z = 0.10.8 TNG100-1 0.6[°] 2/2 0.4 0.2 Ettori + 2015 (NCC, z < 0.2)non-cool core cool core Ettori + 2015 (CC, z < 0.2)0.0 և 10⁻² 10^{-1} 10⁰ $/r_{
m 500,crit}$

Vogelsberger et al. (2017)

Magnetic field results

Our MHD simulations of galaxy formation predict the amplification of primordial fields in halos and galaxies

MAGNETIC FIELD STRENGTH IN A SMALL REGION OF ILLUSTRIS-TNG



The IllustrisTNG simulations predict the magnetic field strength and field topology in galaxies

STELLAR MASS, GAS DENSITY AND MAGNETIC FIELD STRENGTH IN A TYPICAL SPIRAL



Marinacci et al. (2018)

The AURIGA simulations yield disk galaxies with stellar mass and spatial structure consistent with Milky Way-like galaxies

FACE-ON STELLAR MASS DENSITY PROJECTIONS OF TWO SIMULATIONS FROM THE AURIGA SUITE



Grand, Springel, Pakmor, et al. (2017)

A small-scale dynamo is active at very high redshift

EVOLUTION OF THE VOLUME-WEIGHTED RMS B-FIELD STRENGTH INSIDE 10 KPC



Pakmor et al. (2017)

Amplification of B-field occurs through turbulent dynamo

VELCOITY FIELD AND EVOLUTION OF VELOCITY AND B-FIELD POWER SPECTRA



The predicted magnetic field strength agrees quite well with observations

PROFILES OF MAGNETIC FIELD STRENGTH IN SIMULATIONS AND OBSERVATIONS



Cosmic rays and their impact on galactic outflows

In isolated disk galaxy formation simulations, magnetic fields drive magnetically driven small fountain like flows out of the disk slices THROUGH THE GAS DENSITY AND THE MAGNETIC FIELD



Pakmor & Springel (2012)

The CR dynamics is coupled to magnetic fields permeating the gas INTERACTIONS OF COSMIC RAYS AND MAGNETIC FIELDS

Cosmic rays scatter on magnetic fields – this lets them exert a pressure on the thermal gas, and diffuse relative to its rest frame.

Streaming instability:



- CRs can in principle move rapidly along field lines (with c), which acts to reduce any gradient in their number density.
- But if $c_s > v_A$, CR excite Alfven waves (streaming instability)
- scattering off this wave field in turn limits the CR bulk speed to a much smaller, effective streaming speed v_{str}

• streaming speed:
$$\mathbf{v}_{str} = -v_{str} \frac{\nabla P_{cr}}{|\nabla P_{cr}|}$$
 $v_{str} = \lambda \max(c_S, v_A)$
 $\lambda \sim 1$

The CR transport complicates fluid dynamics considerably COSMIC RAY DYNAMICS WITHOUT SOURCE AND SINK TERMS

 $\frac{\partial U}{\partial t} + \nabla \cdot \mathbf{F} = S$ $\boldsymbol{U} = \begin{pmatrix} \rho \\ \rho \boldsymbol{v} \\ \rho \boldsymbol{v} \\ \varepsilon \\ \varepsilon_{\rm cr} \\ \boldsymbol{B} \end{pmatrix}, \quad \boldsymbol{F} = \begin{pmatrix} \rho \boldsymbol{v} \\ \rho \boldsymbol{v} \boldsymbol{v}^{\rm T} + P \mathbf{1} - \boldsymbol{B} \boldsymbol{B}^{\rm T} \\ (\varepsilon + P) \boldsymbol{v} - \boldsymbol{B} (\boldsymbol{v} \cdot \boldsymbol{B}) \\ \varepsilon_{\rm cr} \boldsymbol{v} + (\varepsilon_{\rm cr} + P_{\rm cr}) \boldsymbol{v}_{\rm st} - \kappa_{\varepsilon} \boldsymbol{b} (\boldsymbol{b} \cdot \nabla \varepsilon_{\rm cr}) \\ \boldsymbol{B} \boldsymbol{v}^{\rm T} - \boldsymbol{v} \boldsymbol{B}^{\rm T} \end{pmatrix}, \quad \boldsymbol{S} = \begin{pmatrix} \boldsymbol{0} \\ \boldsymbol{0} \\ P_{\rm cr} \nabla \cdot \boldsymbol{v} - \boldsymbol{v}_{\rm st} \cdot \nabla P_{\rm cr} + \Lambda_{\rm th} + \Gamma_{\rm th} \\ -P_{\rm cr} \nabla \cdot \boldsymbol{v} + \boldsymbol{v}_{\rm st} \cdot \nabla P_{\rm cr} + \Lambda_{\rm tr} + \Gamma_{\rm cr} \\ \boldsymbol{0} \end{pmatrix}$ $\boldsymbol{v}_{\rm st} = -\frac{\boldsymbol{B}}{\sqrt{\rho}}\operatorname{sgn}(\boldsymbol{B}\cdot\boldsymbol{\nabla}\boldsymbol{P}_{\rm cr})$ $\varepsilon = \varepsilon_{\rm th} + \frac{\rho v^2}{2} + \frac{B^2}{2}$ $P = P_{\rm th} + P_{\rm cr} + \frac{B^2}{2}$

cosmic ray streaming, nasty(!) numerically

Energy equation:

$$\frac{\partial \boldsymbol{\varepsilon}_{\rm cr}}{\partial t} + \boldsymbol{\nabla} \cdot \left[\boldsymbol{\varepsilon}_{\rm cr}(\boldsymbol{v} + \boldsymbol{v}_{\rm st}) - \kappa_{\varepsilon} \boldsymbol{b} \left(\boldsymbol{b} \cdot \boldsymbol{\nabla} \boldsymbol{\varepsilon}_{\rm cr} \right) \right] = -P_{\rm cr} \, \boldsymbol{\nabla} \cdot \left(\boldsymbol{v} + \boldsymbol{v}_{\rm st} \right) + \Lambda_{\rm cr} + \Gamma_{\rm cr}$$

anisotropic diffusion

Transport processes of CRs are critical for driving galactic winds

COMPARISON OF DISK GALAXY EVOLUTION WITH DIFFERENT COSMIC RAY PHYSICS



The efficiency of cosmic ray driven winds is a strong function of halo mass

OUTFLOW VELOCITIES IN ISOLATED GALAXIES OF DIFFERENT SIZE

Jacob et al. (2017)

The mass loading of CR-driven winds depends strongly on halo mass, whereas the energy loading is flat

PROPERTIES OF CR DRIVEN WINDS AS A FUNCTION OF HALO MASS

Jacob et al. (2017)

The unknown CR diffusion constant and injection efficiency represent important uncertainties

MASS LOADING OF CR-DRIVEN WINDS FOR DIFFERENT MODEL PARAMETERS

Jacob et al. (2017)

Take home points

- Hydrodynamic cosmological simulations of ACDM can now quite successfully predict galaxy morphologies, clustering, and quenching, as well as diffuse gas properties in the CGM, IGM, and ICM, from high-z to the present
- Black hole feedback processes are critical for galaxy formation and evolution
- A small-scale dynamo driven by feedback generated turbulence can efficiently amplify tiny seed fields already at high redshift to the observed strength
- Cosmic rays may play a key role in explaining galactic outflows