



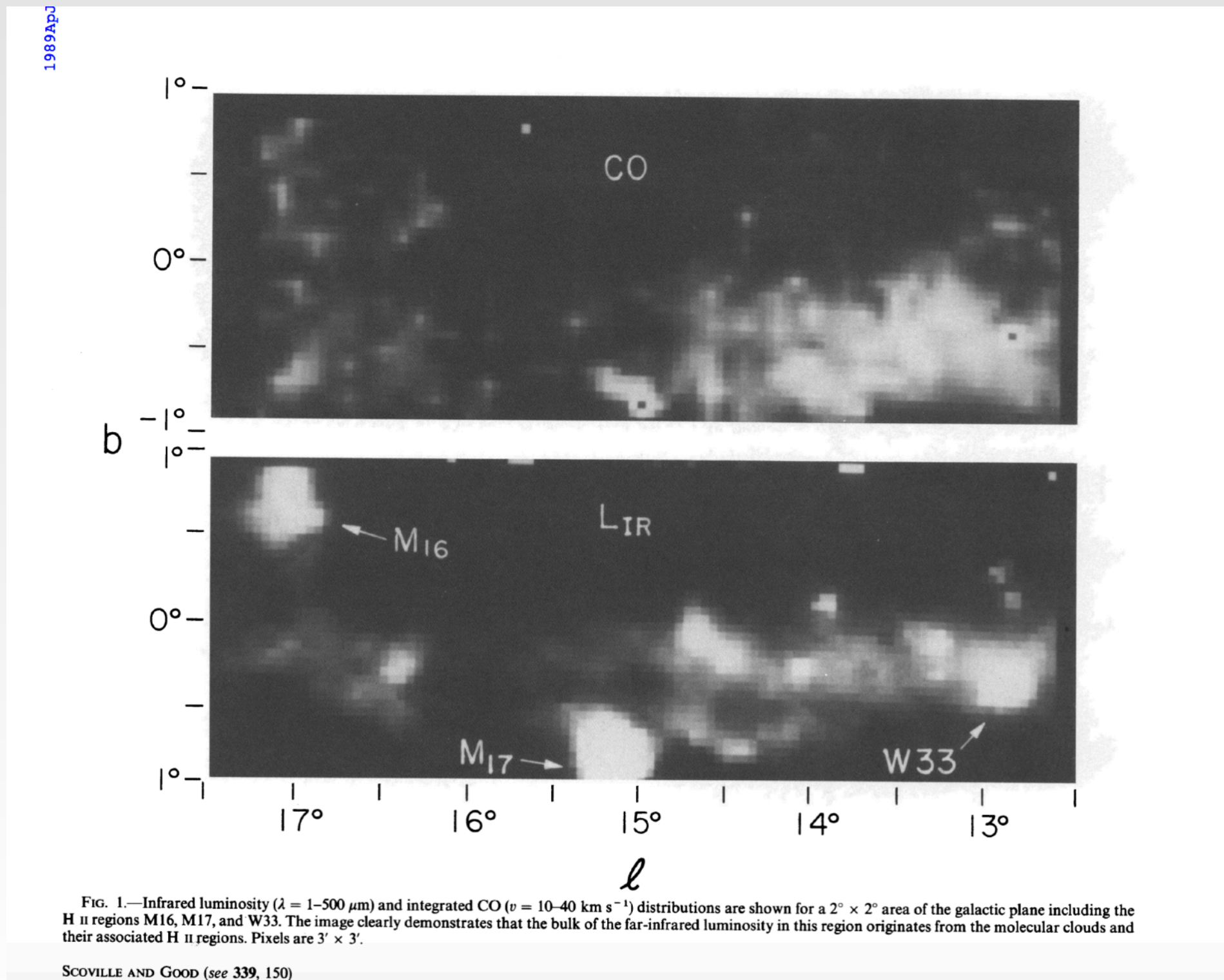
The dynamics and time-dependent mass accretion rates of protostars

Marc-Antoine Miville-Deschenes, Eve Lee, Phil Chang, Daniel Murray, Gunjan
Lakhlani, Laura Keating
The FIRE collaboration
Stars, Planets, & Galaxies, 16 April 2018

Outline

- Introduction: Star formation and giant molecular clouds (GMCs)
- A new GMC catalog, capturing 98% of the flux in the Dame map
- Large scatter in $F_{\text{free-free}}/F_{\text{CO}}$ and $F_{\text{FIR}}/F_{\text{CO}}$, inconsistent with a constant star formation rate
- Feedback from stars is crucial to halt a runaway
- Feedback will likely also drive winds
- FIRE simulations find winds; reproduce the $M_{\text{star}}-M_{\text{halo}}$ relation
- FIRE predicts large variations in $L_{\text{CO}}(z)$; Intensity Mapping will test this
- Conclusions

Stars Form in GMCs



The star formation rate

No. 1, 1988

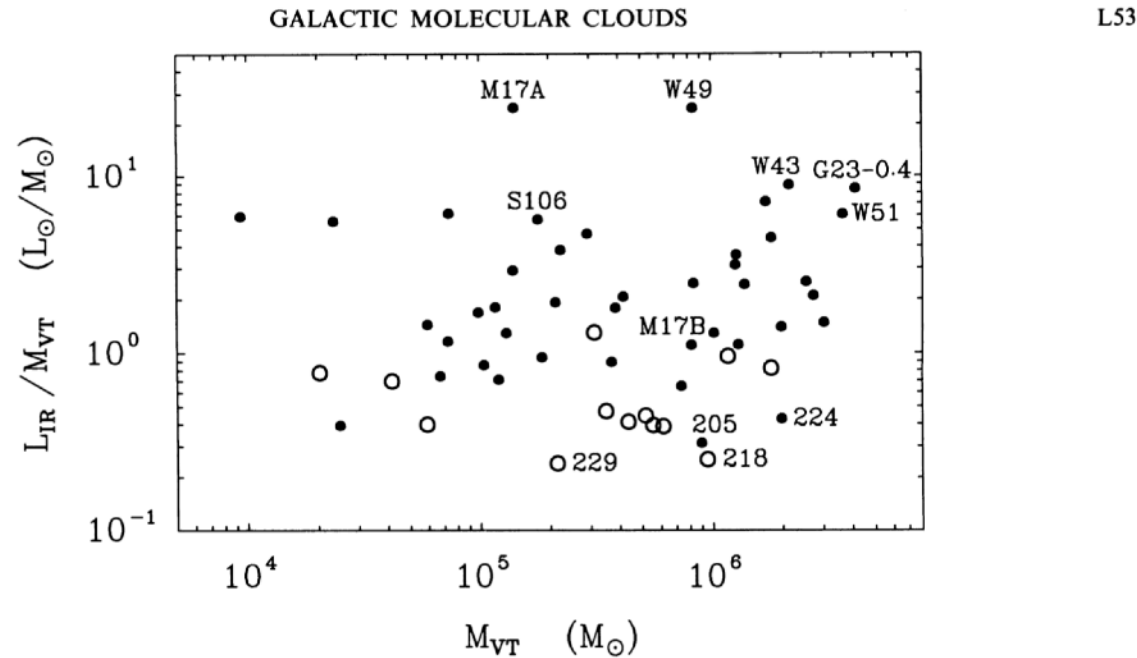


FIG. 2.—The ratio of far-infrared luminosity to cloud mass as a function of cloud mass. The solid and open circles are the same as in Fig. 1. The star formation rate per unit of available molecular mass as measured by $L_{\text{IR}}/M_{\text{VT}}$ is independent of the mass of the cloud.

Mooney & Solomon 1988

THE ASTROPHYSICAL JOURNAL, 833:229 (15pp), 2016 December 20

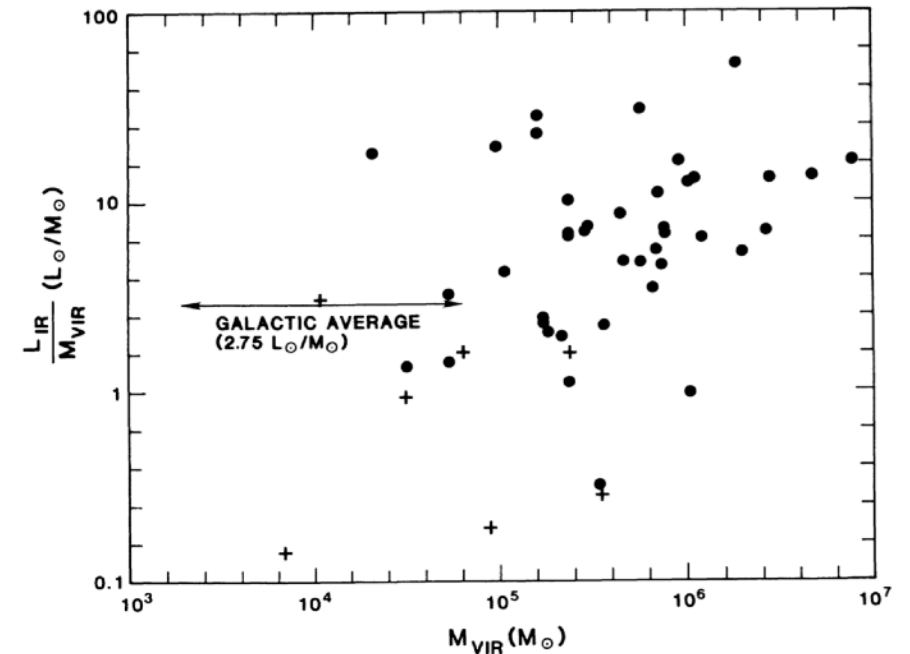
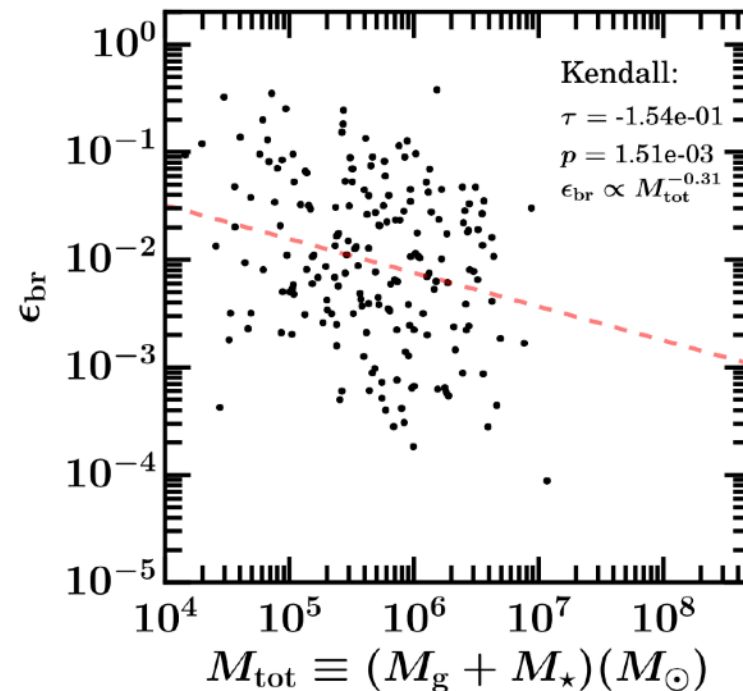
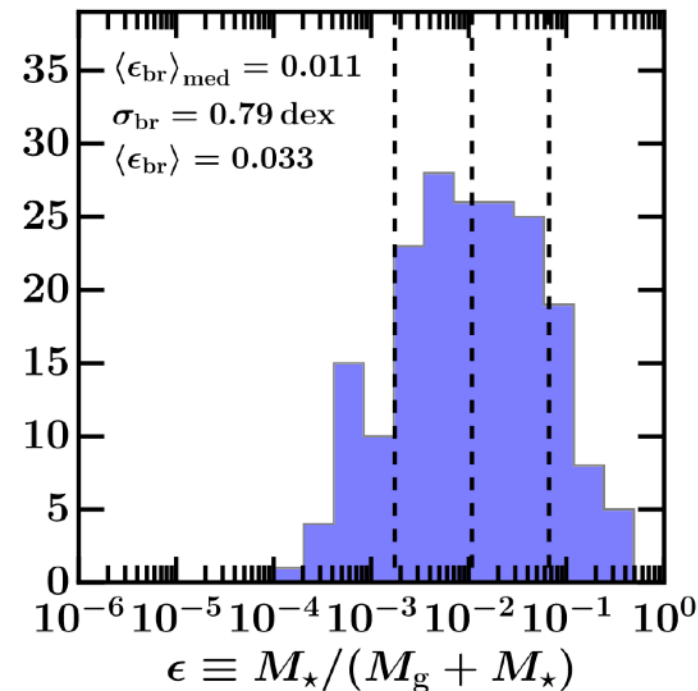


FIG. 9.—The luminosity-to-mass ratio $L_{\text{IR}}/M_{\text{VIR}}$ is shown as a function of virial mass for both cloud samples. No correlation is found between the luminosity-to-mass ratio and cloud masses.

157

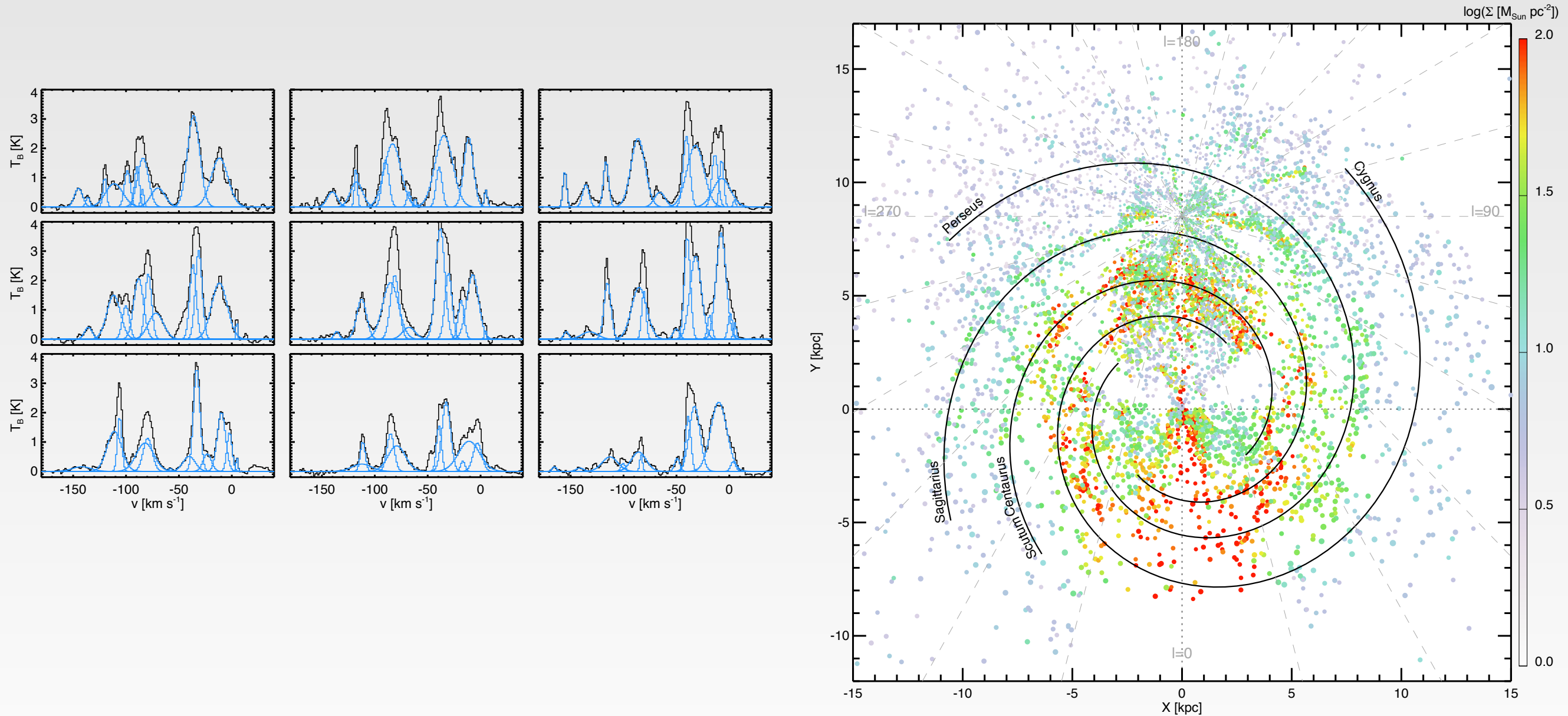
Scoville & Good 1989

LEE, MIVILLE-DESCHÊNES, & MURRAY

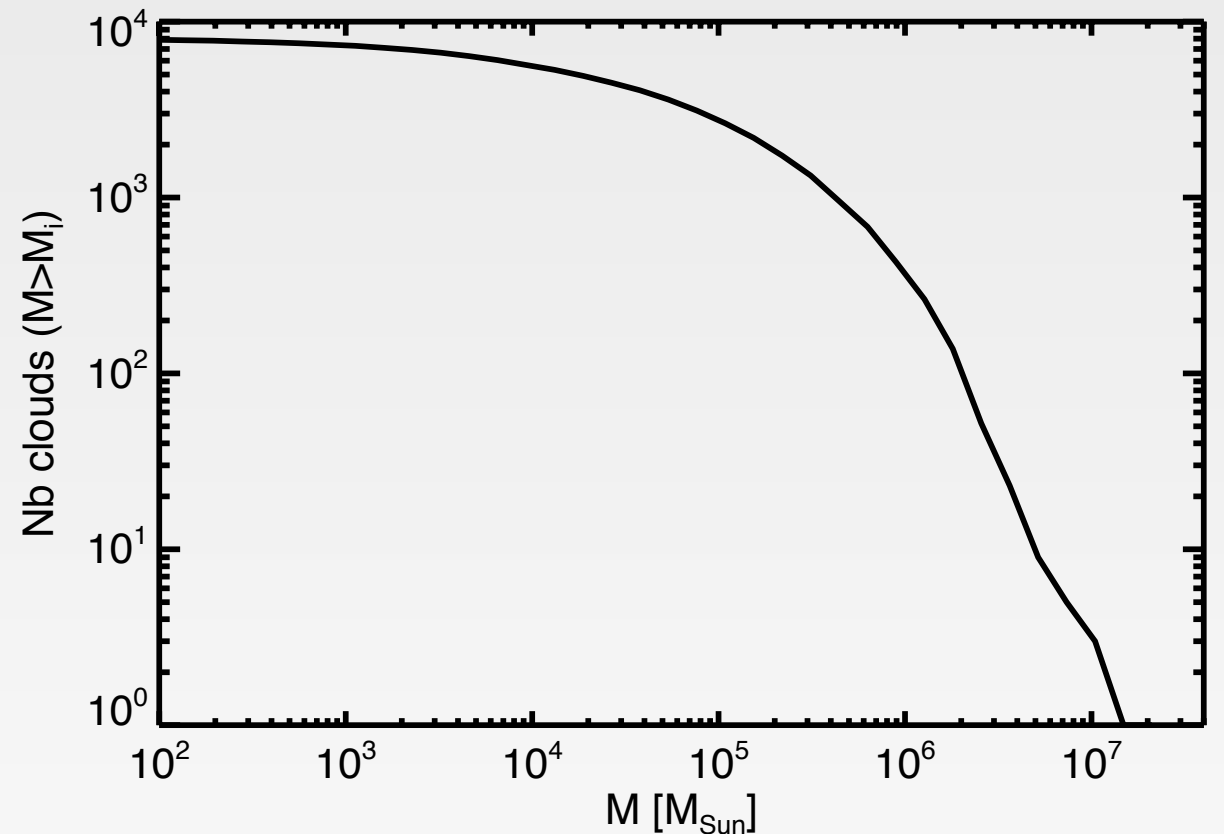
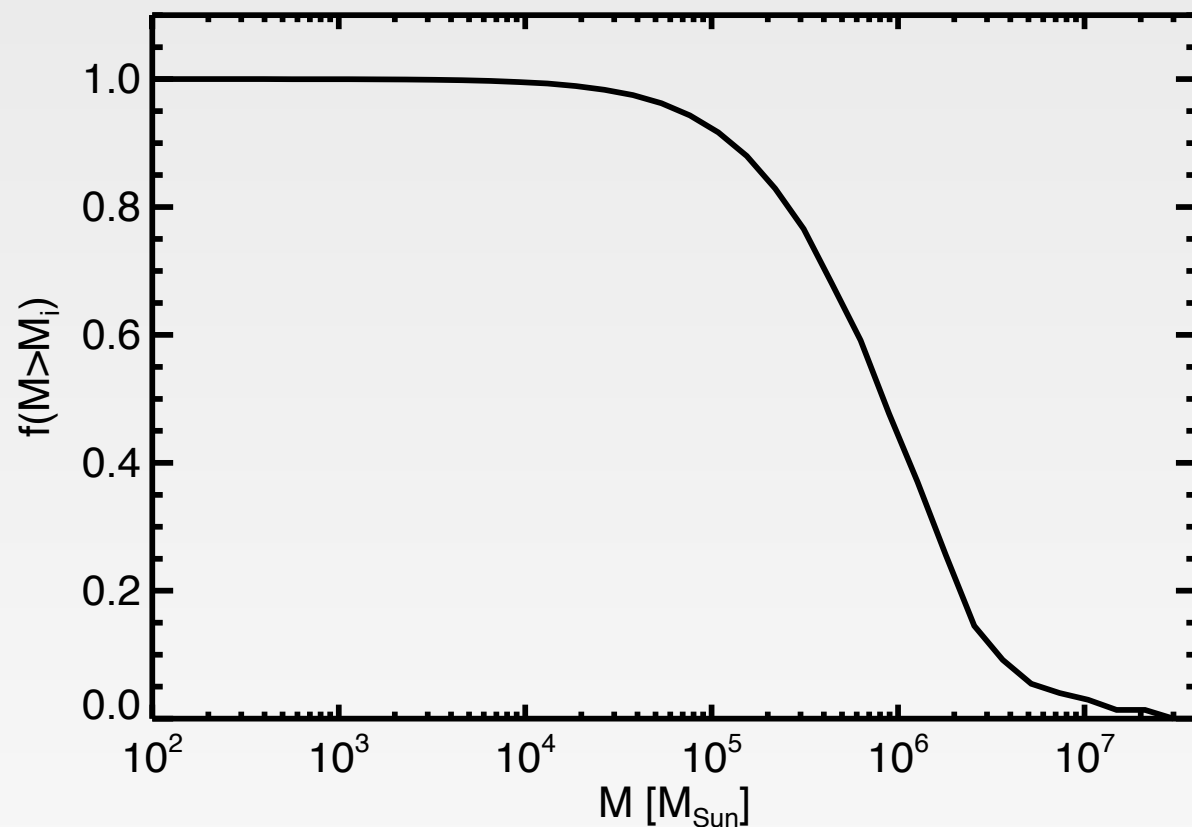


Is $dM_{\star}/dt$ constant?

Introducing a new Milky Way GMC catalog

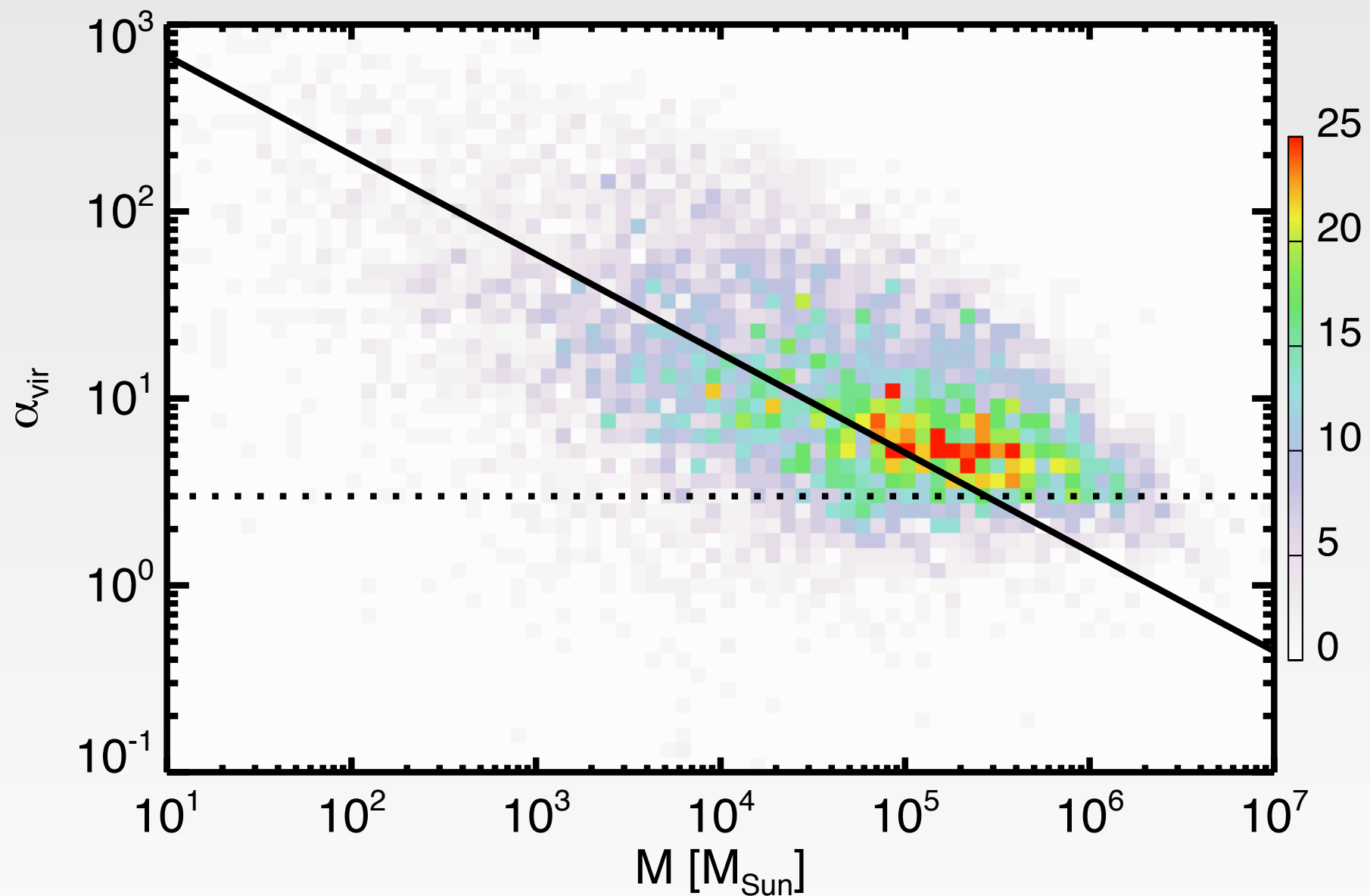


Most of the molecular gas is in about 500 massive clouds



Half the molecular mass is in clouds with $M > 8.4 \times 10^5 M_{\odot}$
There are about 500 such clouds, out of ~ 9000 total
The catalog recovers 98% of the CO flux in the map

Miville-Deschenes, Murray, & Lee



The GMC virial parameter decreases with increasing GMC mass. Many of the $10^6 M_{\odot}$ GMCs are gravitationally bound

Variations with R_{gal}

11

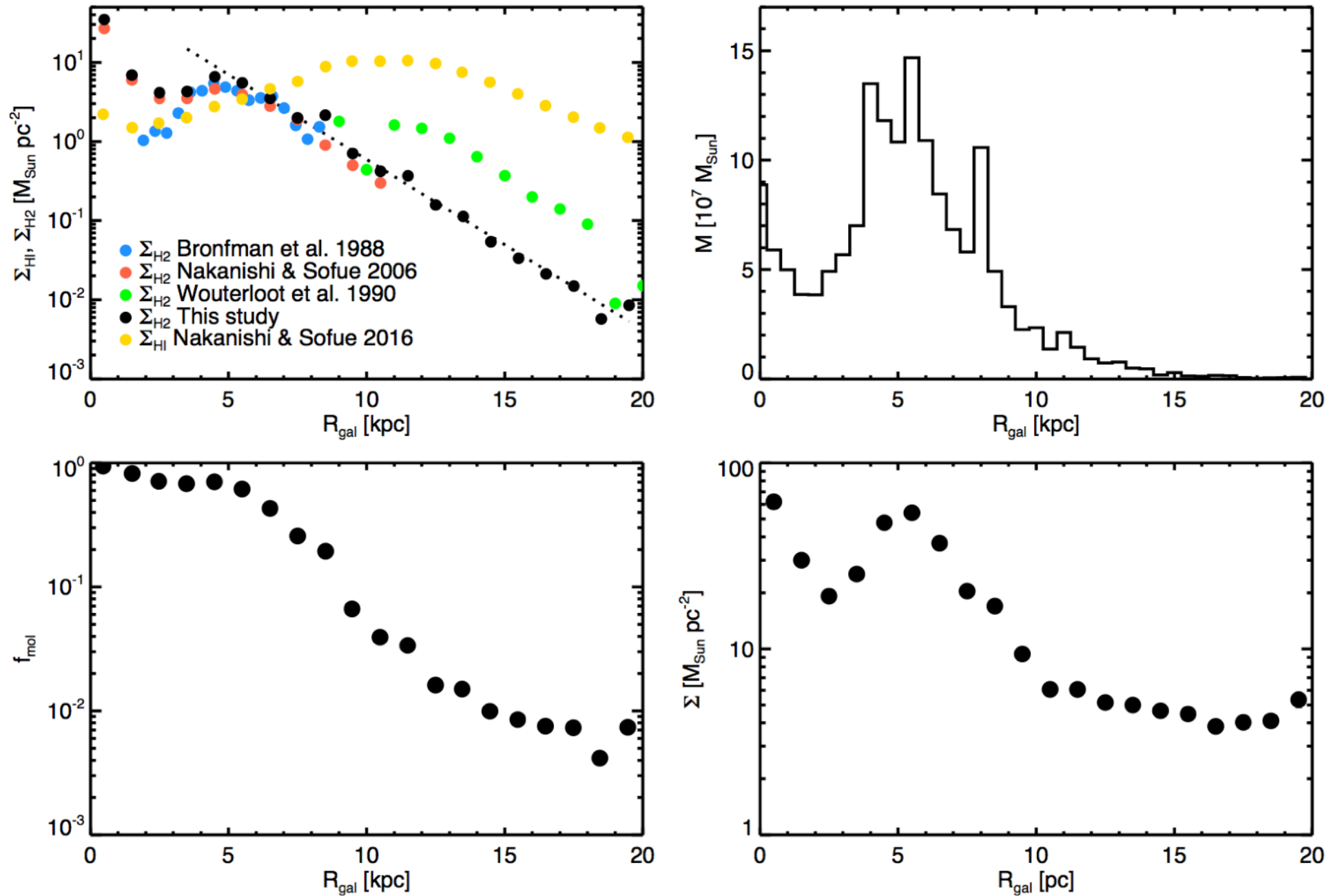


Figure 9. **Top left:** Surface density of H_2 mass as a function of galacto-centric radius (black points). The blue, red and green points correspond respectively to the estimate of Bronfman et al. (1988), Nakanishi & Sofue (2006) and Wouterloot et al. (1990). The dotted line is the exponential fit over the range $4 < R_{\text{gal}} < 17$ kpc: $\Sigma = 83 \exp(-R_{\text{gal}}/2.0)$. The yellow points corresponds to the HI surface density of Nakanishi & Sofue (2016). **Top right:** Total mass of clouds in Galacto-centric rings (thickness 0.5 kpc), as a function of R_{gal} . **Bottom left:** Molecular fraction $f_{\text{mol}} = \Sigma_{\text{H}_2}/(\Sigma_{\text{HI}} + \Sigma_{\text{H}_2})$ built using our estimate of Σ_{H_2} and Σ_{HI} from Nakanishi & Sofue (2016). **Bottom right:** Variation of the median cloud mass surface density Σ as a function of Galactocentric radius.

Mean line width versus R_{gal}

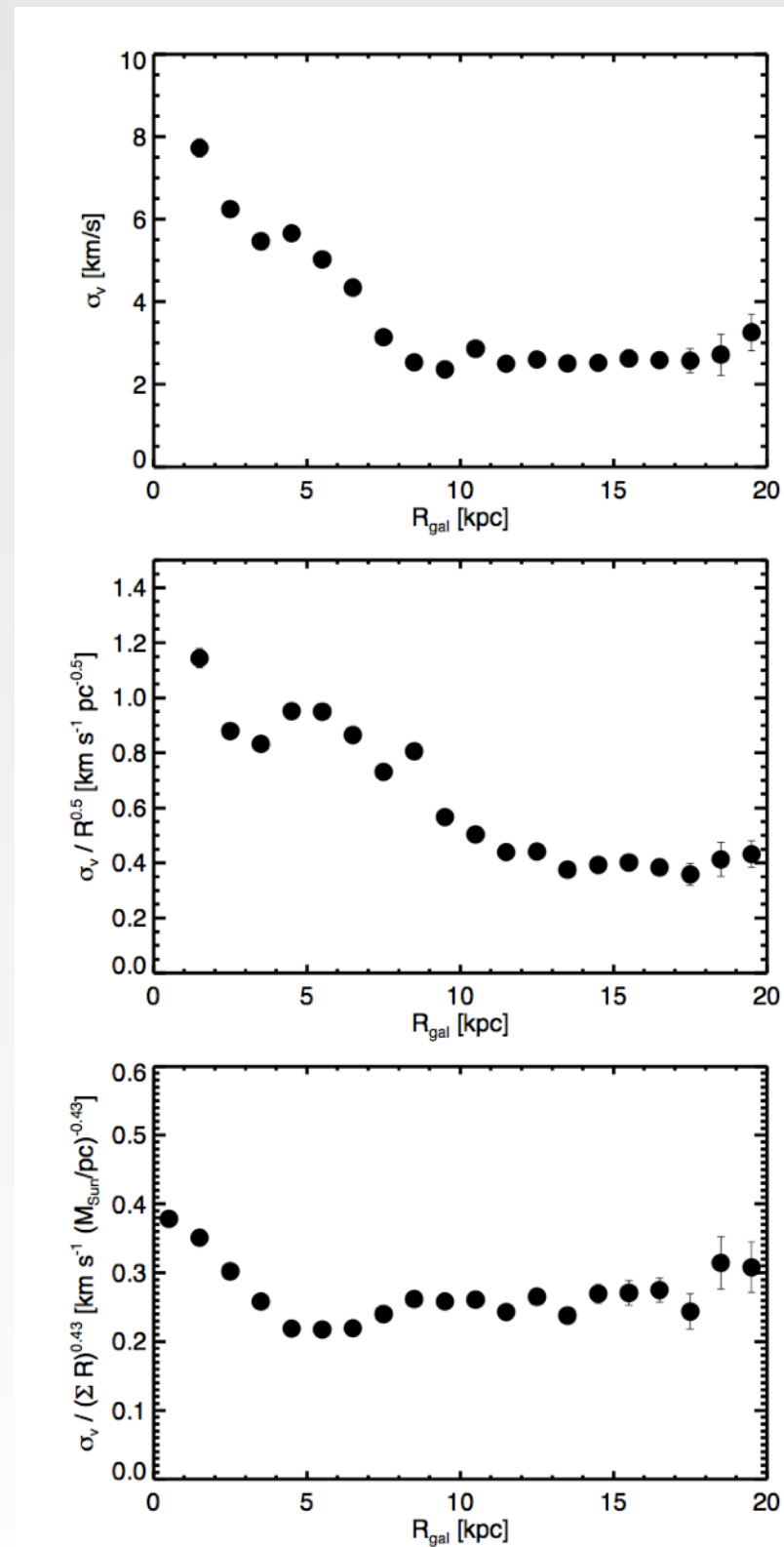
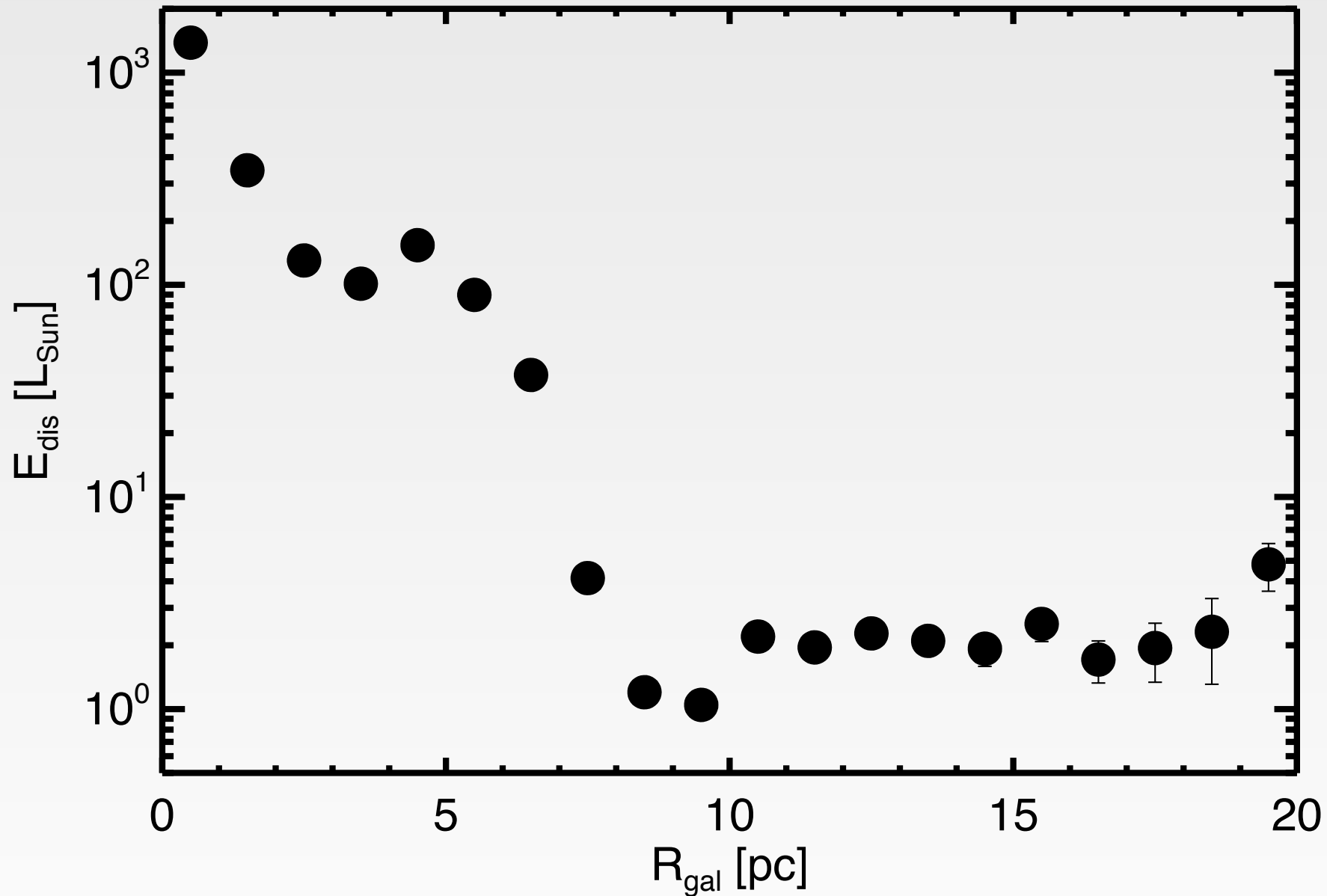


Figure 13. Variation of σ_v (top), $\sigma_v/R^{1/2}$ (middle) and $\sigma_v/(\Sigma R)^{0.43}$ (bottom) as a function of Galactocentric radius. In all panels each dot and its associate error bar indicate respectively the median and 1σ dispersion of the values in a bin of R_{gal} .

Turbulent Luminosity

$$\dot{E}_{\text{dis}} = \frac{1}{2} \frac{M \sigma_v^3}{R}$$

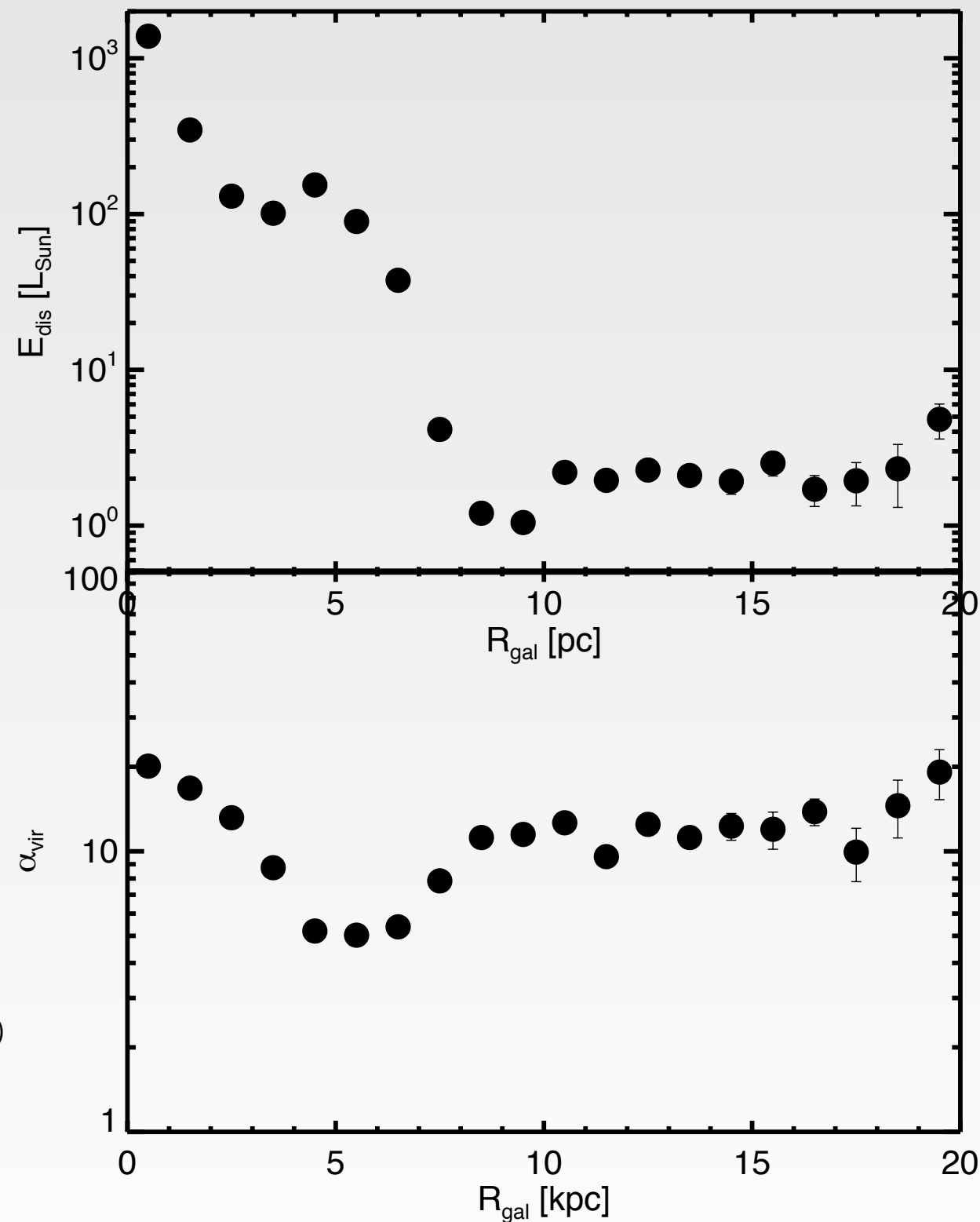


The average energy dissipation rate in annuli. In the outer Galaxy, the dissipation rate $L_{\text{Turb}}=2L_{\odot}$ can be supplied by the flows that create the GMCs. The inner Galaxy is another story, since $L_{\text{Turb}}=100L_{\odot}$; assembly of GMCs does not supply enough power.

What powers the turbulence in individual GMCs?

- Accretion through the Galactic disk (its an accretion disk!) driven by the MRI, or by gravitational torques due to spiral arms
- Accretion onto GMCs
- Contraction of GMCs
- Stellar feedback: protostellar jets, radiation pressure, HII region, stellar winds, supernovae

Virial Parameter vs R_{Gal}



Lee, Murray, & Miville-Deschenes (2016)

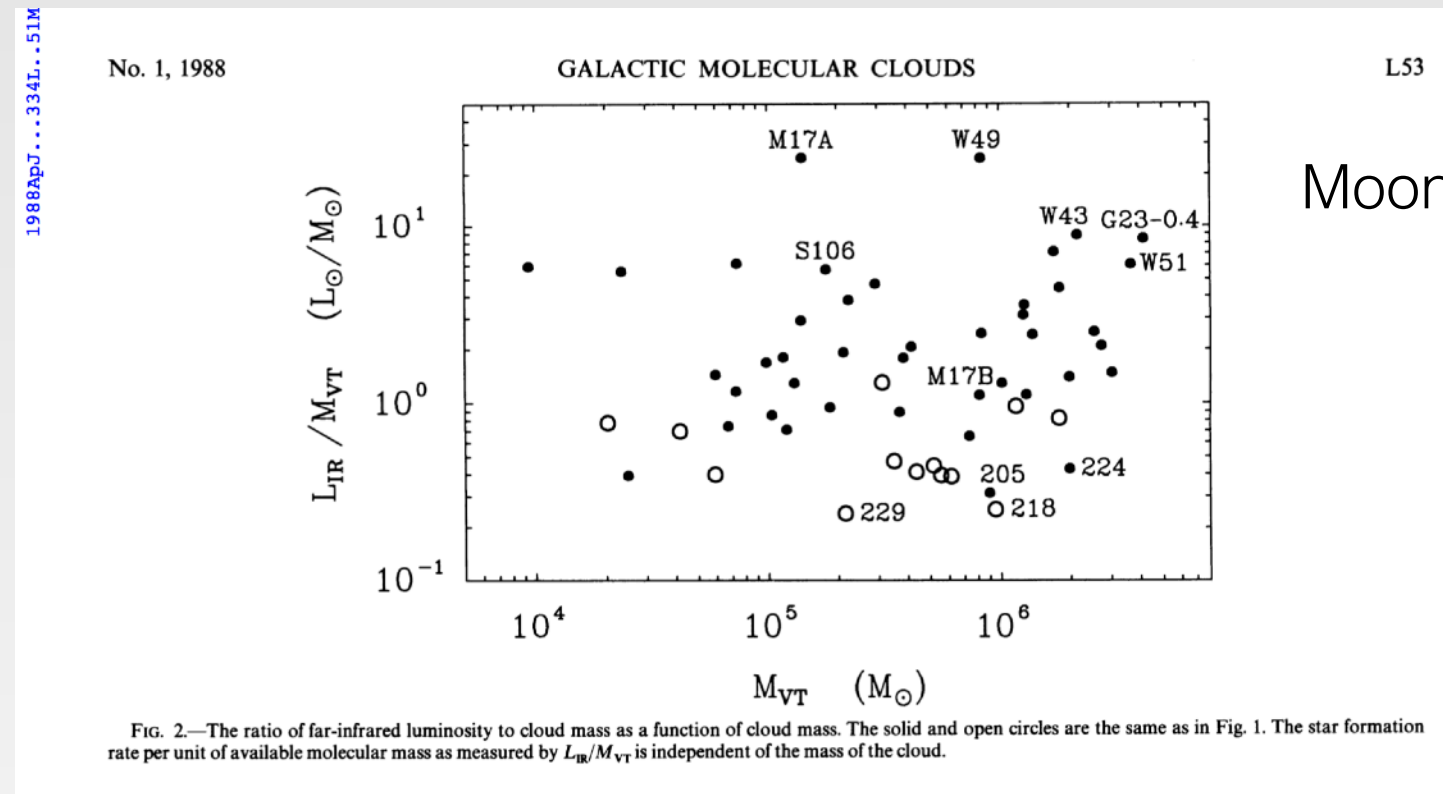
Contraction powers the turbulence in the molecular ring GMCs,
Accretion through the disk $R_{\text{gal}} < 3 \text{ kpc}$

If massive GMCs are bound or collapsing, why don't they all show star formation?

If massive GMCs are bound or collapsing, why don't they all show star formation?

They are collapsing, but the star formation rate starts small, then increases with time

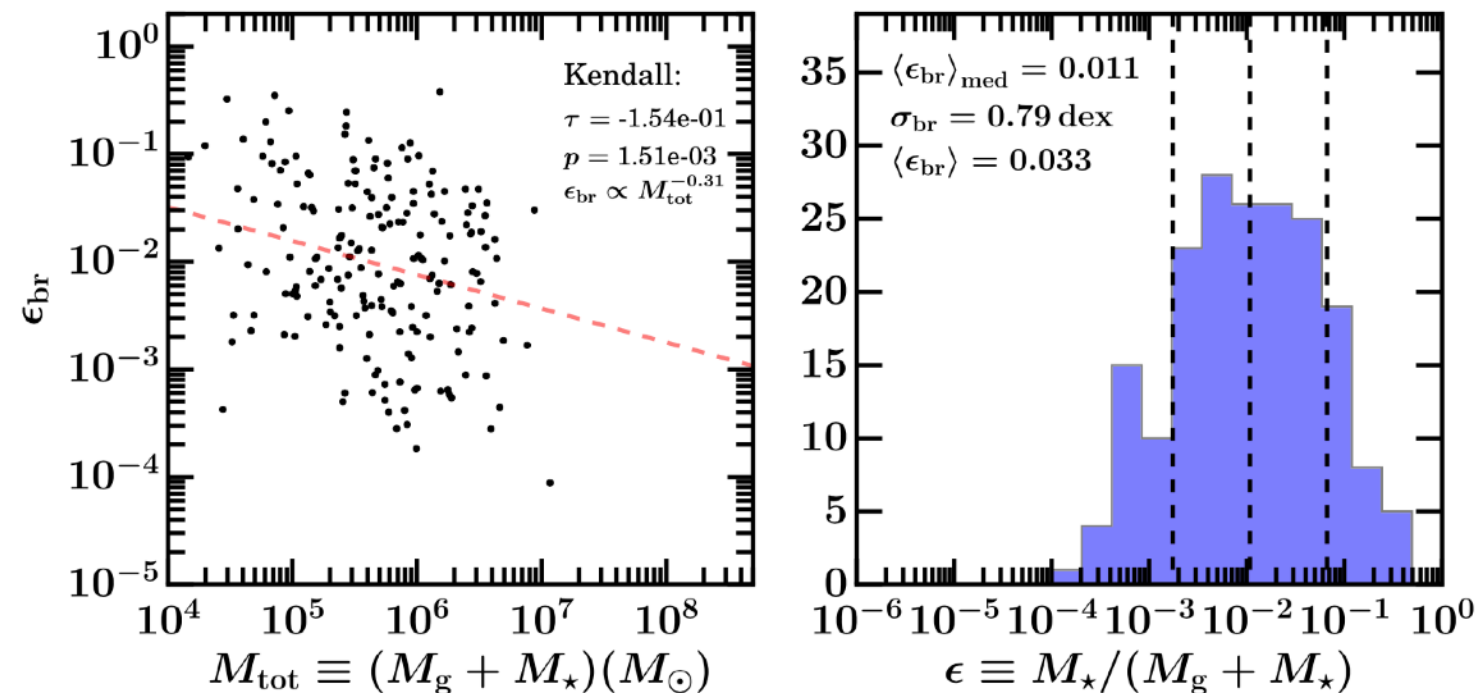
The mass accretion rate



Mooney & Solomon 1988

THE ASTROPHYSICAL JOURNAL, 833:229 (15pp), 2016 December 20

LEE, MIVILLE-DESCHÊNES, & MURRAY



Is $dM_{\star}/dt$ nearly constant?

Spherical Collapse Model

Murray & Chang (2015) ApJ 804 44

$$\frac{\partial \rho}{\partial t} + \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 u_r \rho) = 0,$$

$$\frac{\partial u_r}{\partial t} + u_r \frac{\partial u_r}{\partial r} + \frac{1}{\rho} \frac{\partial \rho v_T^2(r, t)}{\partial r} + \frac{GM(r, t)}{r^2} = 0,$$

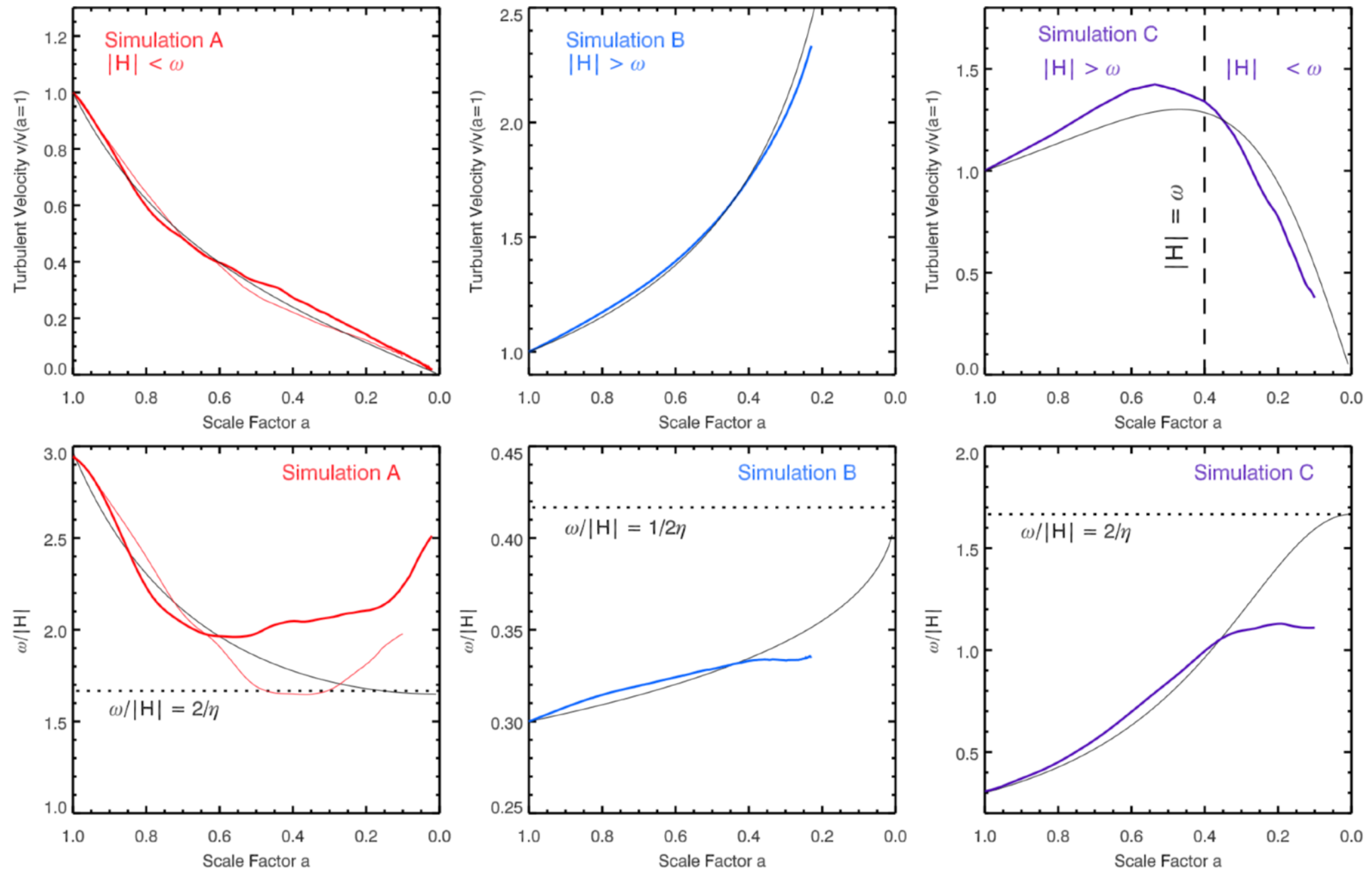
$$\frac{\partial v_T}{\partial t} + u_r \frac{\partial v_T}{\partial r} + \left(1 + \eta \frac{v_T}{u_r} \right) \frac{v_T u_r}{r} = 0.$$

$$M(r, t) = 4\pi \int r^2 \rho(r, t) dr.$$

Robertson & Goldreich

THE ASTROPHYSICAL JOURNAL LETTERS, 750:L31 (5pp), 2012 May 10

ROBERTSON & GOLDRICH



Spherical Collapse Model

Feedback Loop

Murray & Chang (2015) ApJ 804 44

$$\frac{\partial u_r}{\partial t} + u_r \frac{\partial u_r}{\partial r} = - \frac{1}{\rho} \frac{\partial \rho v_T^2}{\partial r} - \frac{GM(r, t)}{r^2},$$

$$\frac{\partial v_T}{\partial t} + u_r \frac{\partial v_T}{\partial r} = + \frac{v_T |u_r|}{r} - \eta \frac{v_T^2}{r}.$$

As a result :

$$v_T(r, t) \approx |u_r(r, t)|$$

Spherical Collapse Model: Stellar sphere of influence

- Star or Star cluster sphere of influence: breaks simple self-similarity

$$M_g(r_*(t), t) \equiv 4\pi \int_{0^+}^{r_*(t)} r^2 \rho(r, t) dr = M_*(t)$$

- Solutions for $r_{\text{disk}} < r < r_*(t)$ differ from those for $r_*(t) < r < R$

Spherical Collapse Model

Murray & Chang (2015) ApJ 804 44

$$\frac{\partial u_r}{\partial t} + u_r \frac{\partial u_r}{\partial r} + \frac{1}{\rho} \frac{\partial \rho v_T^2}{\partial r} + \frac{GM_*(t)}{r^2} = 0,$$

since

$$v_T \sim |u_r|$$

$$|u_r| \sim r^{-1/2}$$

for $r \ll r_*$.

Spherical Collapse Model: the density attractor.

$$\rho(r, t) \sim r^{-k_\rho} \tilde{\rho}(t)$$

$$u_r(r, t) \sim r^p \tilde{u}(t)$$

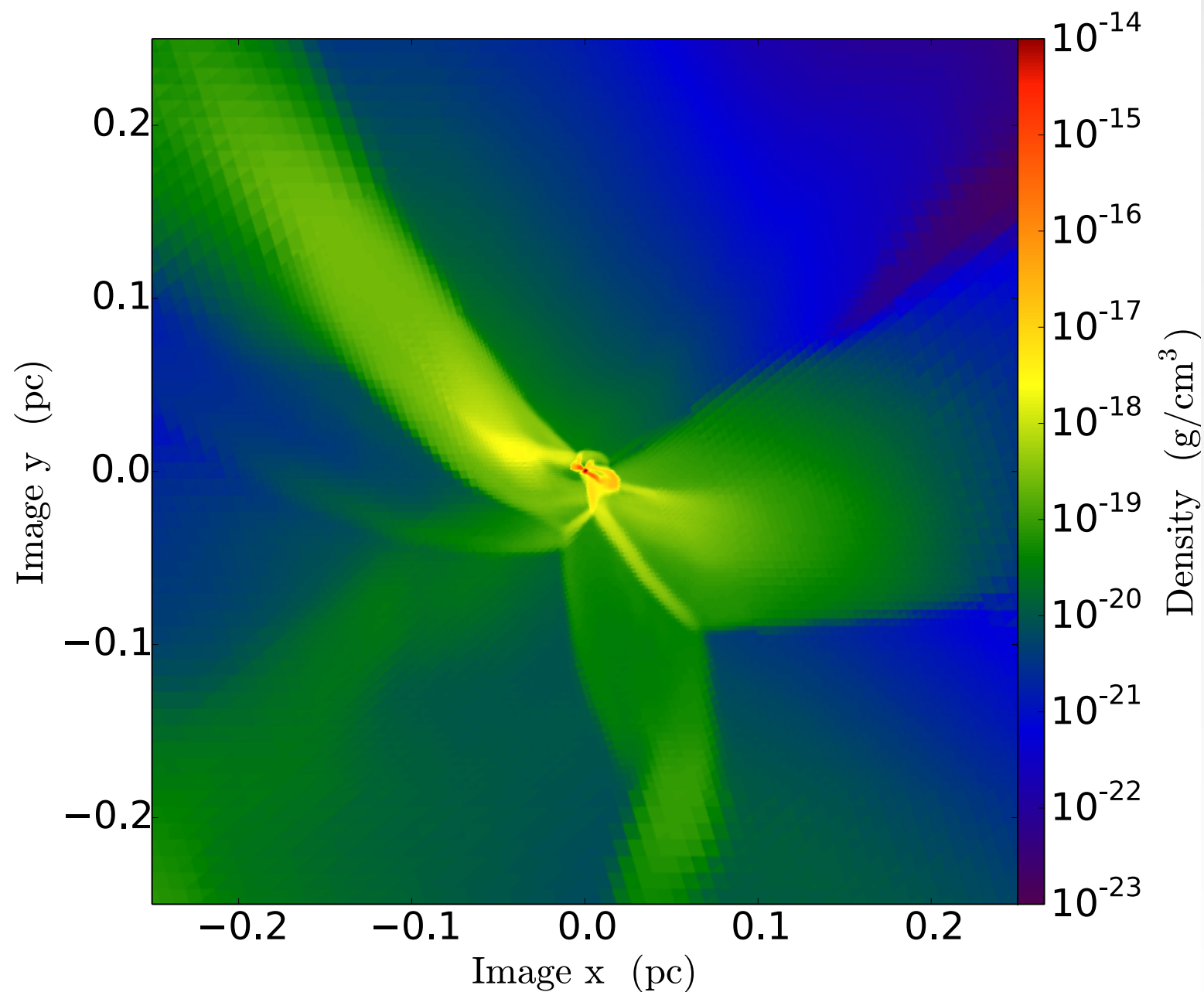
$$\frac{\partial \tilde{\rho}(t)}{\partial t} r^{-k_\rho} + \tilde{\rho}(t) \tilde{u}(t) r^{p-k_\rho-1} [2 + p - k_\rho] = 0,$$

since $p < 0$, at small r the second term grows much faster than the first, so

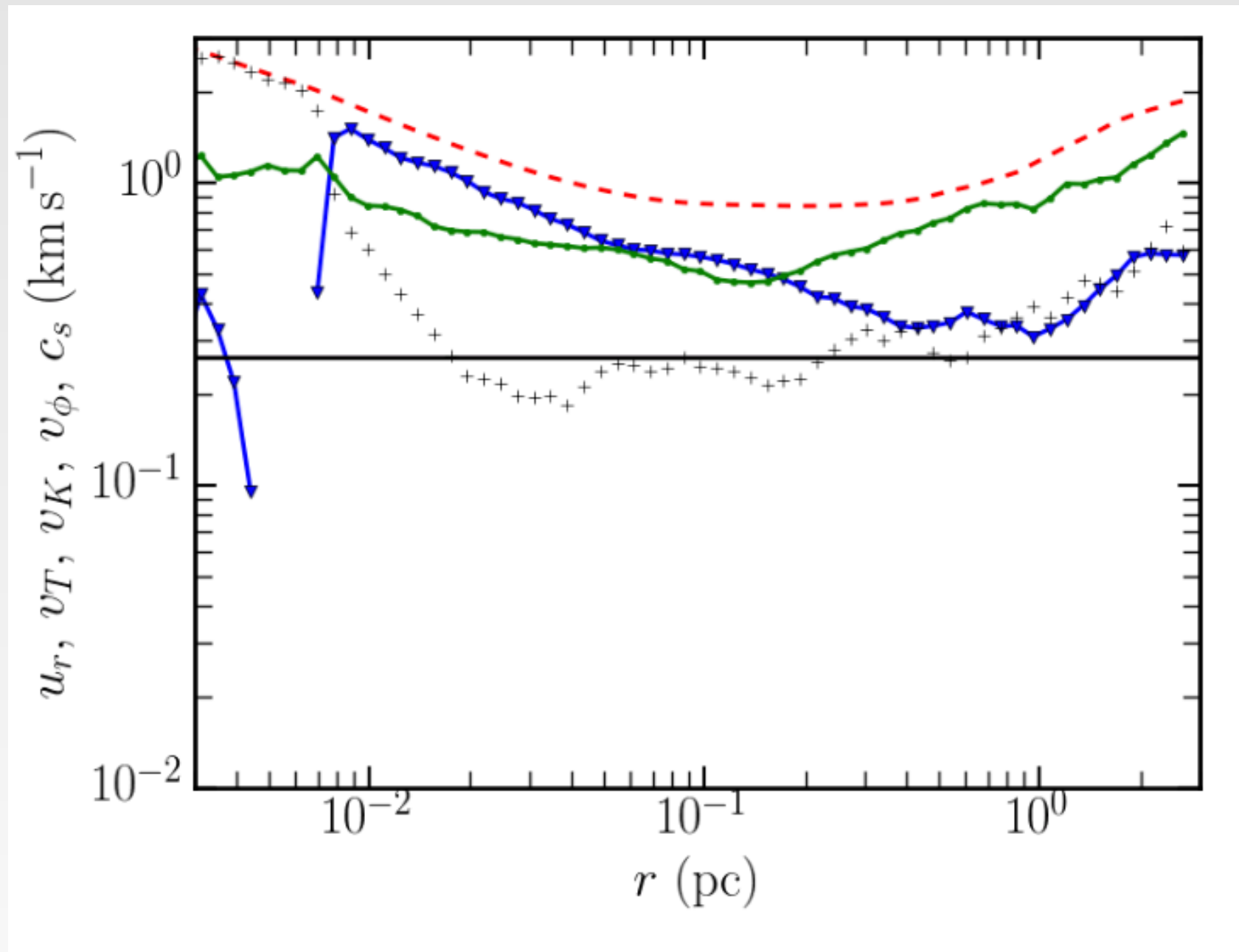
$$k_\rho = 2 + p, \text{ and}$$

$$\frac{\partial \tilde{\rho}(t)}{\partial t} = 0, \text{ i.e., } \rho(r, t) \rightarrow \rho(r).$$

Simulations of Turbulent Collapse



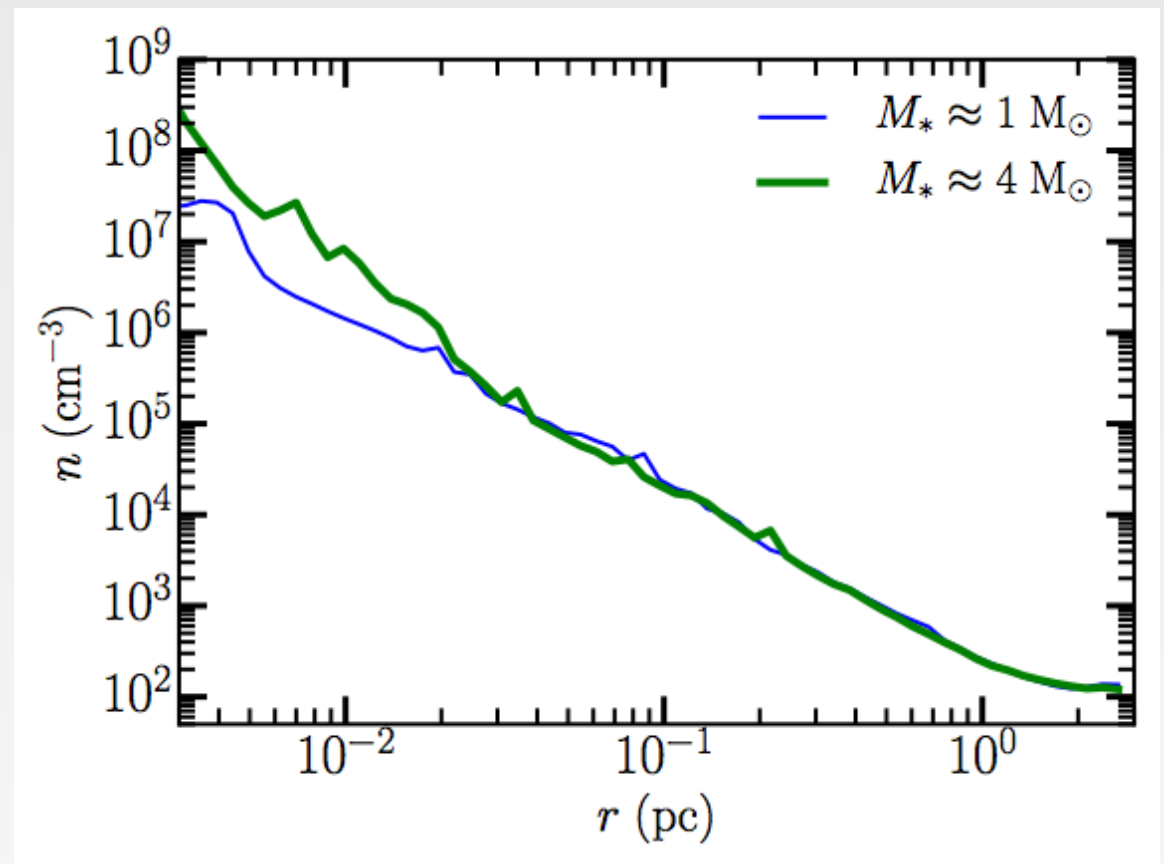
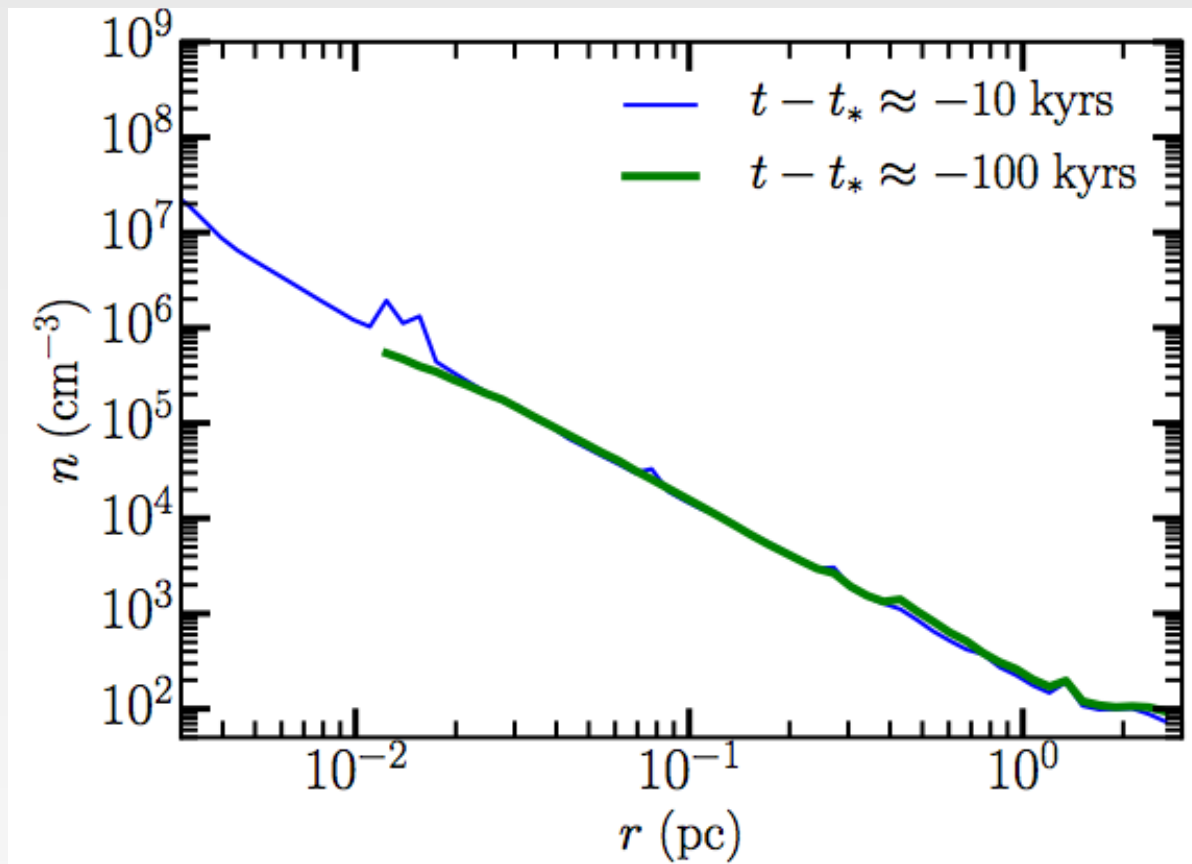
Infall, turbulent, & rotational velocity



Supersonic infall, at $\sim 1/3$ - $1/2 v_K$ (blue line); not HSE
 v_T (green line) increases inward for $r > r_*$, increases inward for
 $r < r_*$

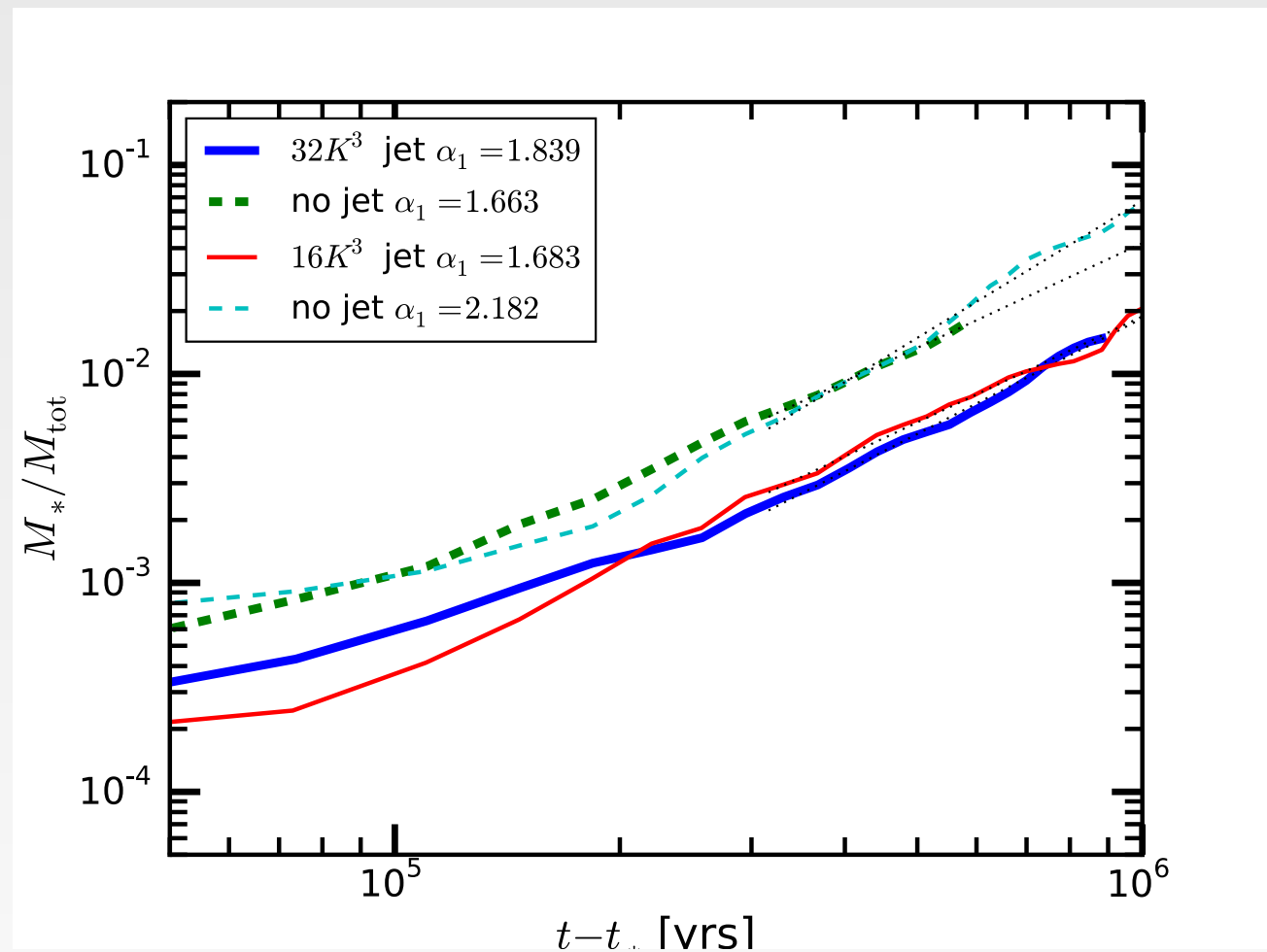
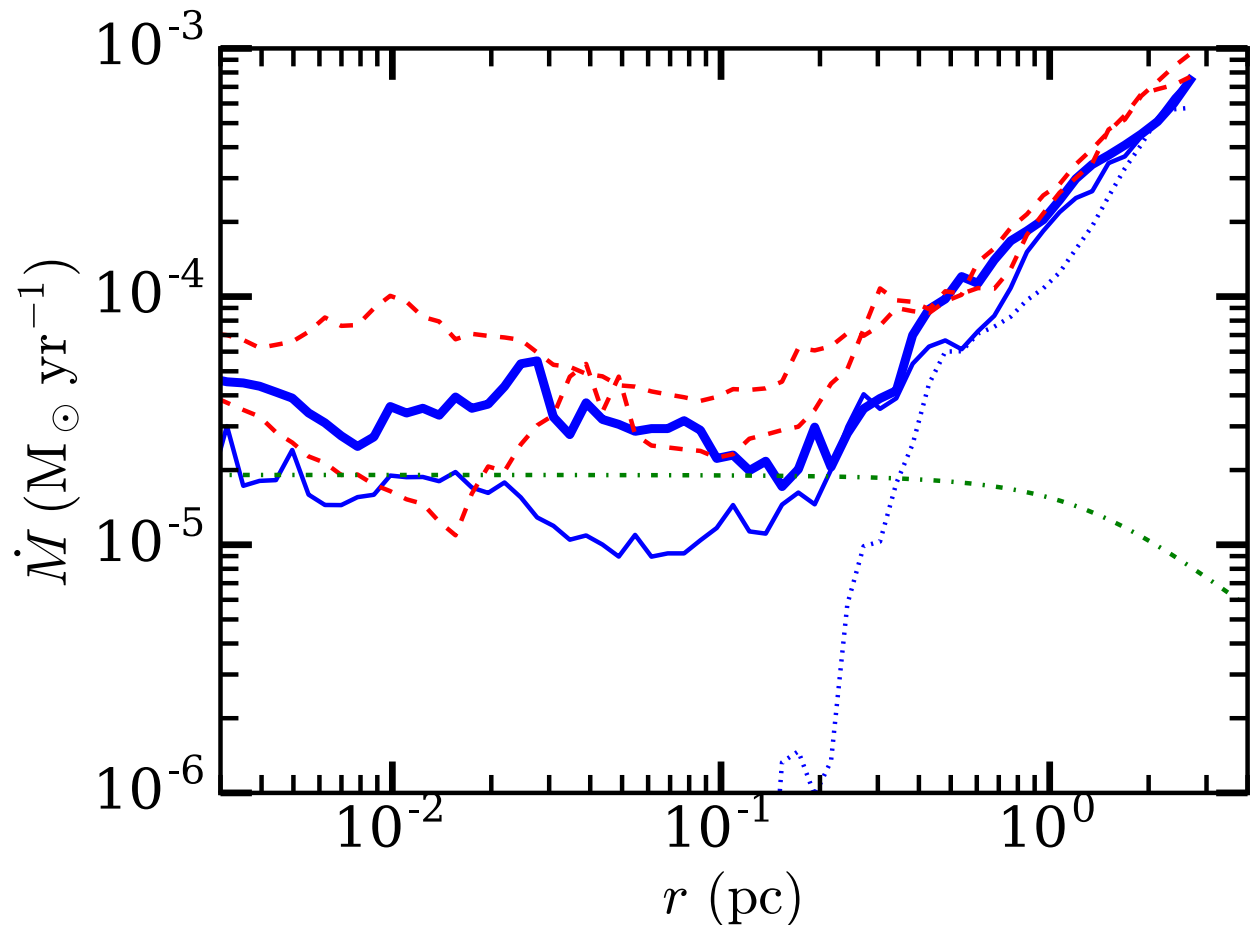
This is when the star has $M_* \approx 3M_\odot$

The density attractor



The density does not vary with time for at least $1/3$ of a large scale dynamical time for $r > r_d$

The mass accretion rate



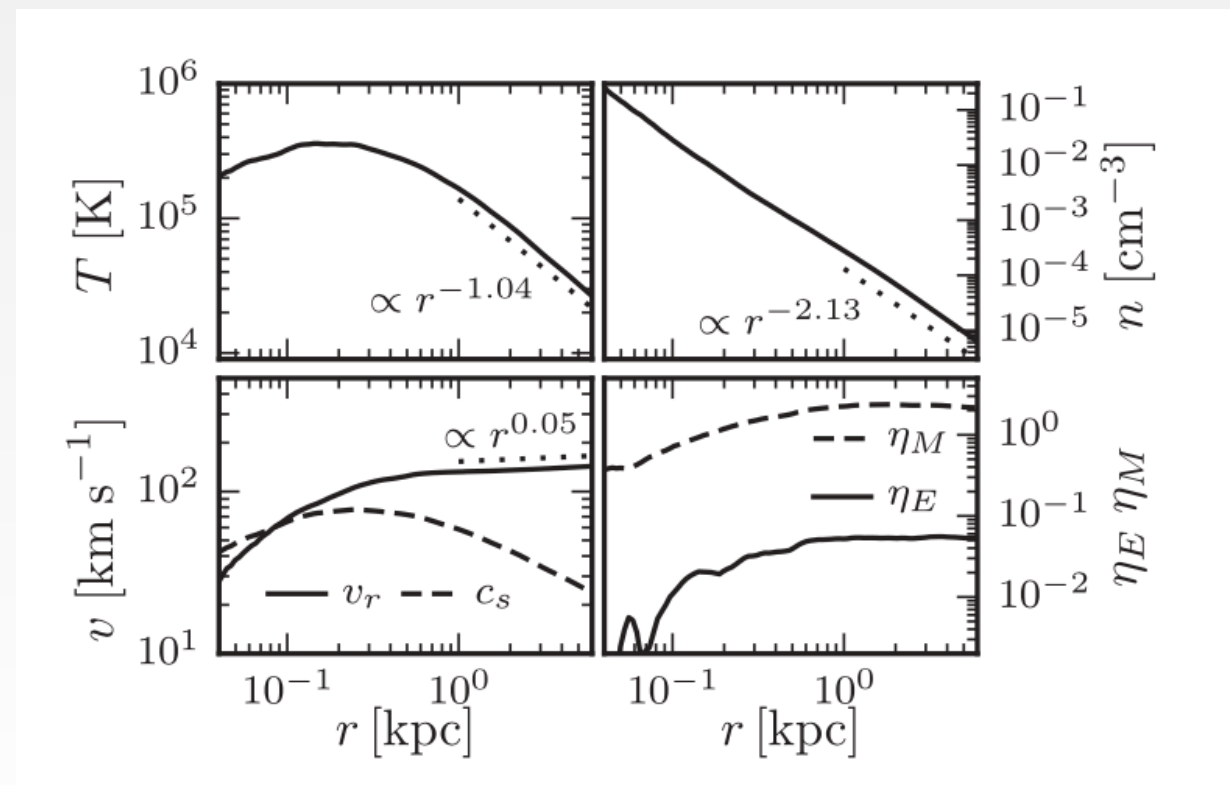
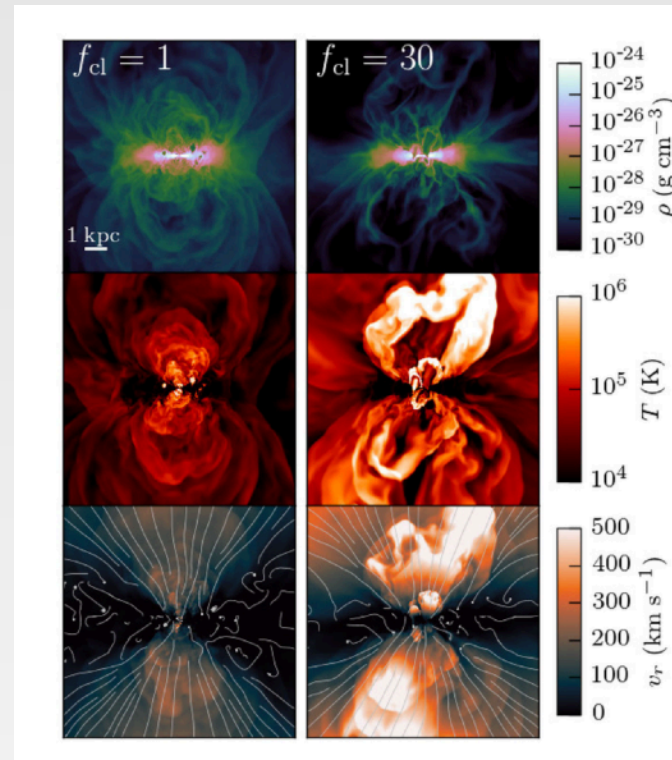
$dM/dt(r)$ is nearly constant for $r < r_*$

$$M_*(t) \sim t^{1.8}$$

Stellar Feedback in action

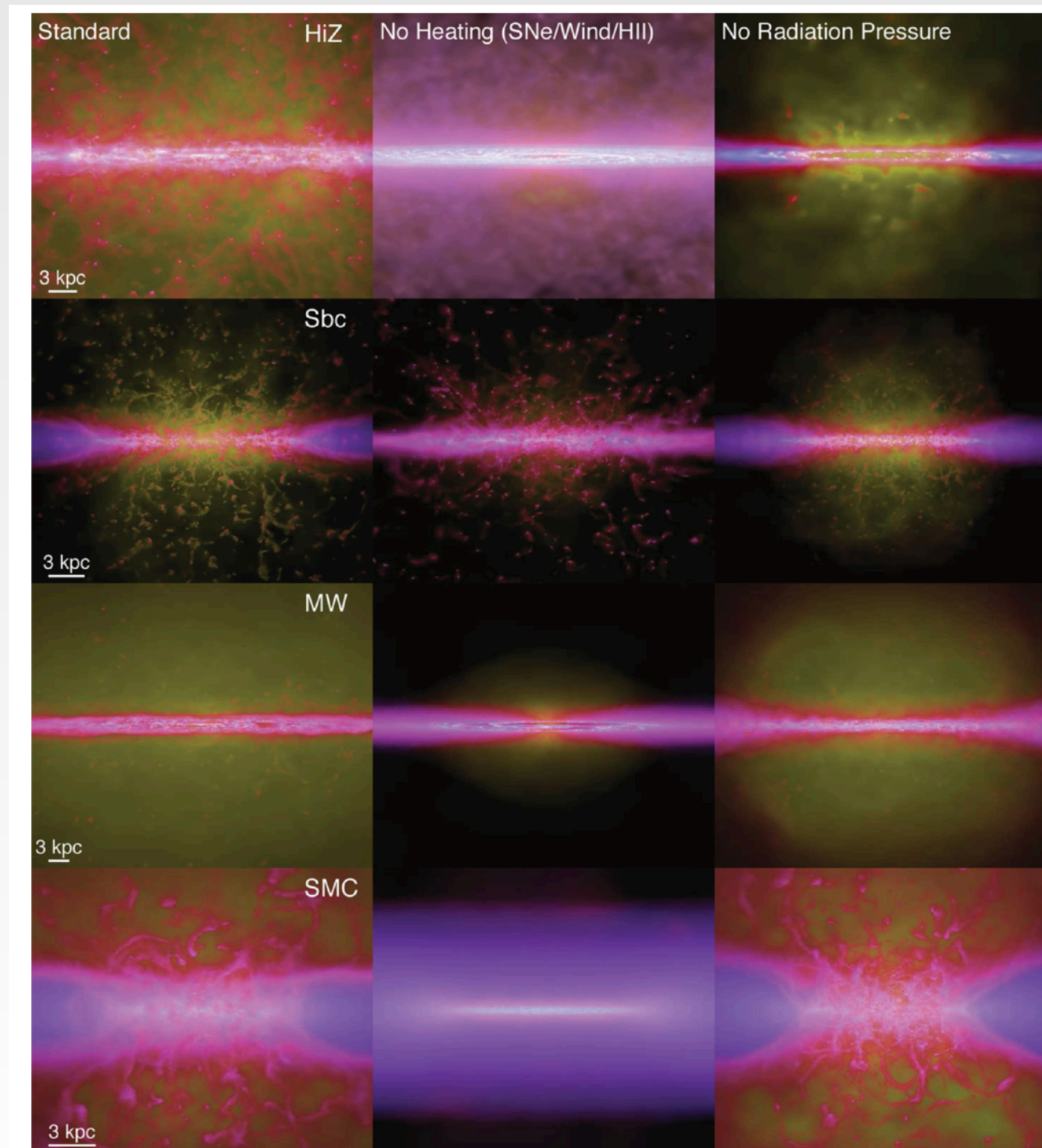


Moving to Larger Scales: 1KPC disk, resolved Sedov-Taylor phase



Moving to Larger Scales: isolated galaxies, mostly resolved Sedov-Taylor phase

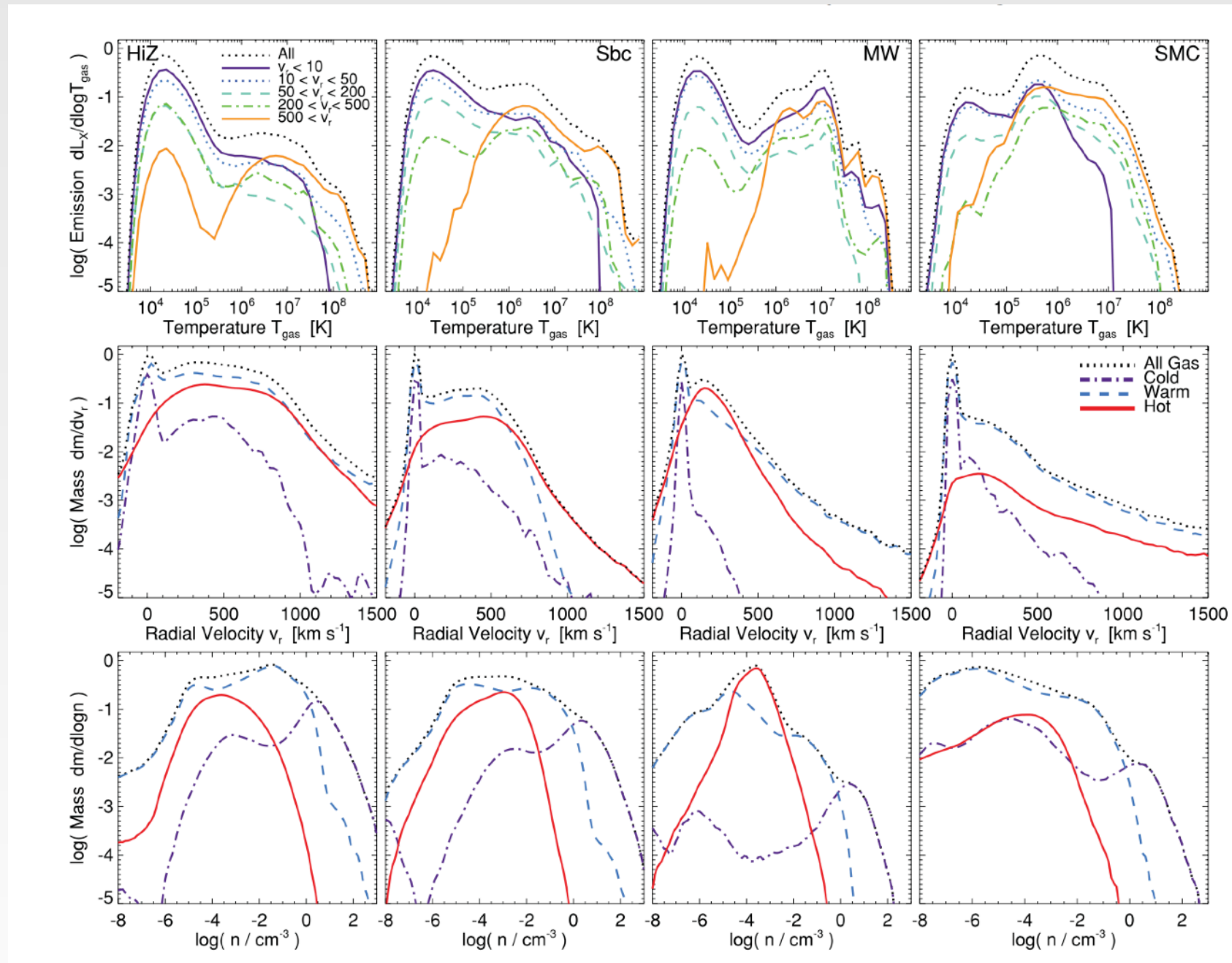
resolution 3pc
100-1000M_⊙



Hopkins, Quataert, & Murray (2012)

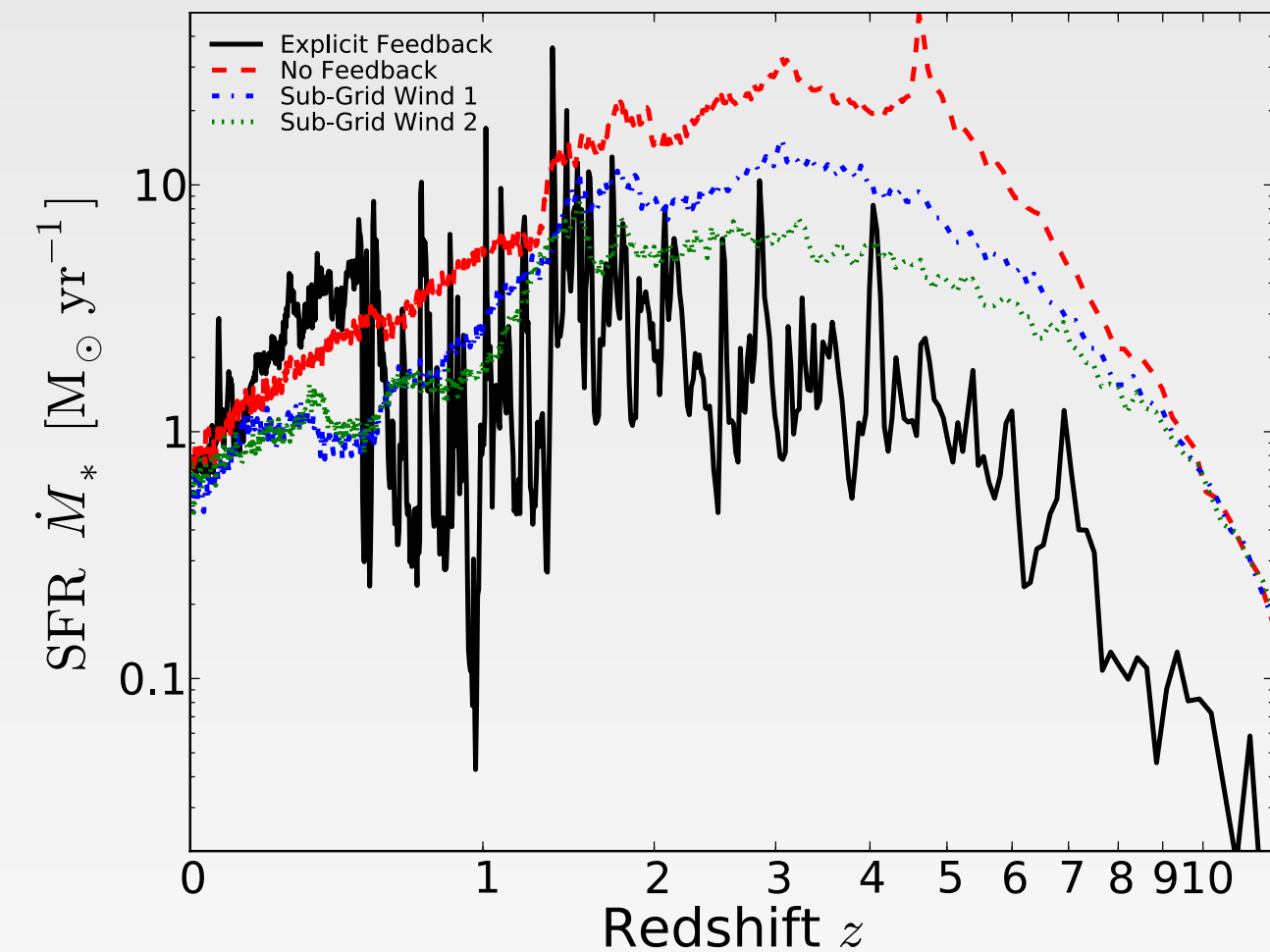
Moving to Larger Scales: isolated galaxies, mostly resolved Sedov-Taylor phase

resolution 3pc
100-1000M_☉



Hopkins, Quataert, & Murray (2012)

Moving to Larger Scales: cosmological zoom galaxies (FIRE)



8

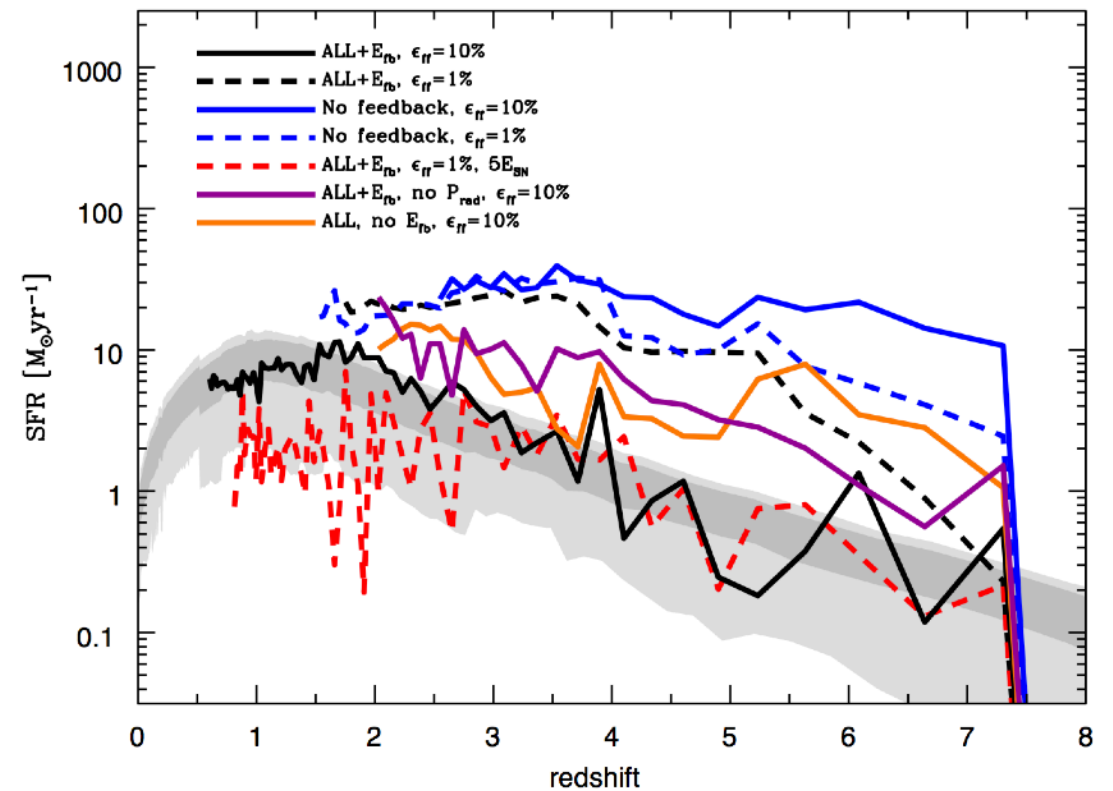
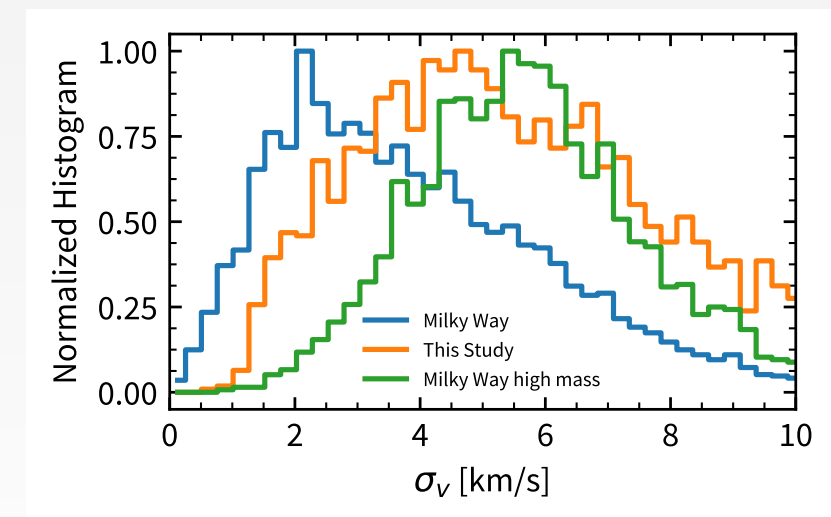
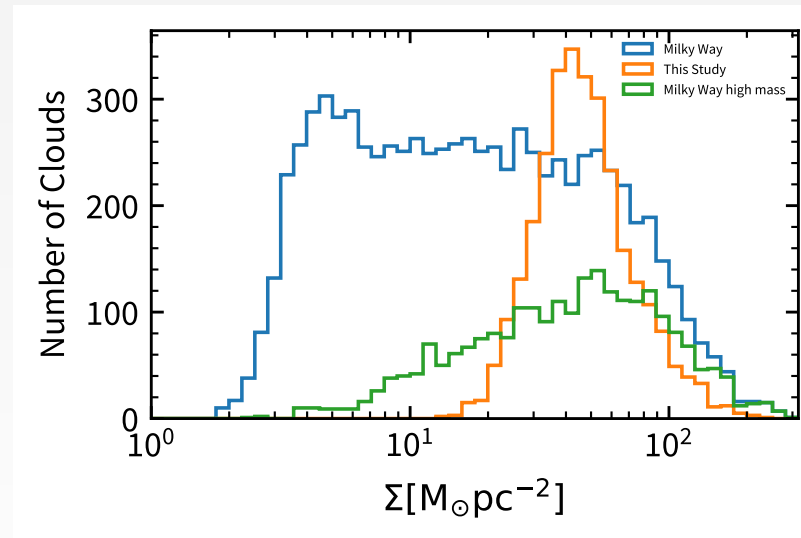
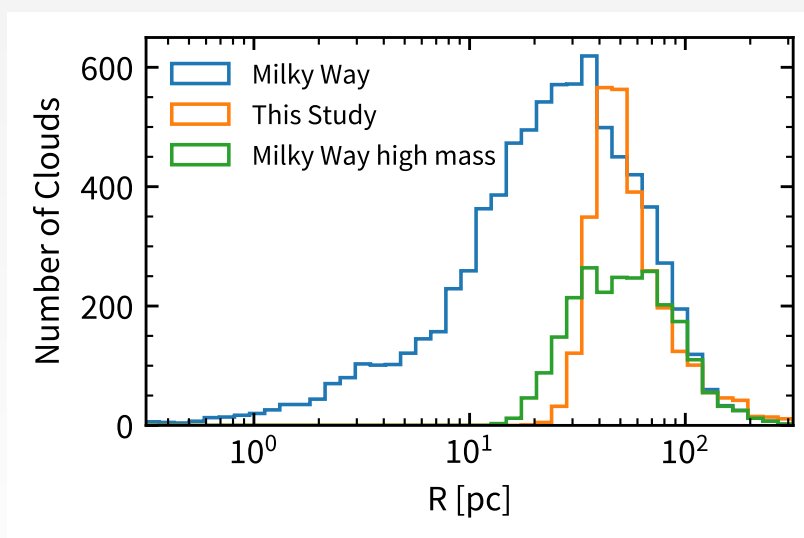
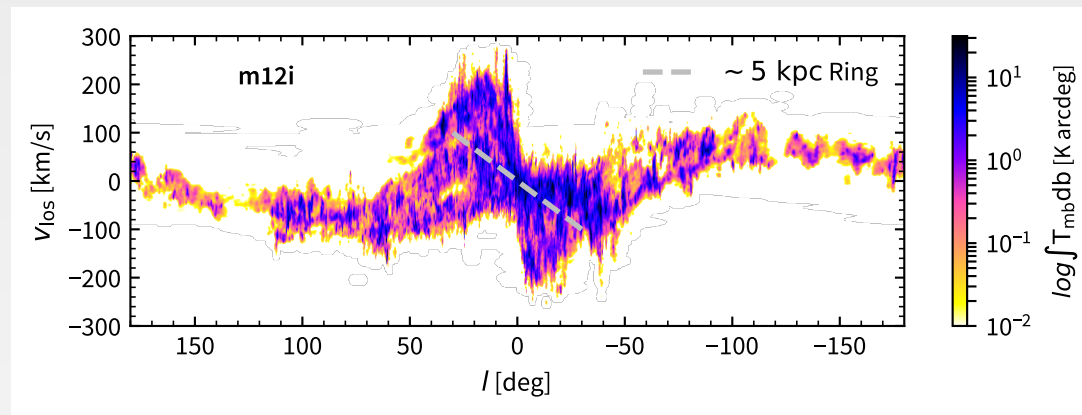
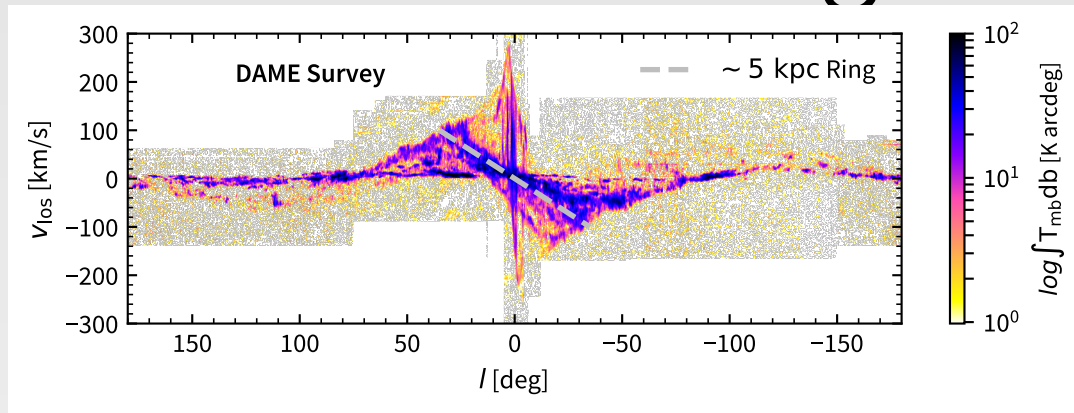


FIG. 4.— Simulated star formation histories compared to the Behroozi et al. (2013) data for $M_{\text{vir}}(z=0) = 10^{12} M_{\odot}$. Dark and light gray shaded areas are one- and two-sigma confidence regions respectively. We adopt bins of size $\Delta t_{\text{SF}} = 100 \text{ Myr}$ for the simulated SFHs. Without feedback, SFHs are overpredicted by at least one order of magnitude at $z > 1$. Efficient feedback in conjunction with $\epsilon_{\text{ff}} \gtrsim 10\%$ (ALL_Efb_e010) renders a star formation history in agreement with the Behroozi et al. data. In simulations with low star formation efficiency of $\epsilon_{\text{ff}} = 1\%$, the effectiveness of feedback diminishes and SFHs are ~ 1 dex higher than expected. Boosting the available SNe feedback (ALL_Efb_e001_5ESN) alleviates this, but leads to a significantly stronger suppression of star formation at $z \lesssim 2.5$. Both radiation pressure and efficient SN feedback appear crucial, as removing any of these feedback sources offsets the SFH by up to ~ 1 dex, as discussed in the text.

Mostly not resolved S-T phase

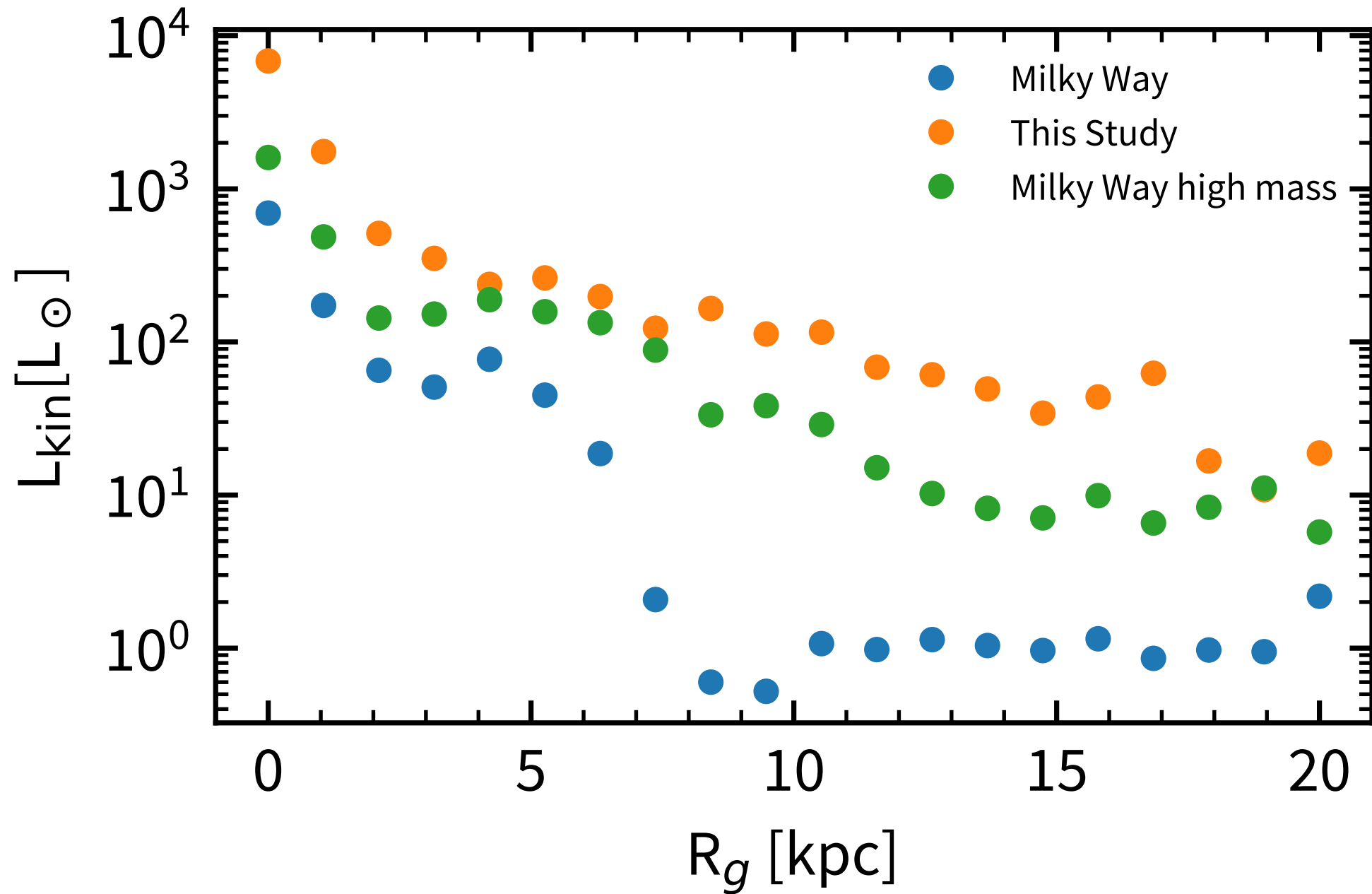
Hopkins et al. (2014); Agertz & Kravtsov (2016)

Moving to Larger Scales: cosmological zoom galaxies (FIRE)



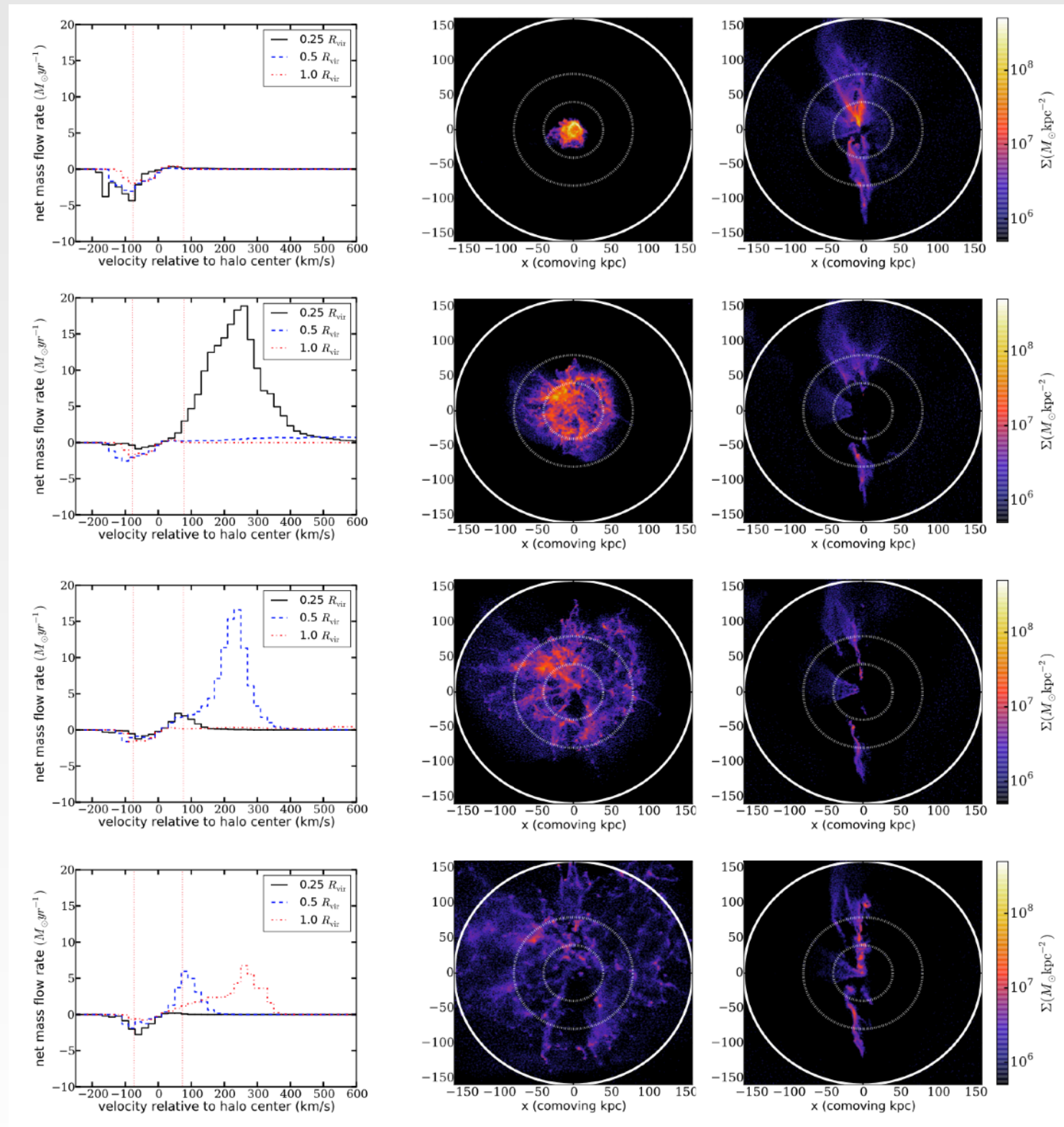
Lakhlani et al. (2018)

Moving to Larger Scales: cosmological zoom galaxies (FIRE)



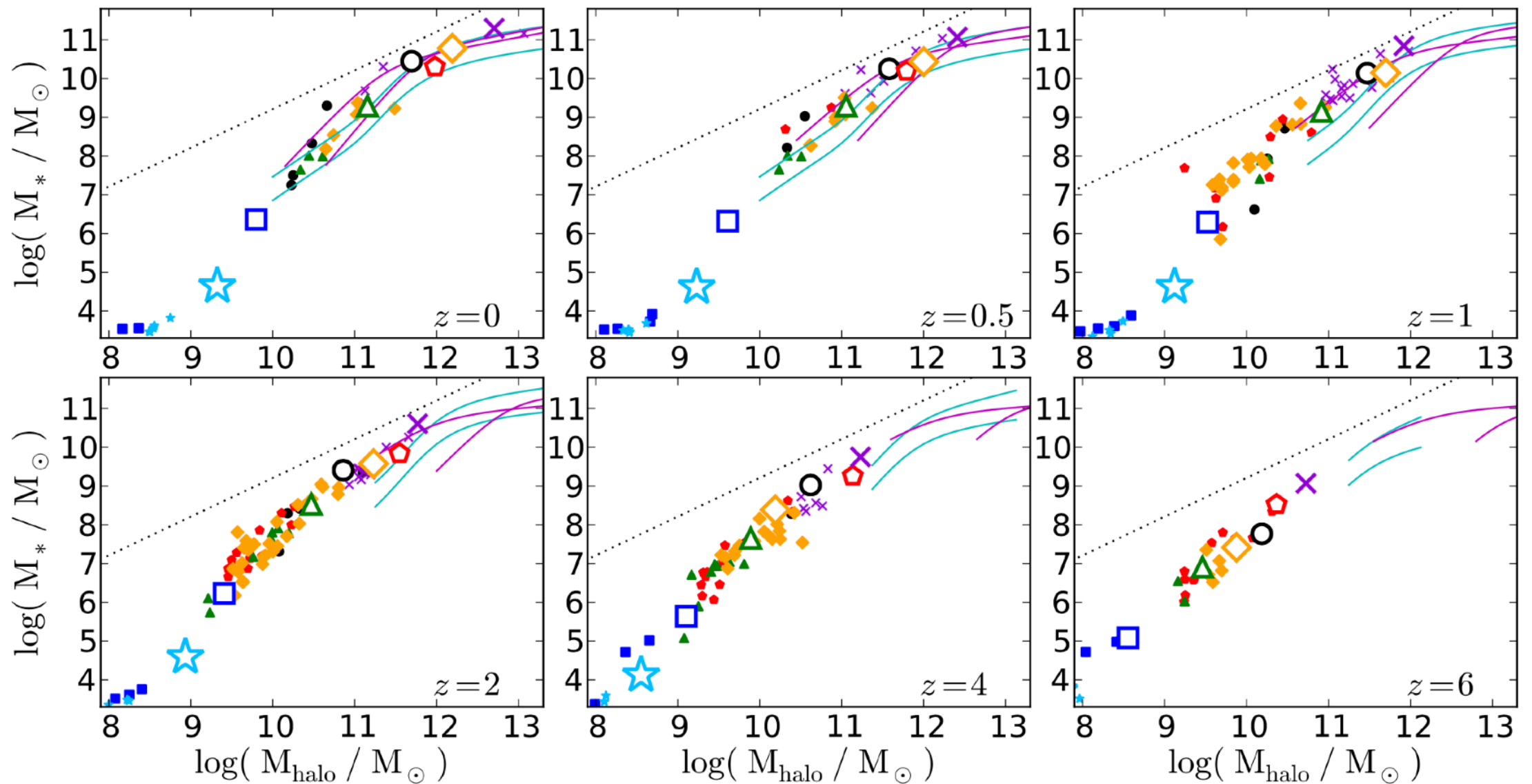
Lakhlani et al. (2018)

Moving to Larger Scales: cosmological zoom galaxies (FIRE)



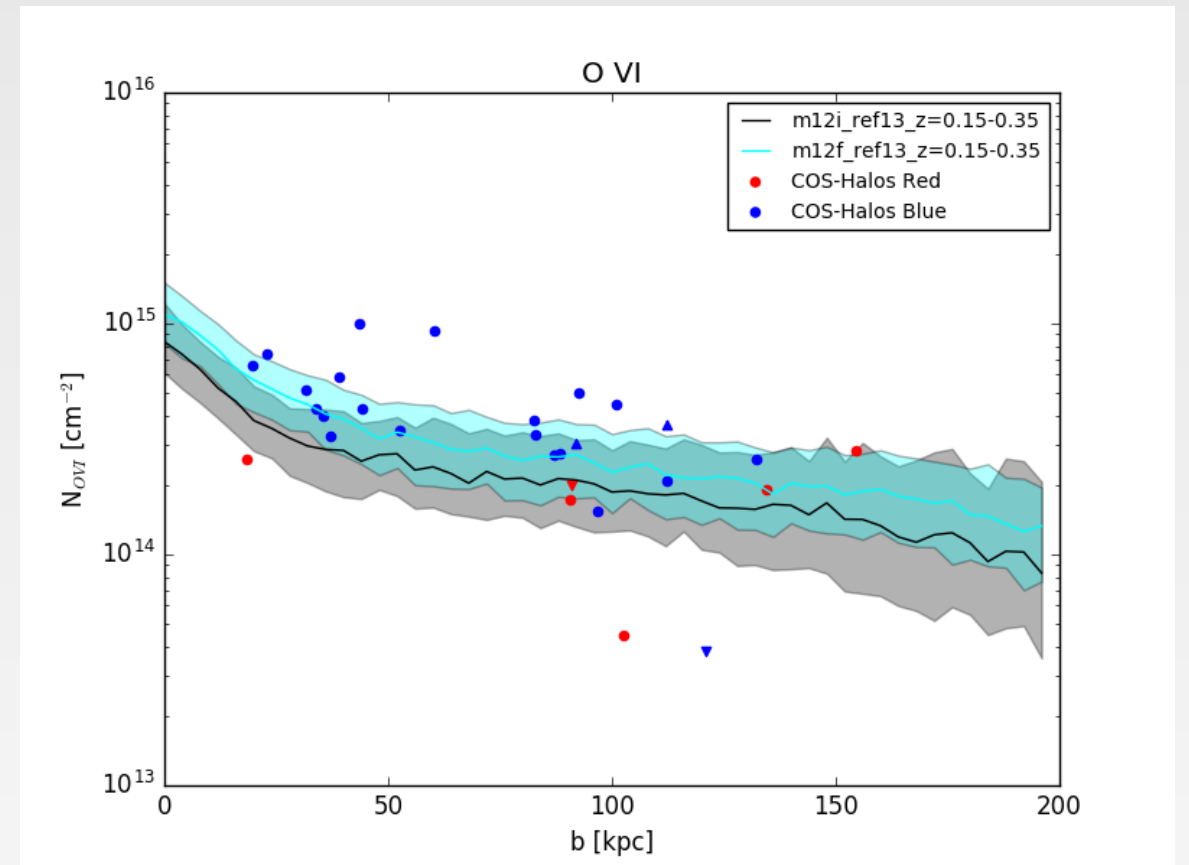
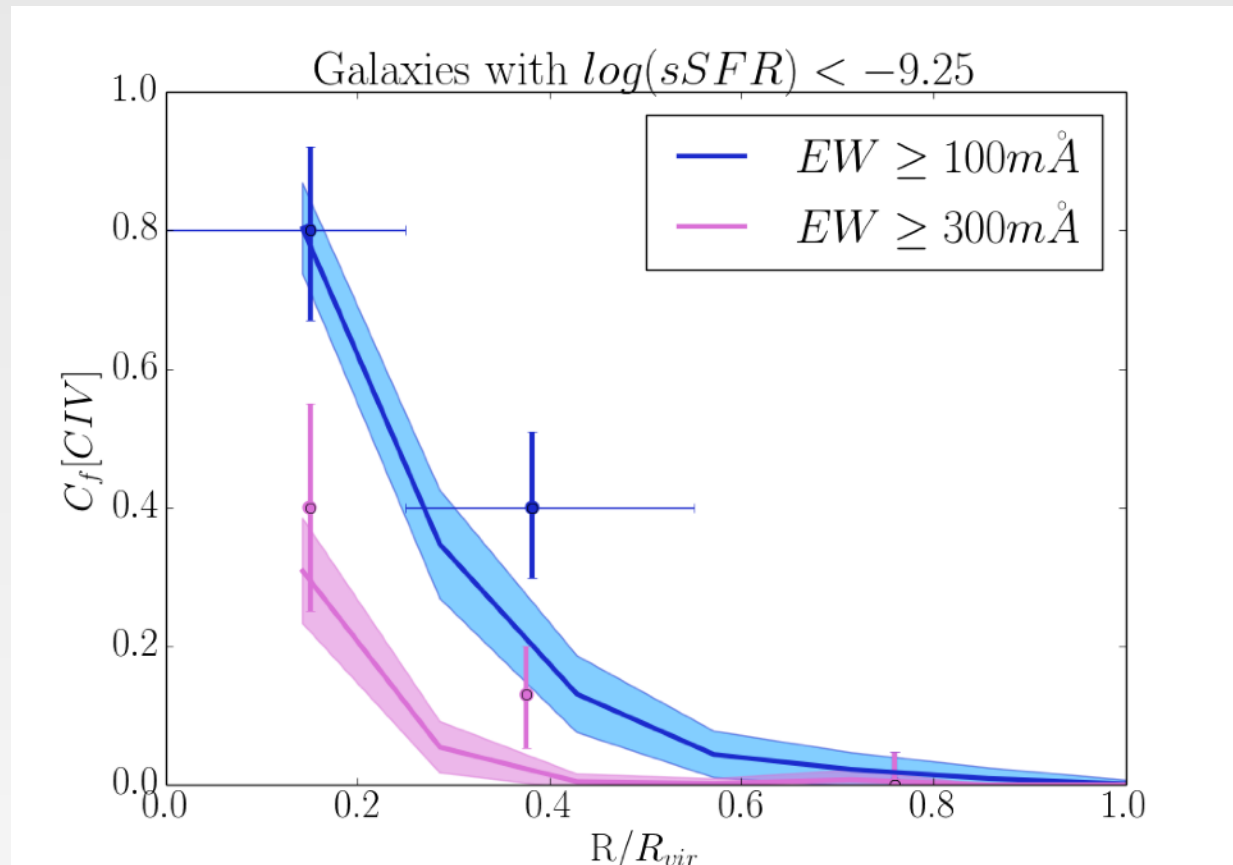
Muratov et al. (2015) MNRAS **454** 2691

Moving to Larger Scales: cosmological zoom galaxies (FIRE)



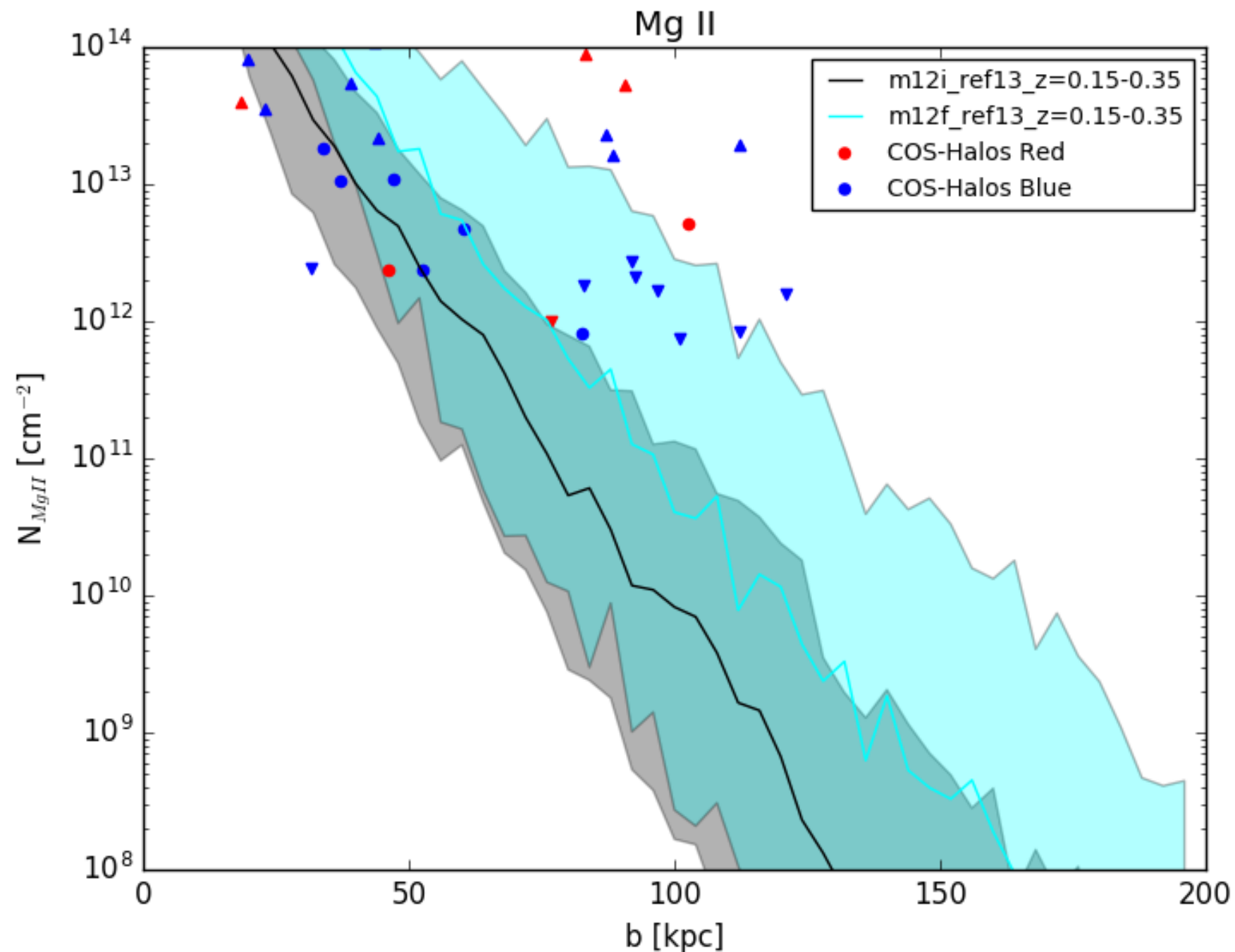
Hopkins et al. (2014)

Moving to Larger Scales: cosmological zoom galaxies (FIRE)



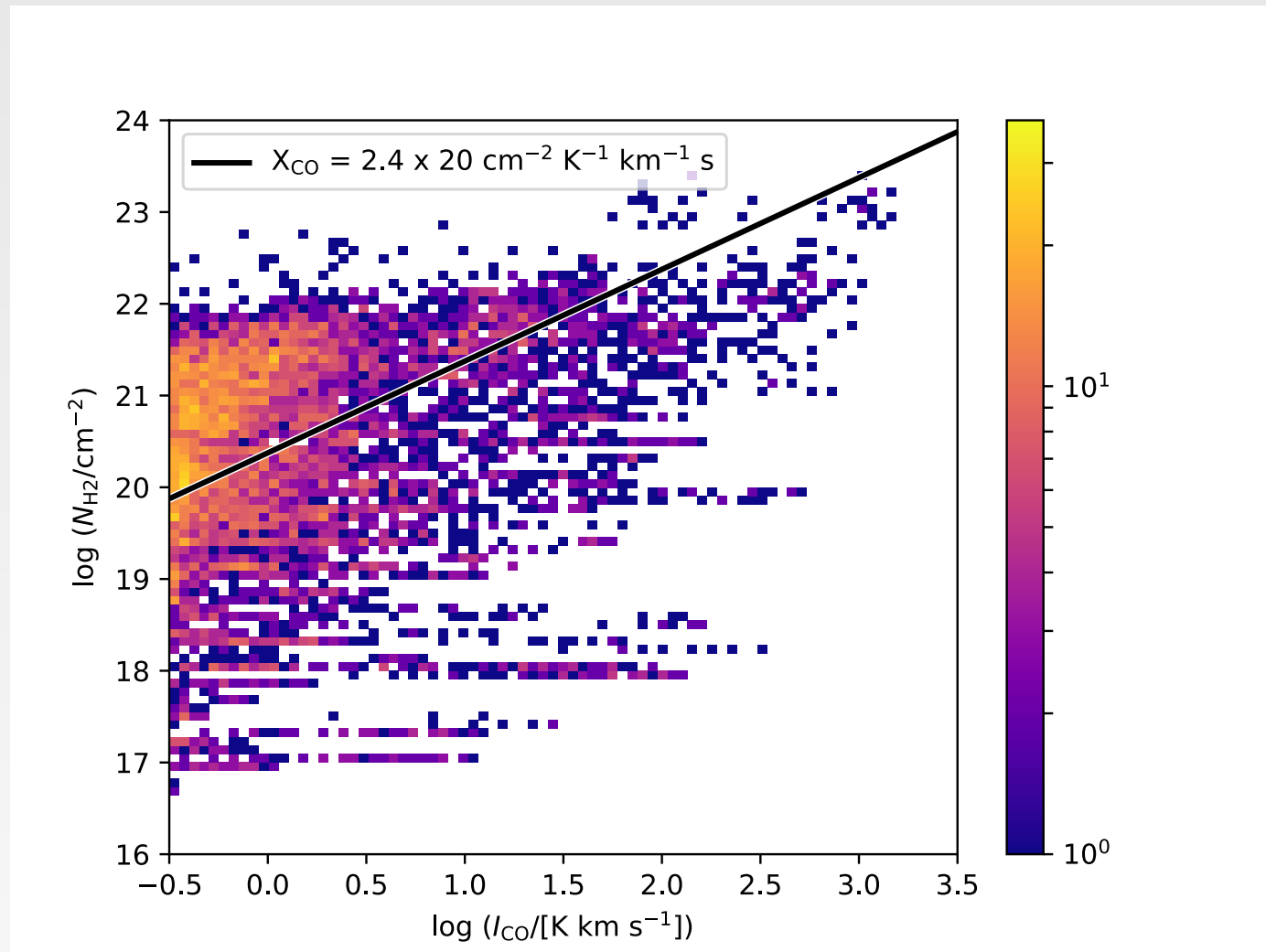
Fei Li, NWM, et al; Hummels, NWM, et al. in prep

Moving to Larger Scales: cosmological zoom galaxies (FIRE)

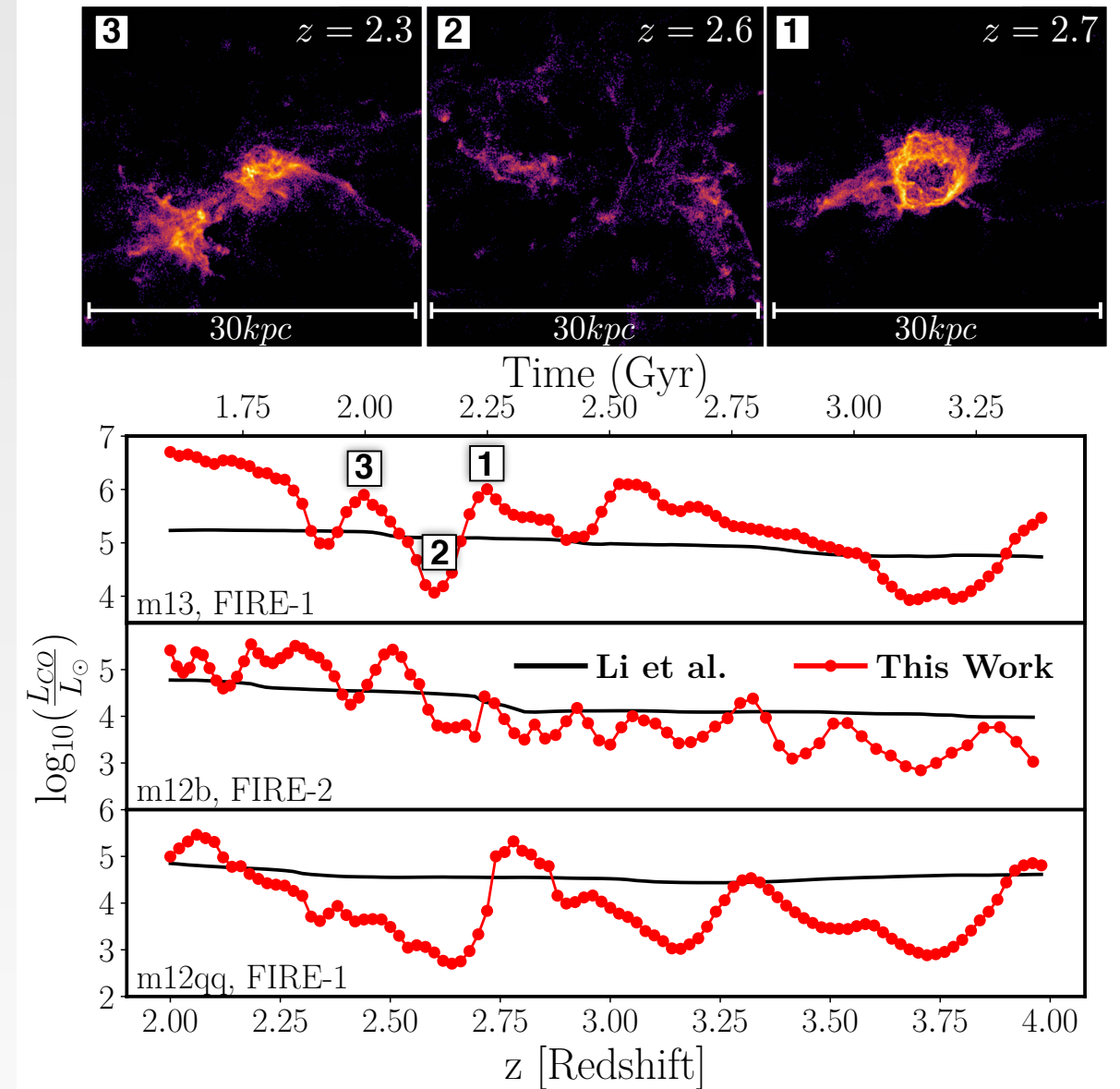


Hummels et al. in prep

CO Luminosity vs z



Laura Keating, Alex Richings



Lakhlani et al (2018b)

L_{CO} in a FIRE simulation

Conclusions

- A new, complete GMC catalog is now available
- Turbulence in inner galaxy GMCs driven by contraction; most GMCs are not forming stars
- These two results suggest that feedback must disrupt the cloud rather than slowing the collapse
- SN feedback drives winds in simulations with a range of resolution; need to resolve S-T, or sub-grid model
- The resulting galaxies have realistic ISMs
- The resulting galaxies lie on the $M_{\text{star}}-M_{\text{halo}}$ relation
- The CGM is not resolved; despite that, CIV, OVI looks ok; MgII does not
- The IGM may be expelled entirely—this can be checked by intensity mapping.