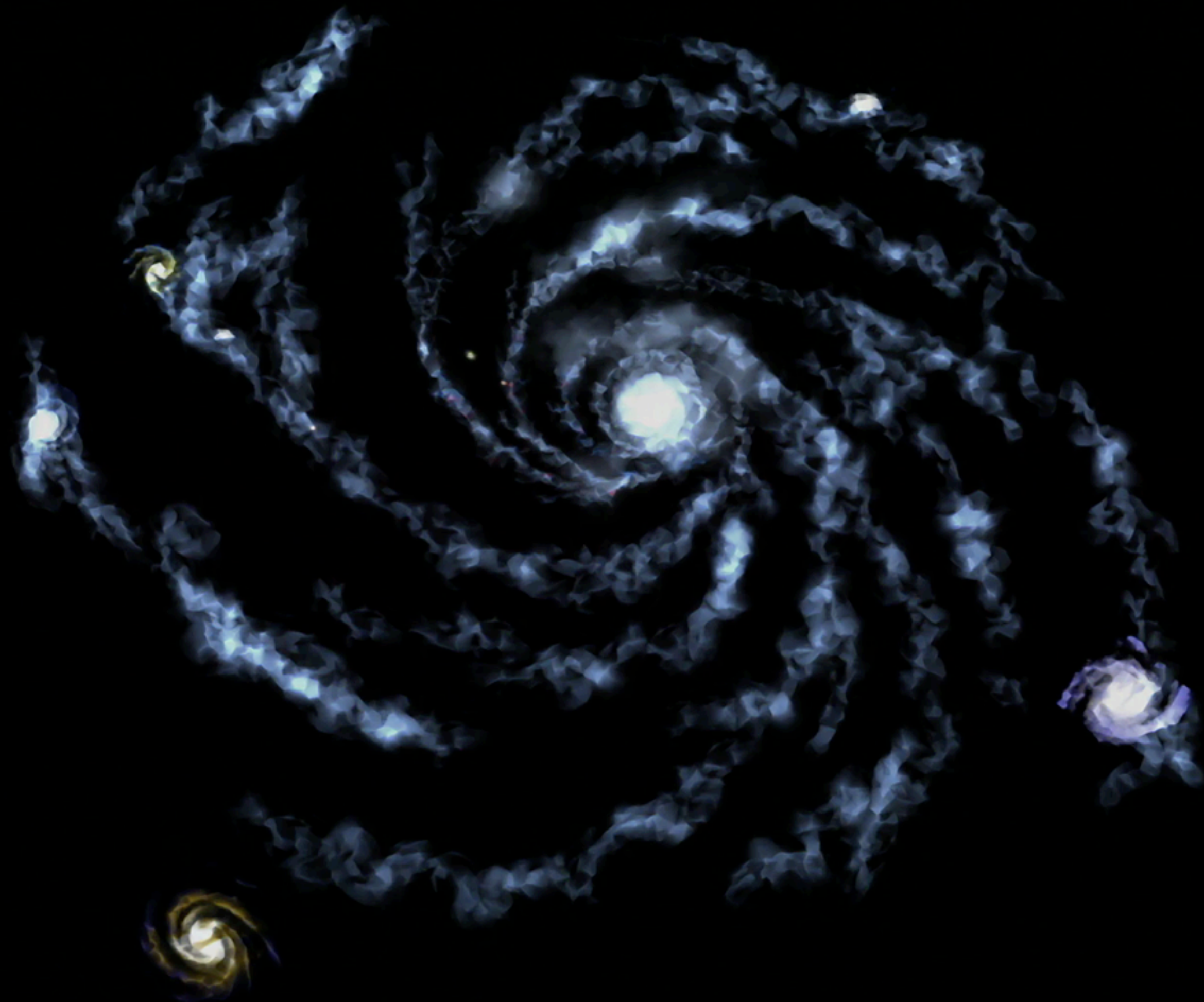
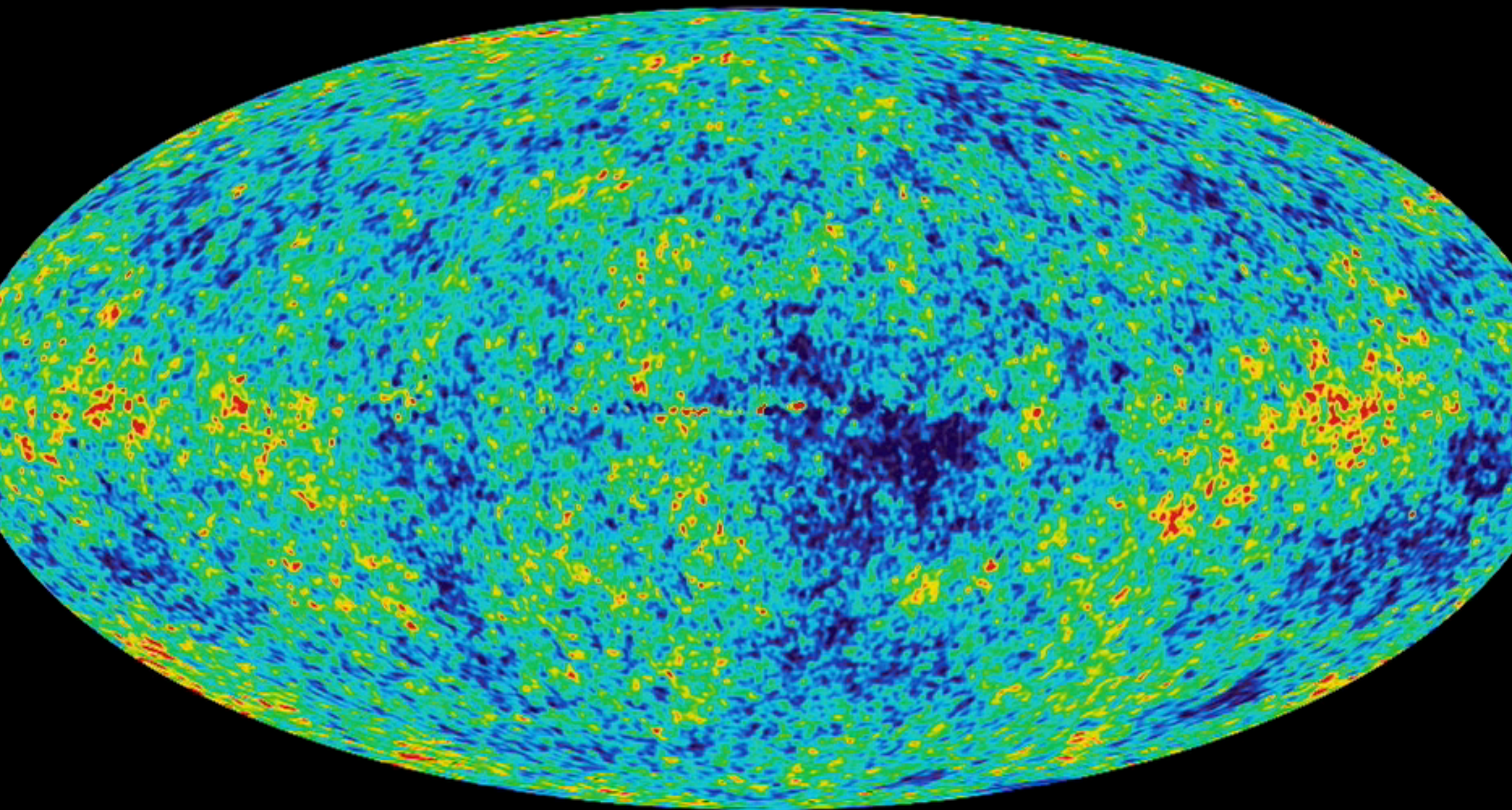


From Initial Conditions to Structure Formation, and Back

Eiichiro Komatsu (Max Planck Institute for Astrophysics)
COSMO-16, Univ. of Michigan, August 8, 2016





Fluctuations in photons: how about the matter distribution?

A Remarkable Story

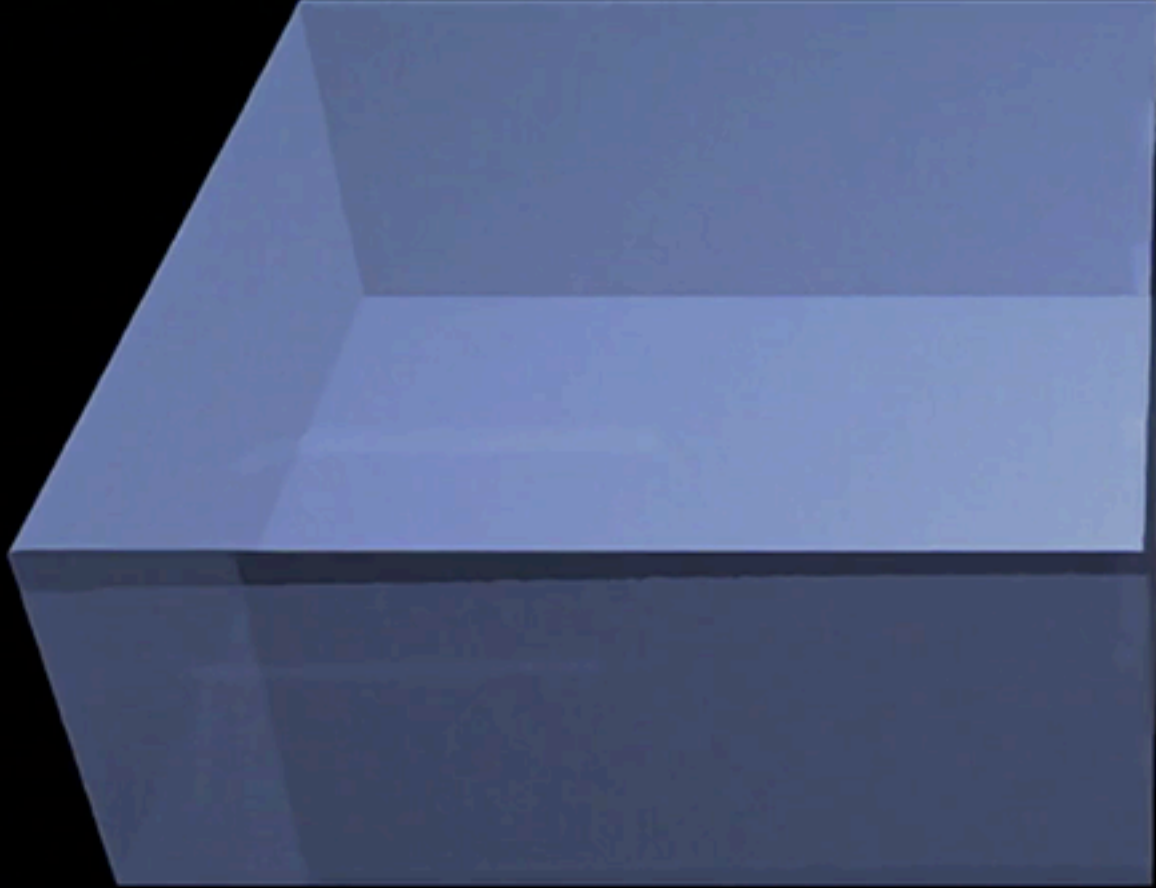
- Observations of CMB taught us that **galaxies, stars, planets, and ourselves originated from tiny fluctuations in the early Universe**

But, how confident are we?



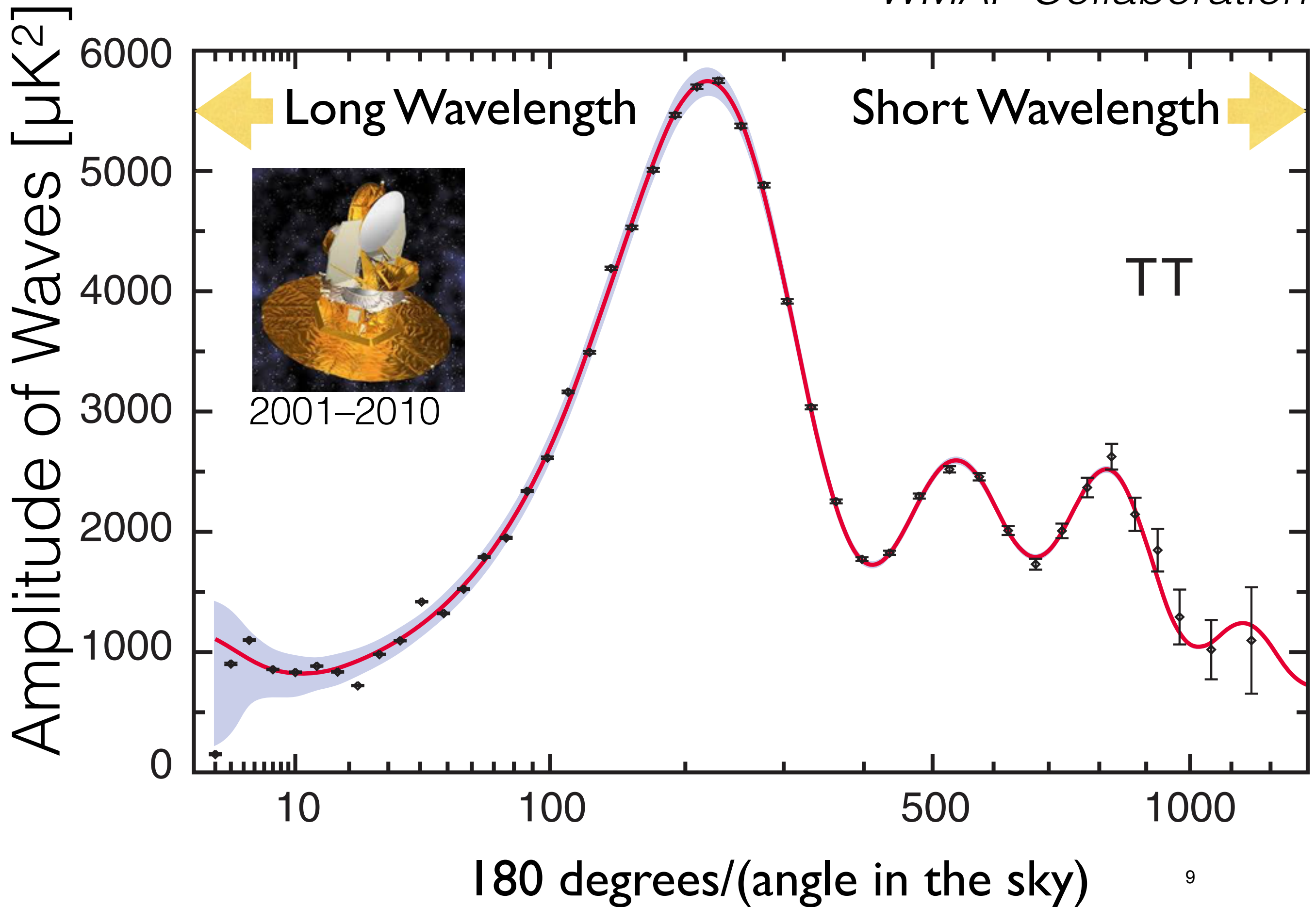
Cosmic Miso Soup

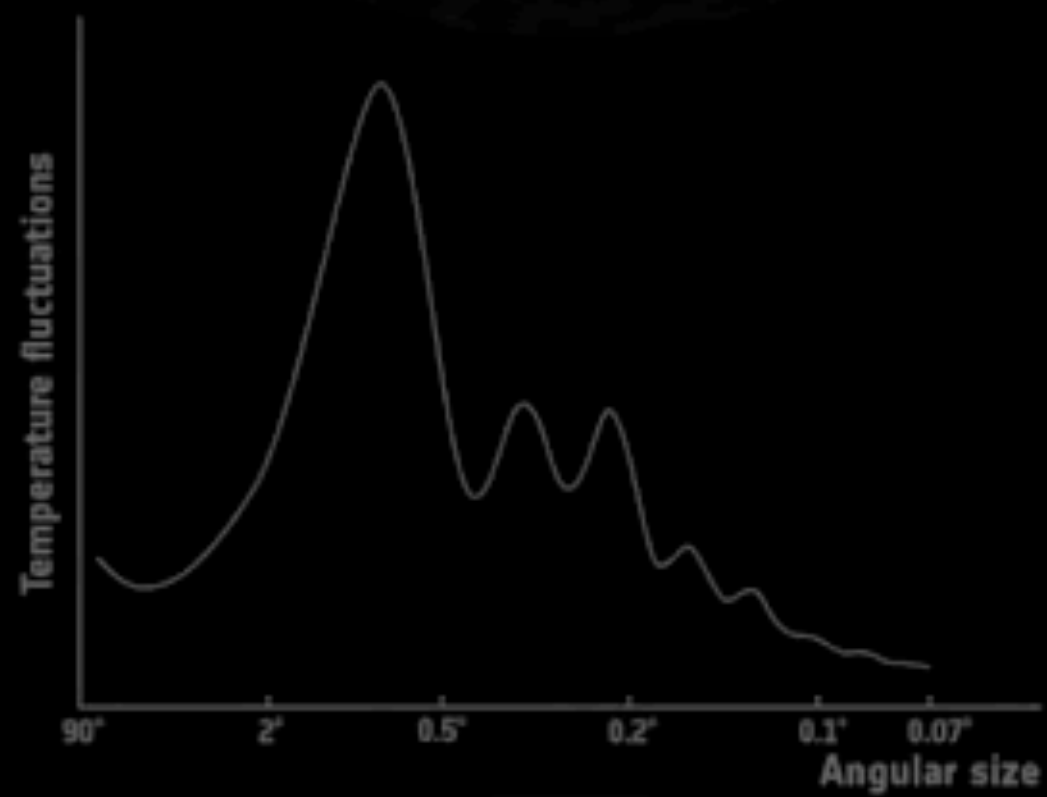
- When matter and radiation were hotter than 3000 K, matter was completely ionised. The Universe was filled with plasma, which behaves just like a soup
- Think about a Miso soup (if you know what it is). Imagine throwing Tofus into a Miso soup, while changing the density of Miso
- And imagine watching how ripples are created and propagate throughout the soup



Power Spectrum Analysis

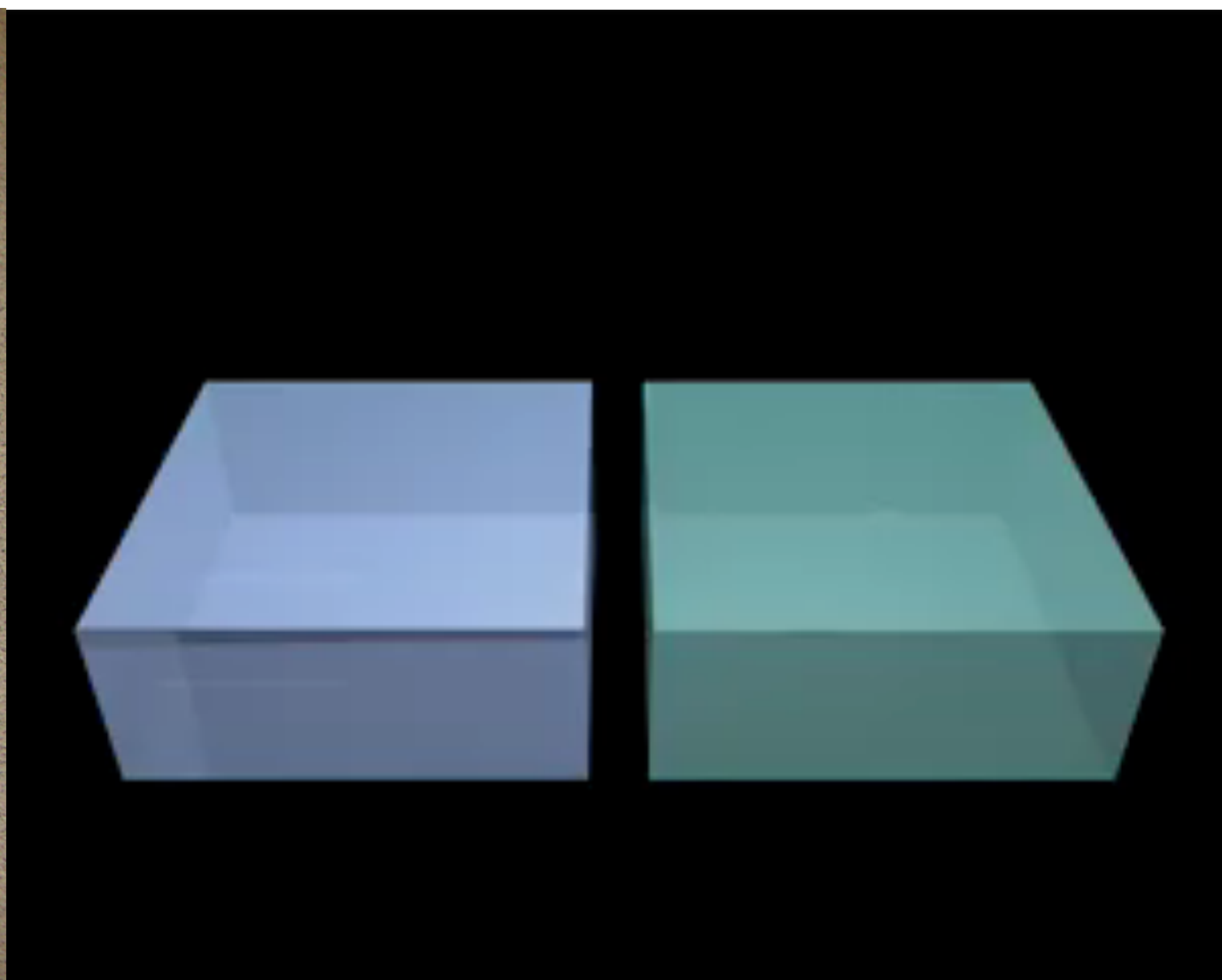
- Decompose temperature fluctuations in the sky into a set of waves with various wavelengths
- Make a diagram showing the strength of each wavelength

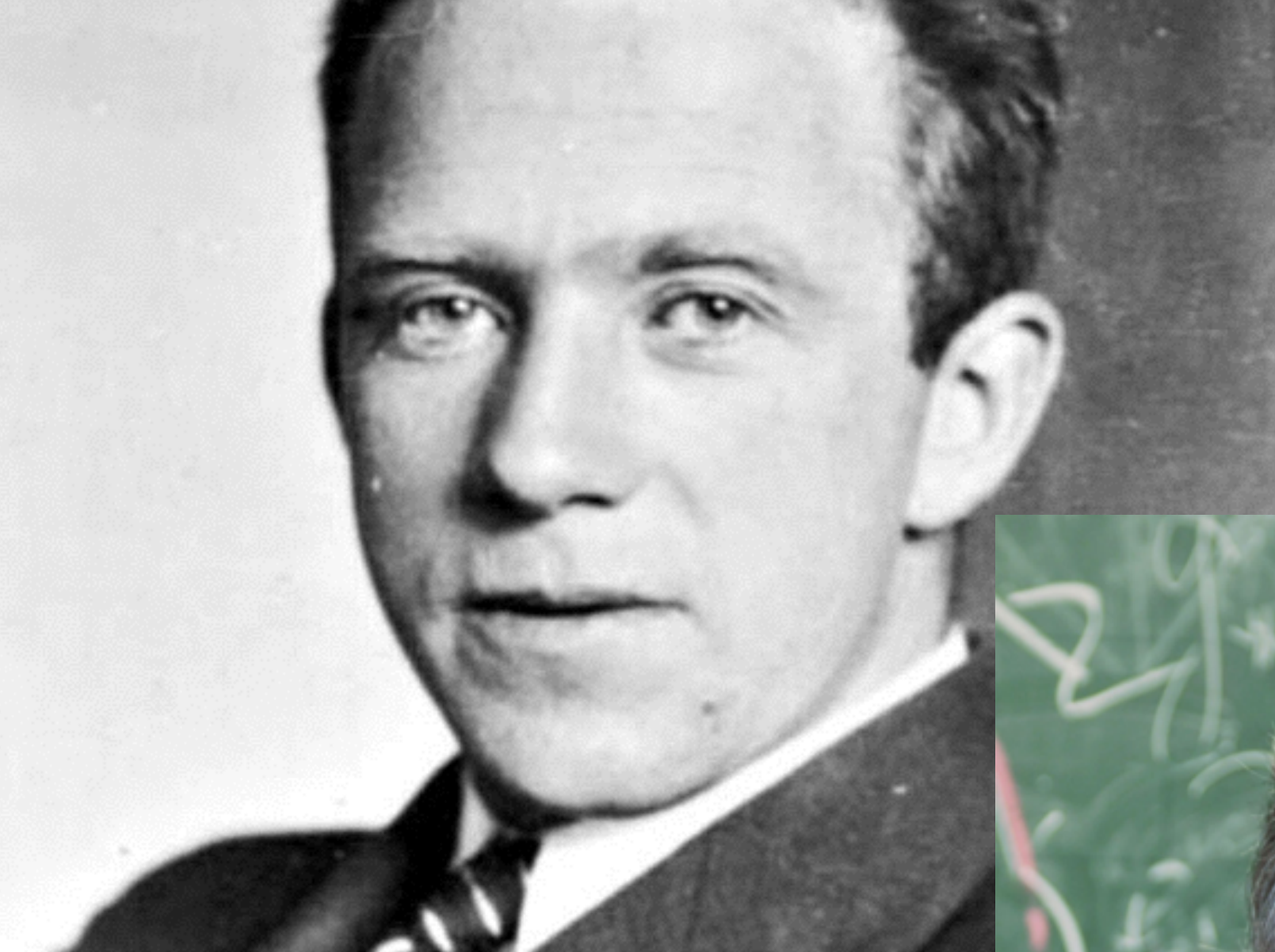




Origin of Fluctuations

- **Who dropped** those Tofus into the cosmic Miso soup?





Werner Heisenberg
(1901–1976)

Slava Mukhanov
(Munich Univ.)



Leading Idea

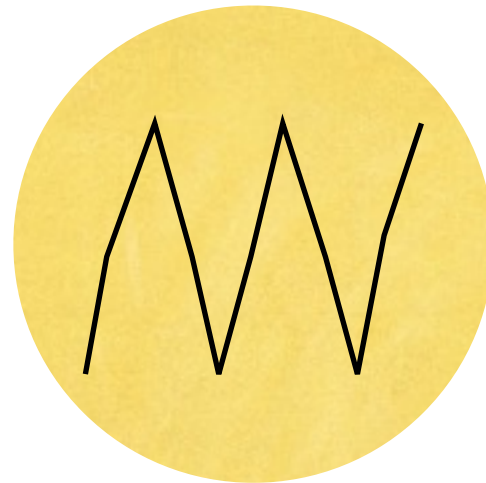
- **Quantum Mechanics at work in the early Universe**
(Mukhanov & Chibisov, 1981)
- **Heisenberg's Uncertainty Principle:**
 - **[Energy you can borrow] x [Time you borrow] $\sim h$**
 - **Time was very short in the early Universe =
You could borrow a lot of energy**
- **Those energies became the origin of fluctuations**
- How did quantum fluctuations on the microscopic scales become macroscopic fluctuations over cosmological sizes?

Cosmic Inflation

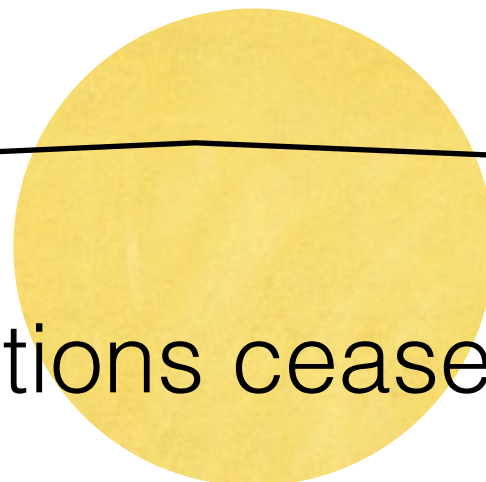
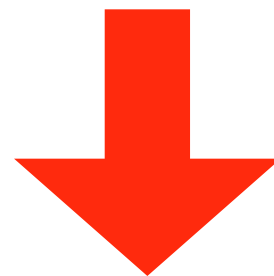
- In a tiny fraction of a second, the size of an atomic nucleus became the size of the Solar System
- In 10^{-36} second, space was stretched by at least a factor of 10^{26}

Stretching Micro to Macro

Quantum fluctuations on
microscopic scales



Inflation!



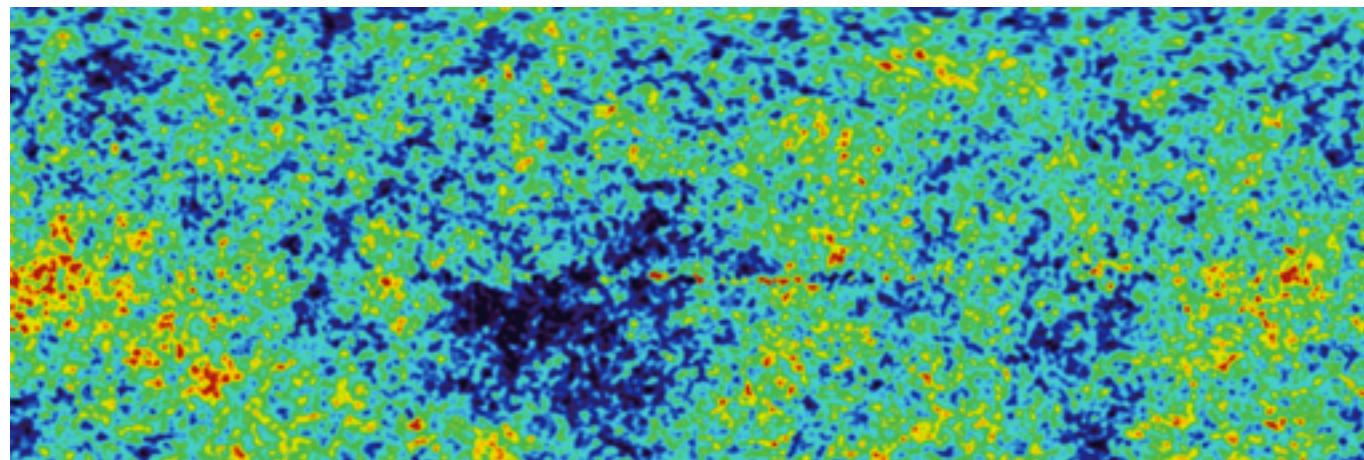
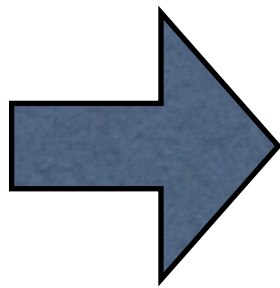
- Quantum fluctuations cease to be quantum
- Become macroscopic, classical fluctuations

Key Predictions of Inflation

 ζ

scalar
mode

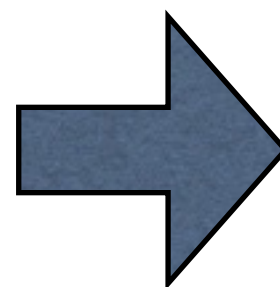
- Fluctuations we observe today in CMB and the matter distribution originate from quantum fluctuations generated during inflation

 h_{ij}

tensor
mode

- There should also be ***ultra-long-wavelength gravitational waves*** generated during inflation

Starobinsky (1979)



We measure distortions in space

- A distance between two points in space

$$d\ell^2 = a^2(t)[1 + 2\zeta(\mathbf{x}, t)][\delta_{ij} + h_{ij}(\mathbf{x}, t)]dx^i dx^j$$

- ζ : “curvature perturbation” (scalar mode)
 - Perturbation to the determinant of the spatial metric
- h_{ij} : “gravitational waves” (tensor mode)
 - Perturbation that does not change the determinant (area)



$$\sum_i h_{ii} = 0$$

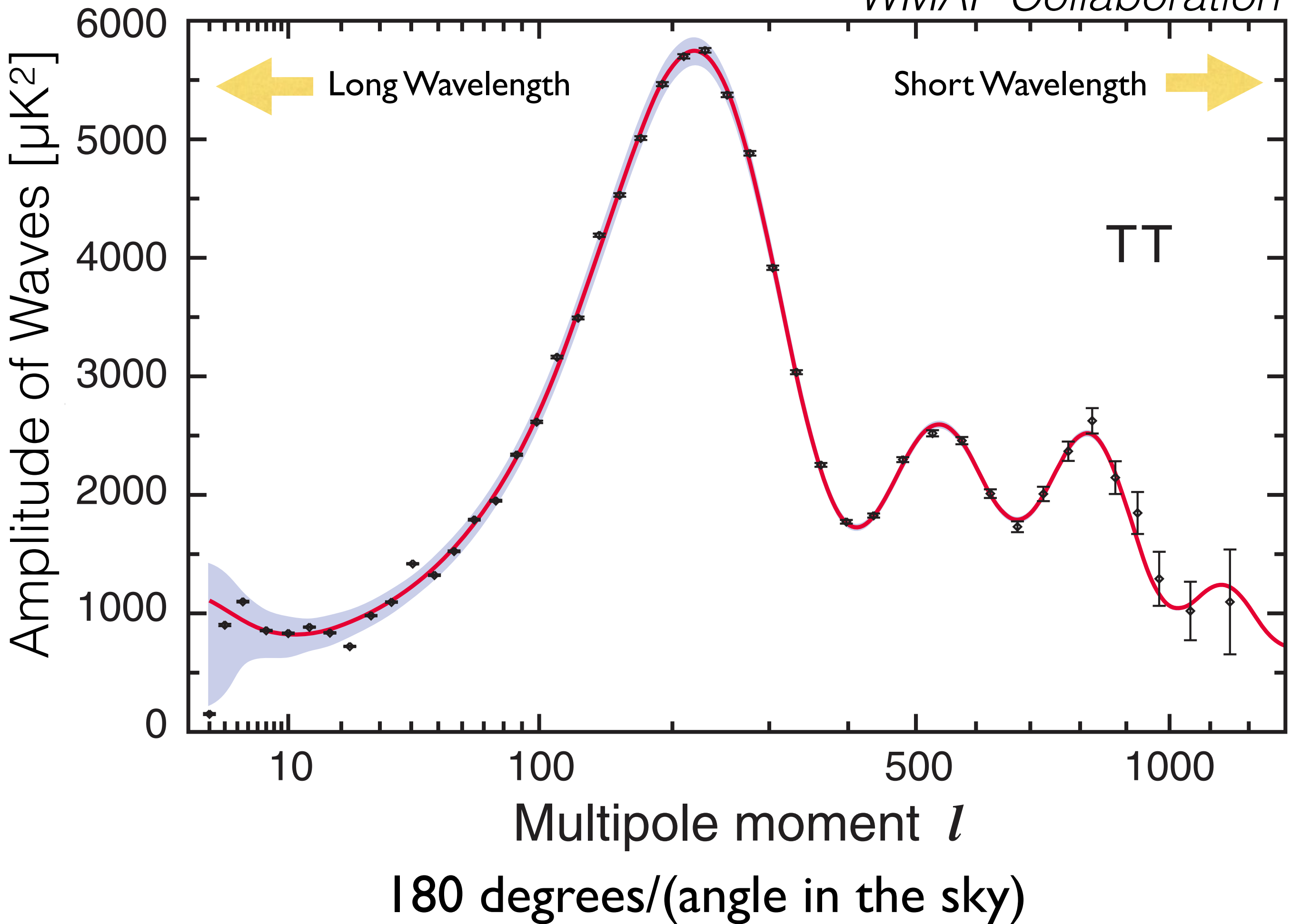
Heisenberg's Uncertainty Principle

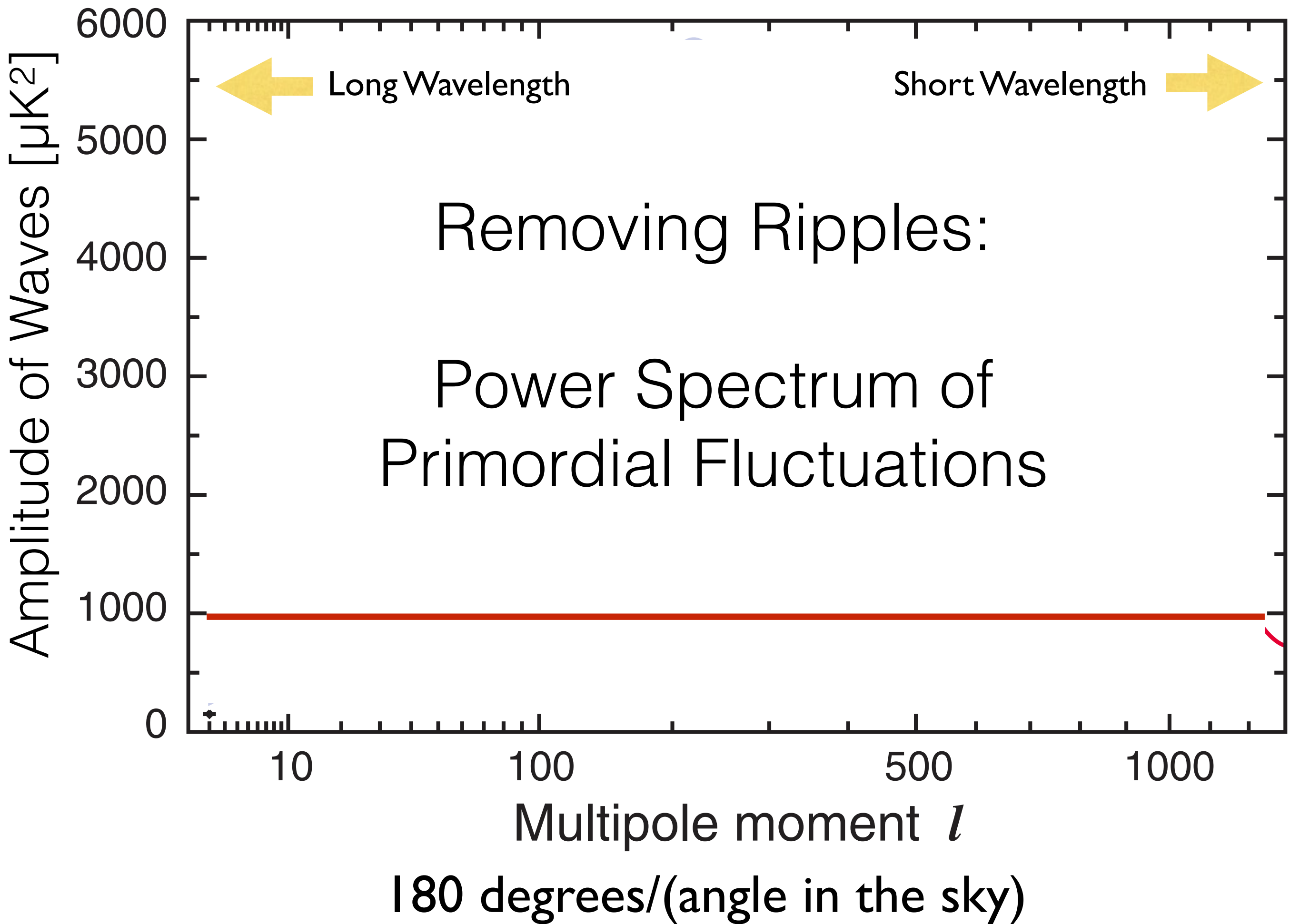
- [Energy you can borrow] x [Time you borrow] = constant
- Suppose that the distance between two points increases in proportion to **$a(t)$** [which is called the scale factor] by the expansion of the universe
- Define the “expansion rate of the universe” as

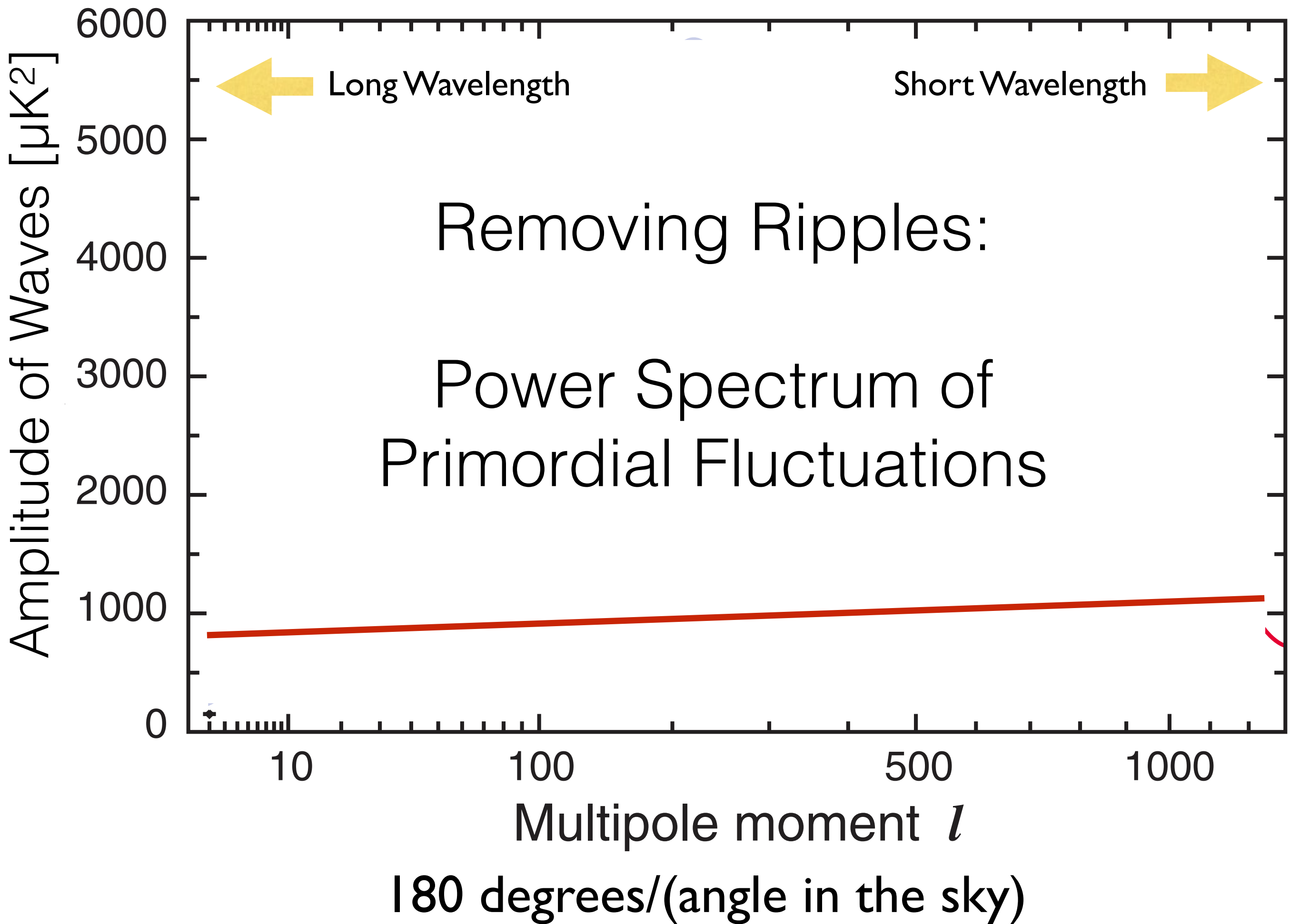
$$H \equiv \frac{\dot{a}}{a} \quad [\text{This has units of } 1/\text{time}]$$

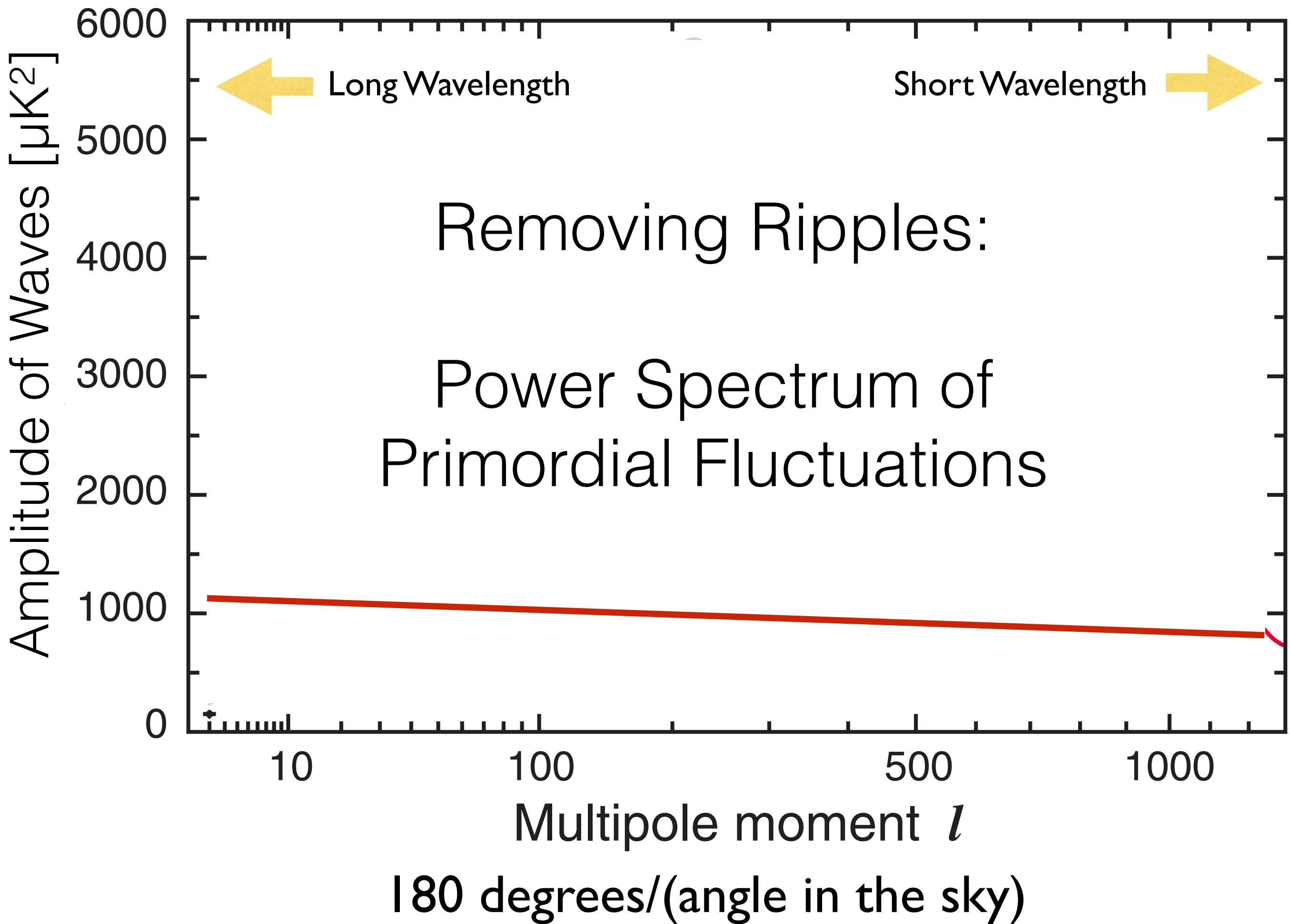
Fluctuations are proportional to H

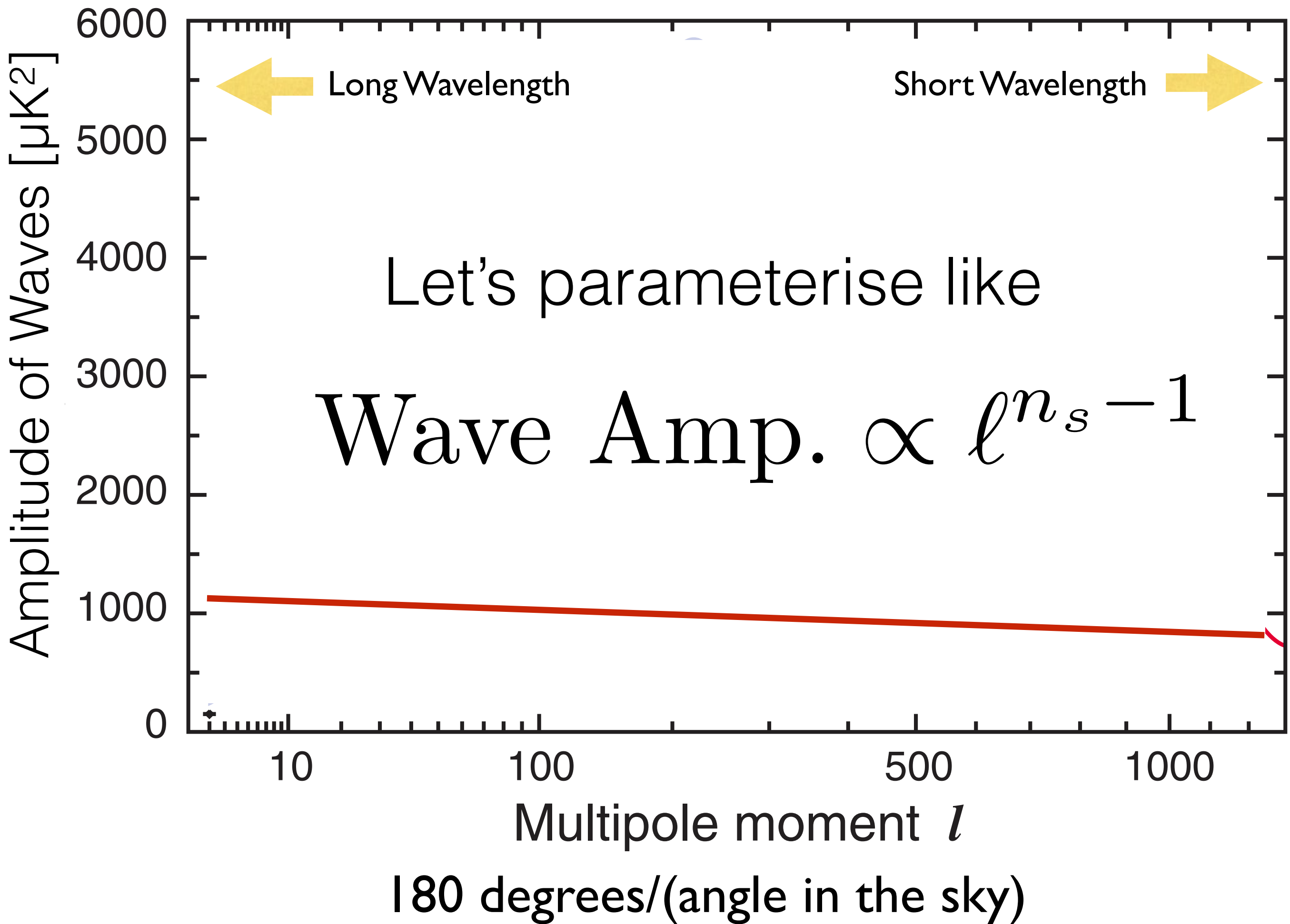
- [Energy you can borrow] x [Time you borrow] = constant
- $H \equiv \frac{\dot{a}}{a}$ [This has units of 1/time]
- Then, **both ζ and h_{ij} are proportional to H**
- Inflation occurs in 10^{-36} second - this is such a short period of time that you can borrow a lot of energy!
 H during inflation in energy units is 10^{14} GeV

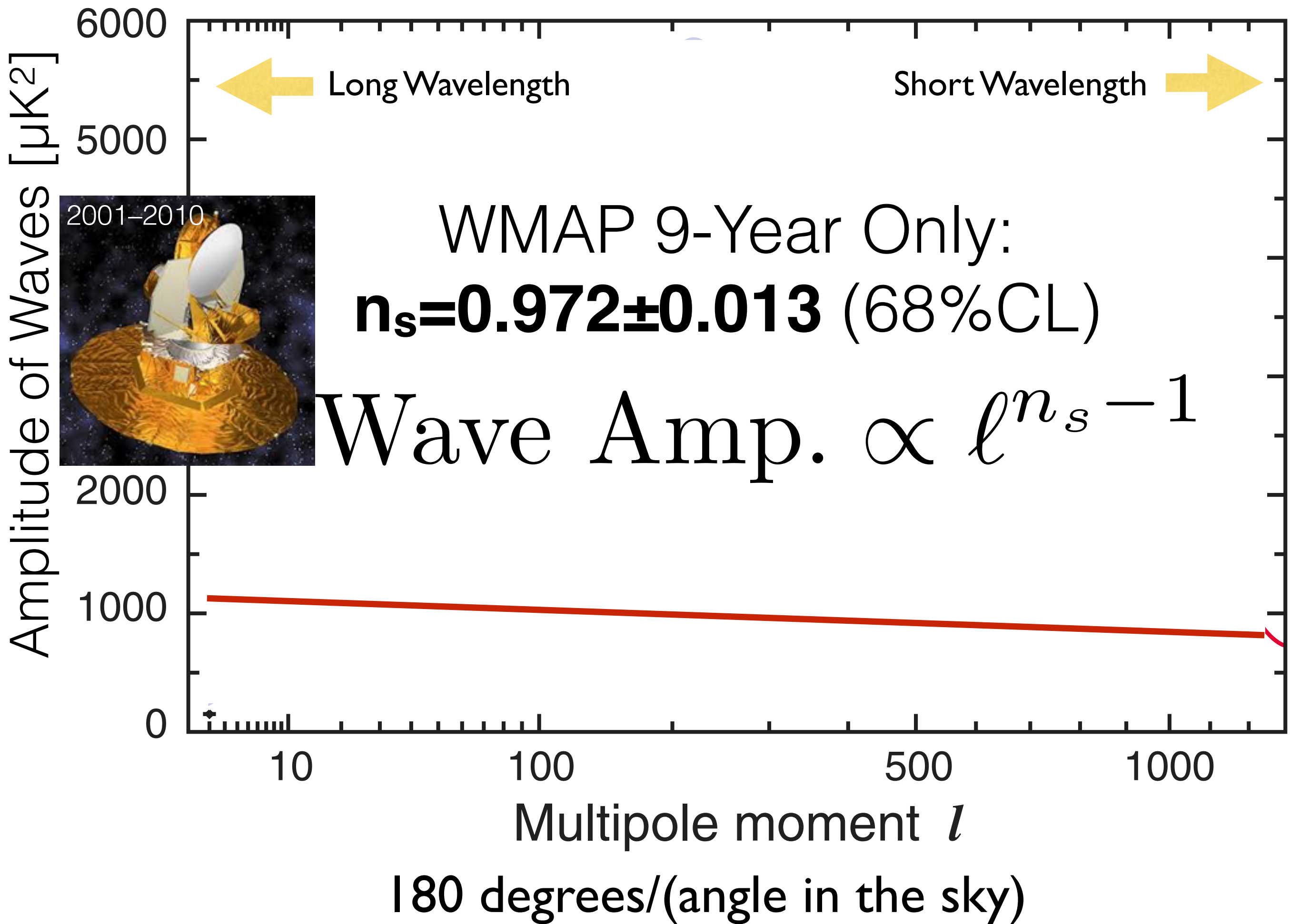


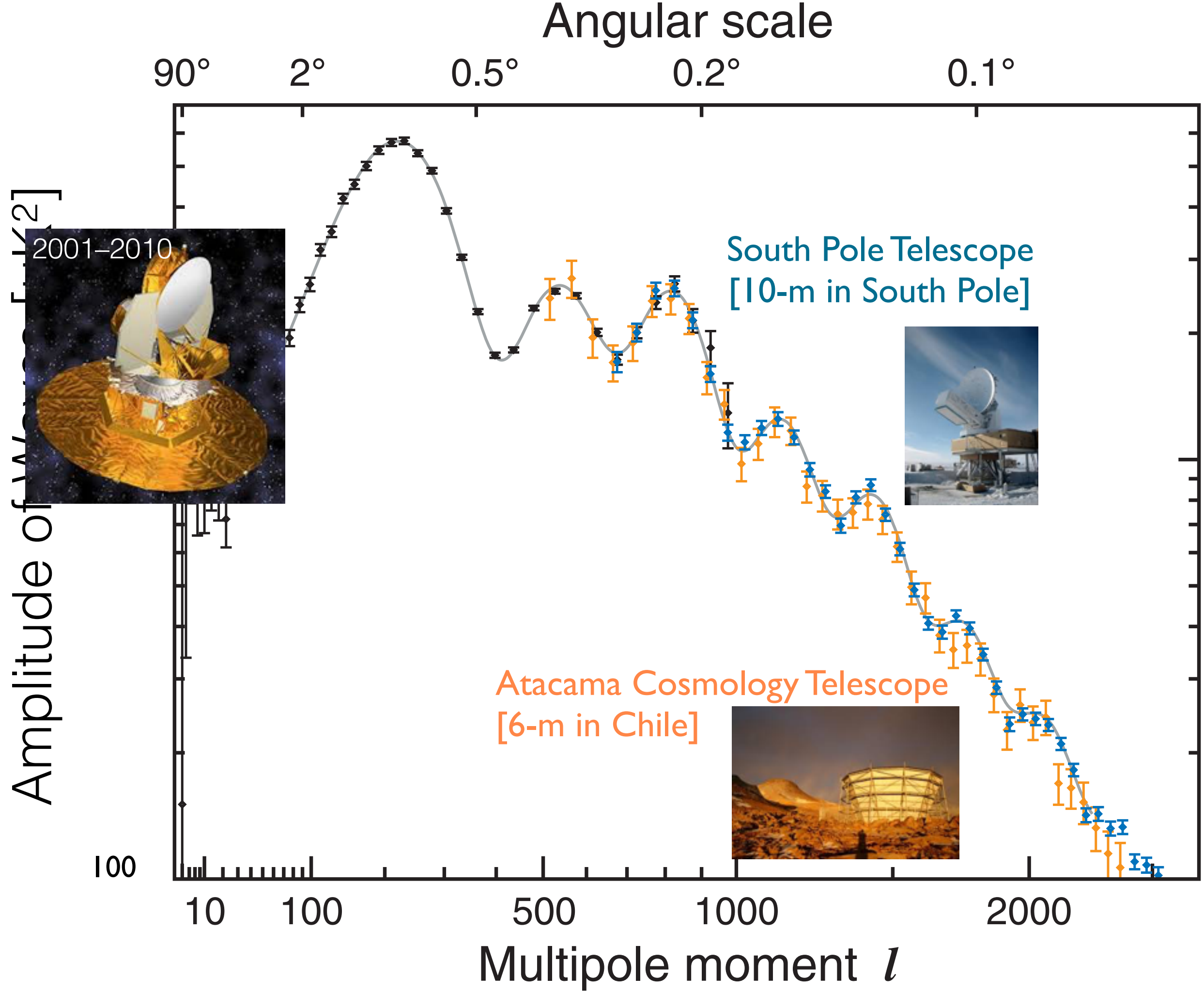


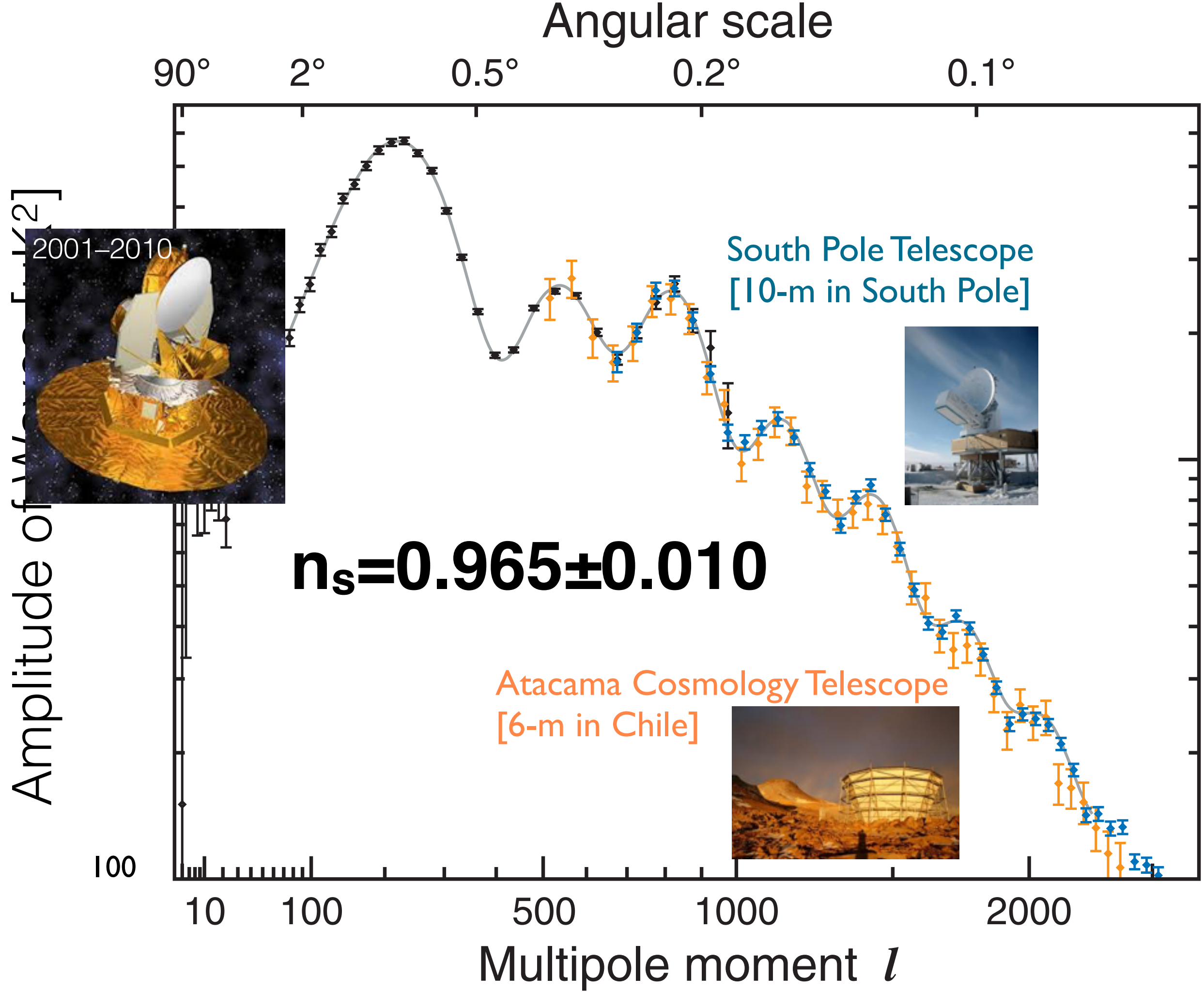


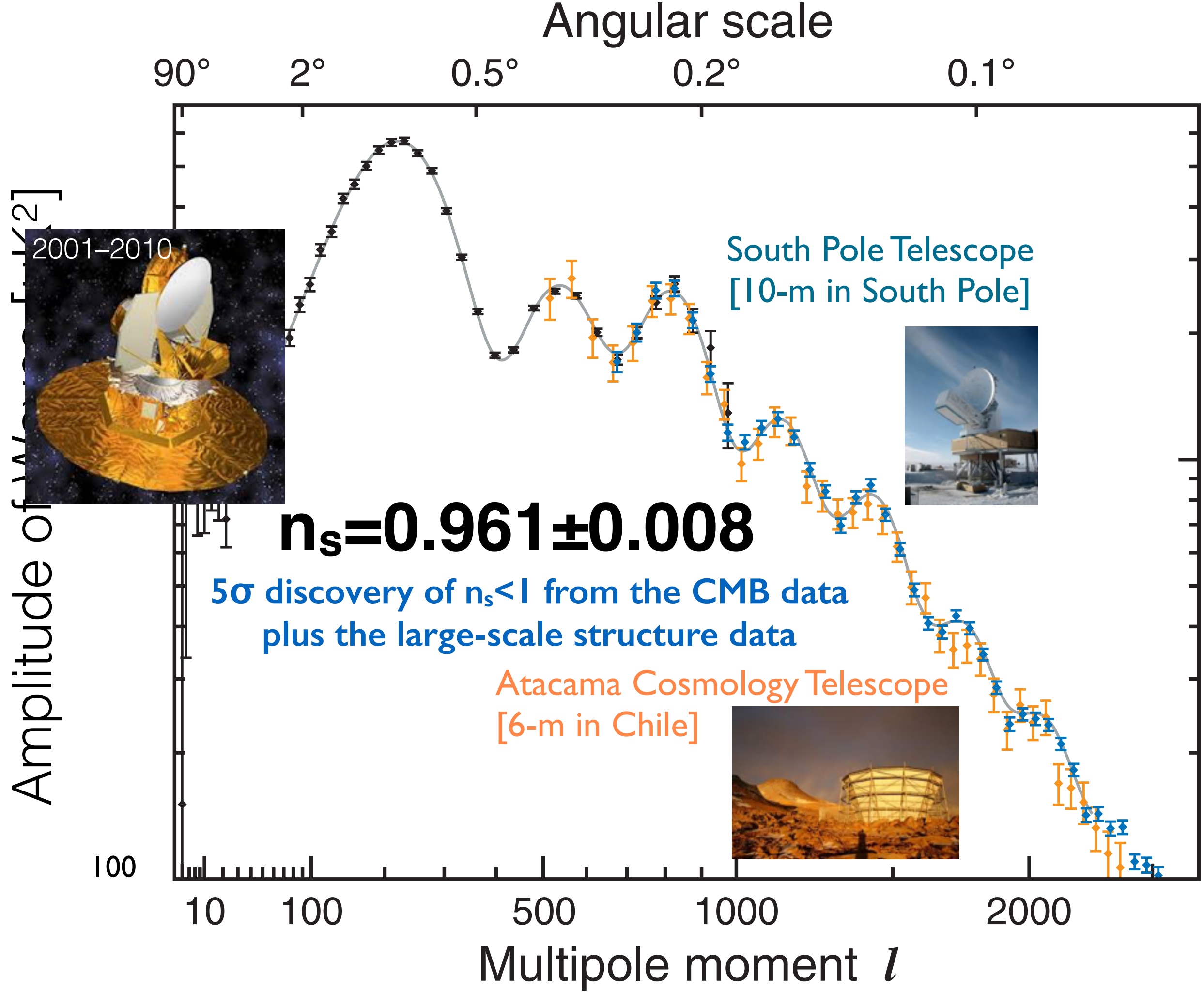












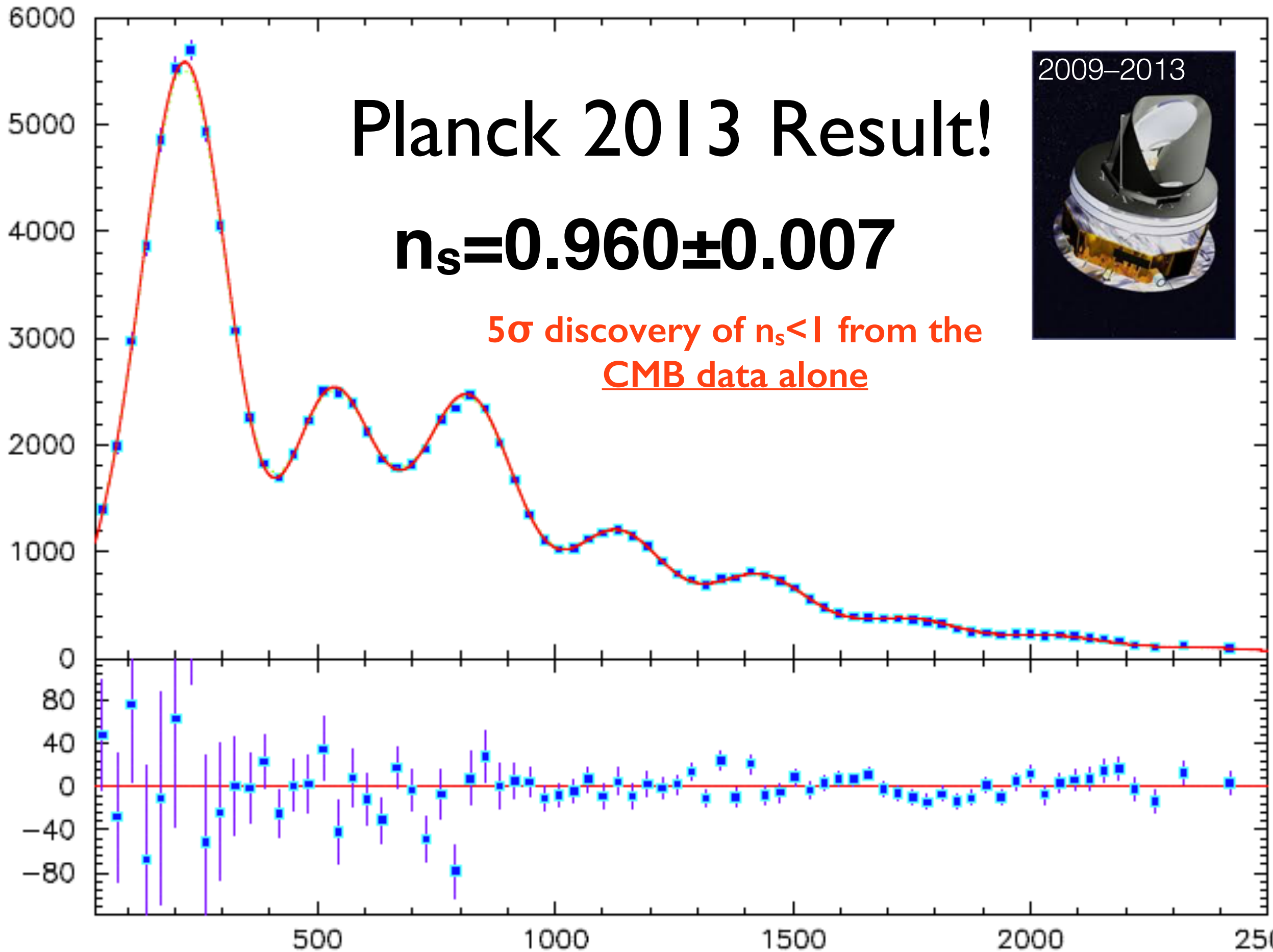
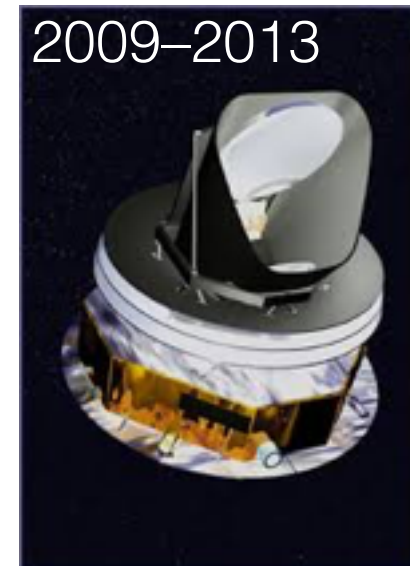
Residual Amplitude of Waves [μK^2]

Planck 2013 Result!

$$n_s = 0.960 \pm 0.007$$

**5 σ discovery of $n_s < 1$ from the
CMB data alone**

2009–2013



l 80 degrees/(angle in the sky)

Predicted in 1981.
Finally discovered in 2013
by WMAP and Planck

- Inflation must end
- Inflation predicts $n_s \sim 1$, but not exactly equal to 1. Usually $n_s < 1$ is expected
- **The discovery of $n_s < 1$ has been a dream of cosmologists since 1992,** when the CMB anisotropy was first discovered and $n_s \sim 1$ (to within 30%) was indicated



Slava Mukhanov said in his 1981 paper that n_s should be less than 1

An implication of $n_s < 1$

- Accelerated expansion during the early universe

$$\frac{\ddot{a}}{a} = \dot{H} + H^2 > 0 \quad \rightarrow \quad \epsilon \equiv -\frac{\dot{H}}{H^2} < 1$$

- When $\epsilon \ll 1$, the universe expands quasi-exponentially.
- If $\epsilon=0$, space-time is exactly de Sitter:
$$ds^2 = -dt^2 + e^{2Ht} d\mathbf{x}^2$$
- *But, inflation never ends if $\epsilon=0$.* When $\epsilon \ll 1$, space-time is nearly, but not exactly, de Sitter:

$$ds^2 = -dt^2 + e^{2 \int dt' H(t')} d\mathbf{x}^2$$

Symmetry of de Sitter Space

$$ds^2 = -dt^2 + e^{2Ht} d\mathbf{x}^2$$

- De Sitter spacetime is invariant under 10 isometries (transformations that keep ds^2 invariant):

- Time translation, followed by space dilation

$$t \rightarrow t - \lambda/H, \quad \mathbf{x} \rightarrow e^\lambda \mathbf{x}$$

- Spatial rotation, $\mathbf{x} \rightarrow R\mathbf{x}$

- Spatial translation, $\mathbf{x} \rightarrow \mathbf{x} + c$

- Three more transformations irrelevant to this talk

Symmetry of de Sitter Space

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discovered in 2012/13

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- Three more transformations irrelevant to this talk

Rotational Invariance?

$$ds^2 = -dt^2 + e^{2Ht} d\mathbf{x}^2$$

- De Sitter spacetime is invariant under 10 isometries (transformations that keep ds^2 invariant):

- Time translation, followed by space dilation



$$t \rightarrow t - \lambda/H, \quad \mathbf{x} \rightarrow e^\lambda \mathbf{x}$$

discovered in 2012/13

- Spatial rotation, $\mathbf{x} \rightarrow R\mathbf{x}$

Is this symmetry valid?

- Spatial translation, $\mathbf{x} \rightarrow \mathbf{x} + c$

- Three more transformations irrelevant to this talk

Anisotropic Expansion

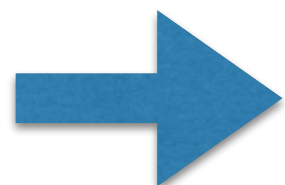
$$ds^2 = -dt^2 + e^{2Ht} \left[e^{-2\beta(t)} dx^2 + e^{2\beta(t)} (dy^2 + dz^2) \right]$$

- How large can $\dot{\beta}/H$ be during inflation?
- In single scalar field theories, Einstein's equation gives

$$\dot{\beta} \propto e^{-3Ht}$$

- But, the presence of anisotropic stress in the stress-energy tensor can source a **sustained** period of anisotropic expansion:

$$T_j^i = P\delta_j^i + \pi_j^i \quad \text{with} \quad \pi_1^1 = -\frac{2}{3}\mathcal{V}, \quad \pi_2^2 = \pi_3^3 = \frac{1}{3}\mathcal{V}$$


$$\ddot{\beta} + 3H\dot{\beta} = \frac{1}{3}\mathcal{V}$$

Inflation with a vector field

- Consider that there existed a vector field at the beginning of inflation:

$$A_\mu = (0, u(t), 0, 0)$$

A_1 : Preferred direction in space at the initial time

- You might ask where A_μ came from. Well, if we have a scalar field and a tensor field (gravitational wave), why not a vector?
- The conceptual problem of this setting is not the existence of a vector field, but that it requires A_1 that is homogeneous over a few Hubble lengths before inflation
- But, this problem is common with the original inflation, which requires ϕ that is homogeneous over a few Hubble lengths, in order for inflation to occur in the first place!

Observational Consequence

- Broken rotational invariance makes **the power spectrum depend on a direction of wavenumber**

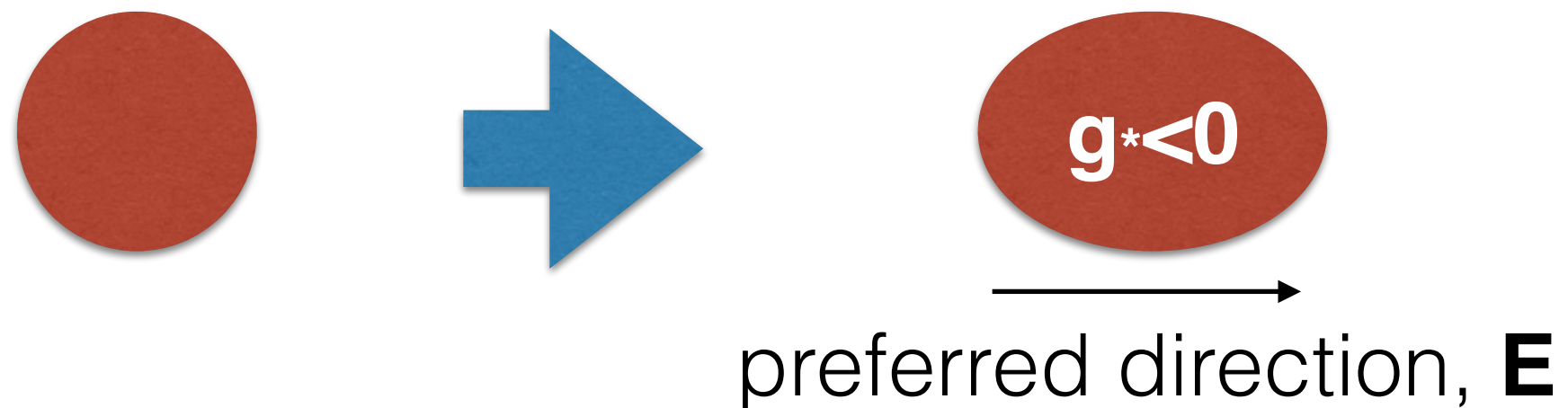
$$P(k) \rightarrow P(\mathbf{k}) = P_0(k) \left[1 + g_*(k) (\hat{k} \cdot \hat{E})^2 \right]$$

where $\hat{E}$ is a preferred direction in space

- A model predicts: $g_*(k) = -\mathcal{O}(1) \times 24 I_k N_k^2$
- $I = [\text{energy density of a vector}]/[\text{energy density of a scalar}]$, divided by ε
- $N \sim 60$ is the number of e-fold of inflation counted from the end of inflation

Signatures in CMB

- This “quadrupolar” modulation of the power spectrum turns a circular hot/cold spot of CMB into an elliptical one



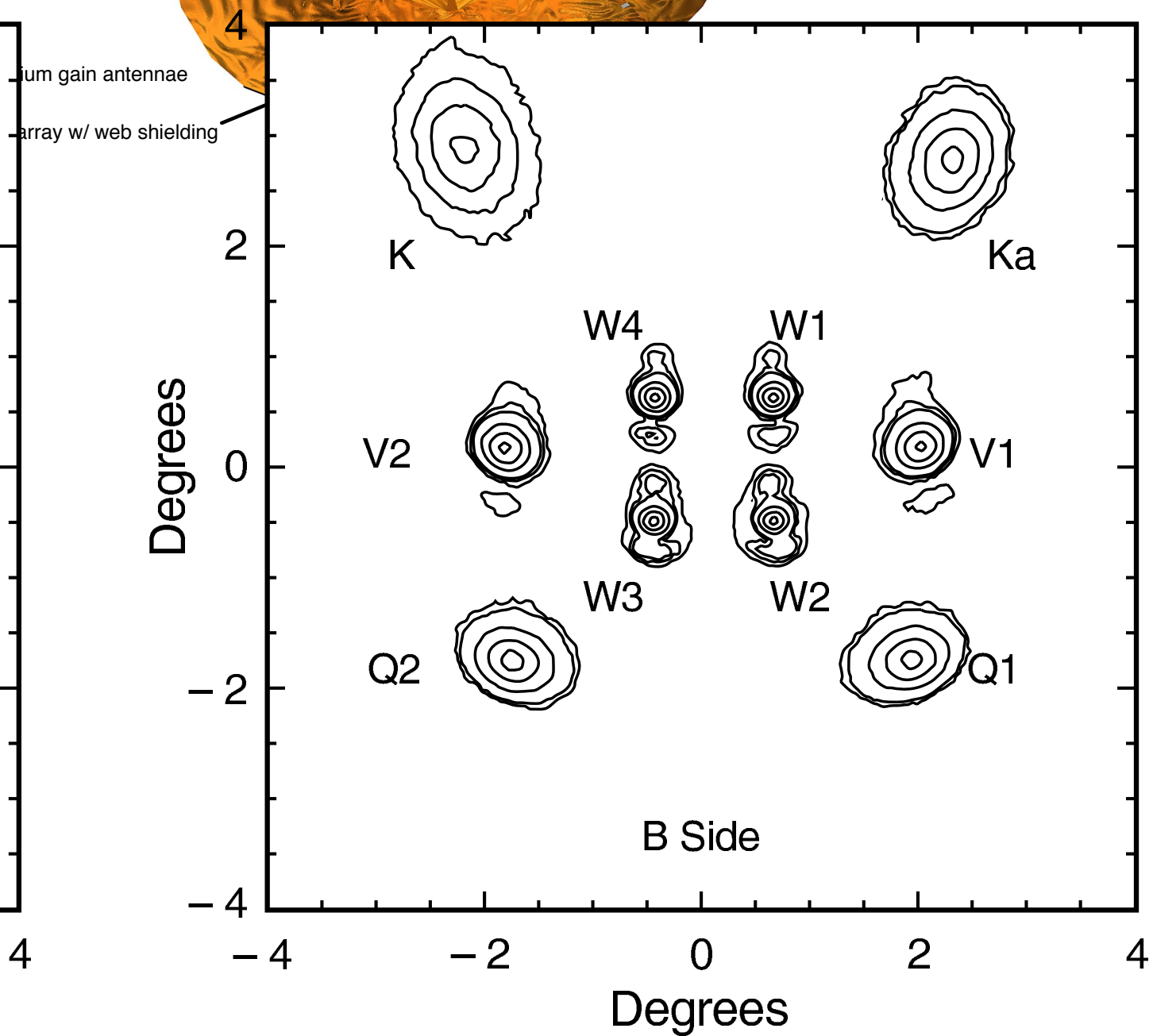
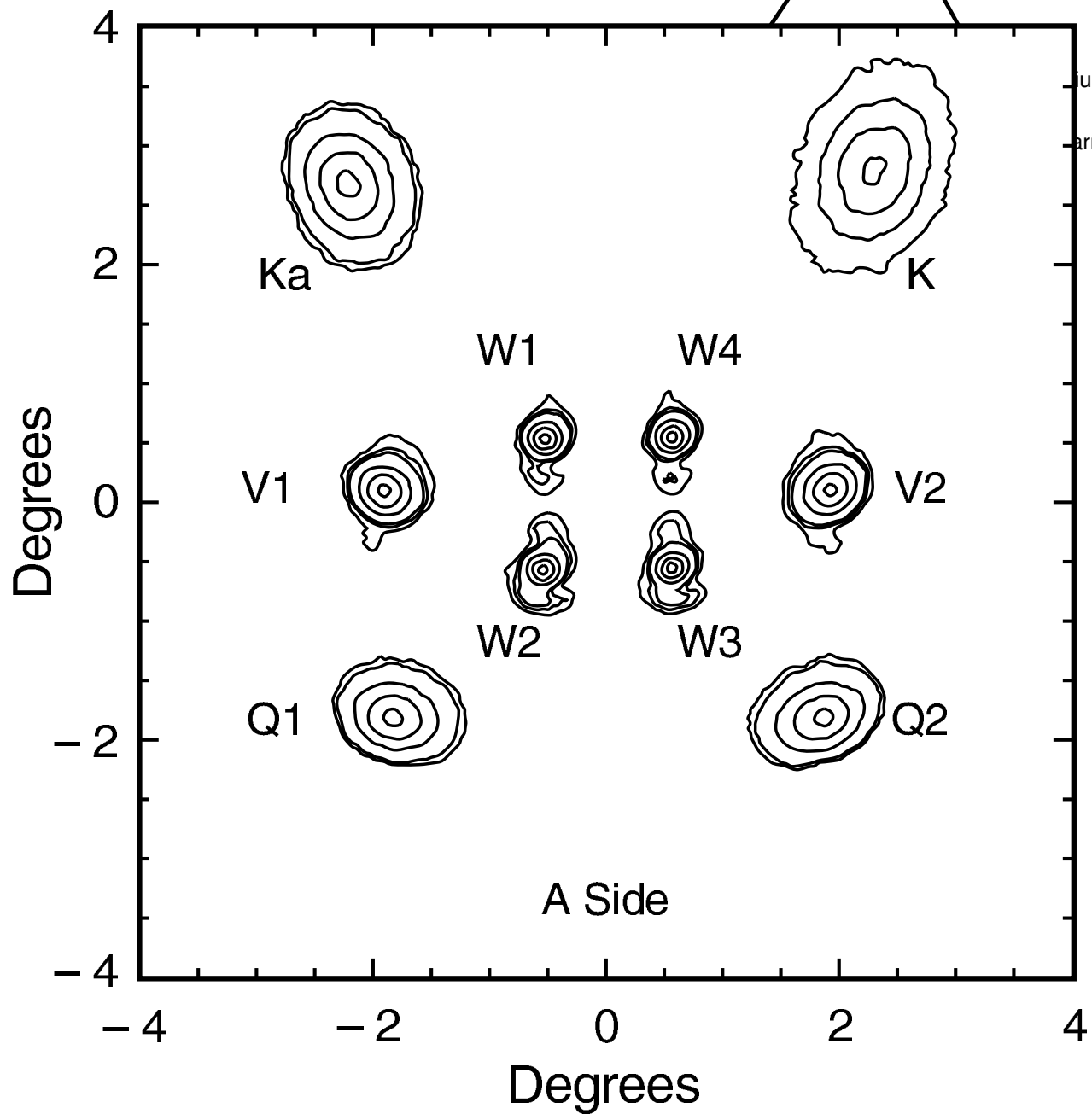
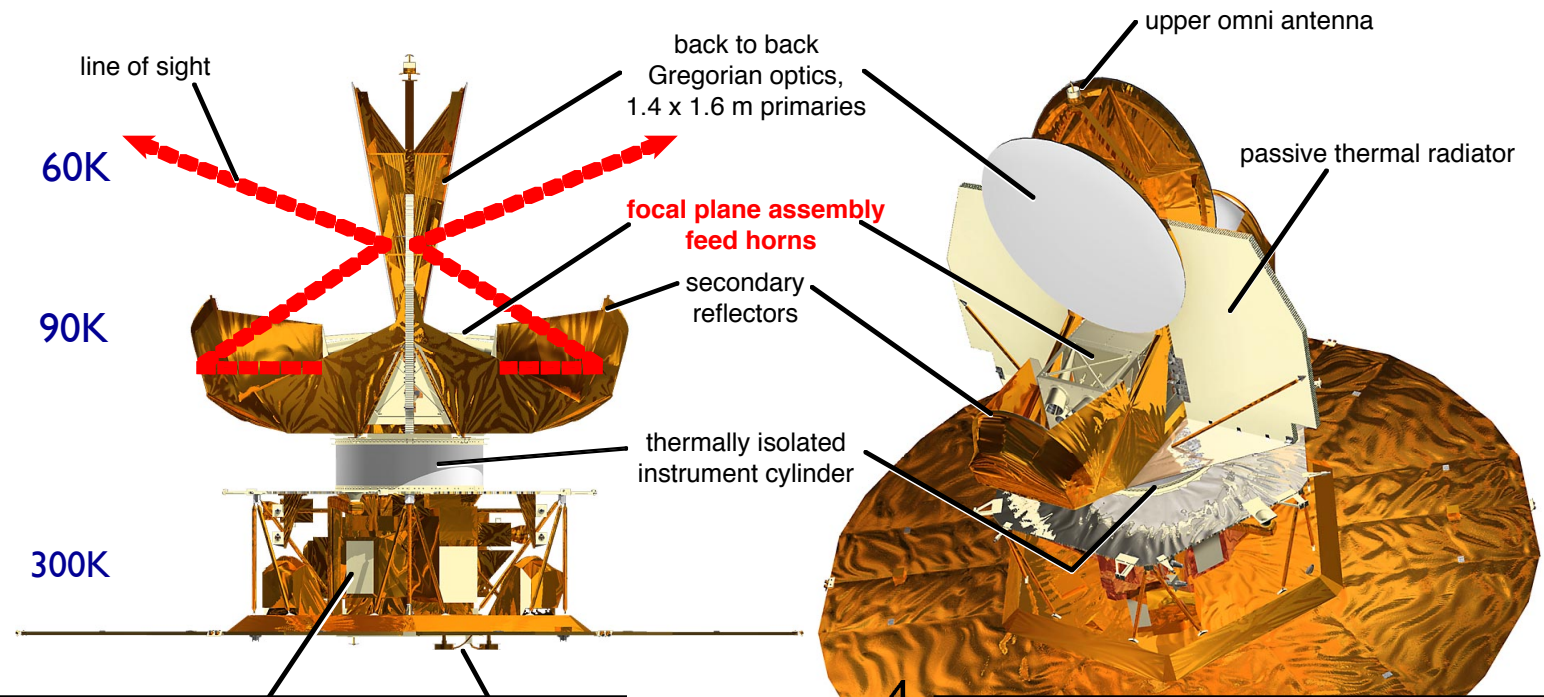
- This is a local effect, rather than a global effect: the power spectrum measured at *any location* in the sky is modulated by $(\hat{k} \cdot \hat{E})^2$

A Beautiful Story

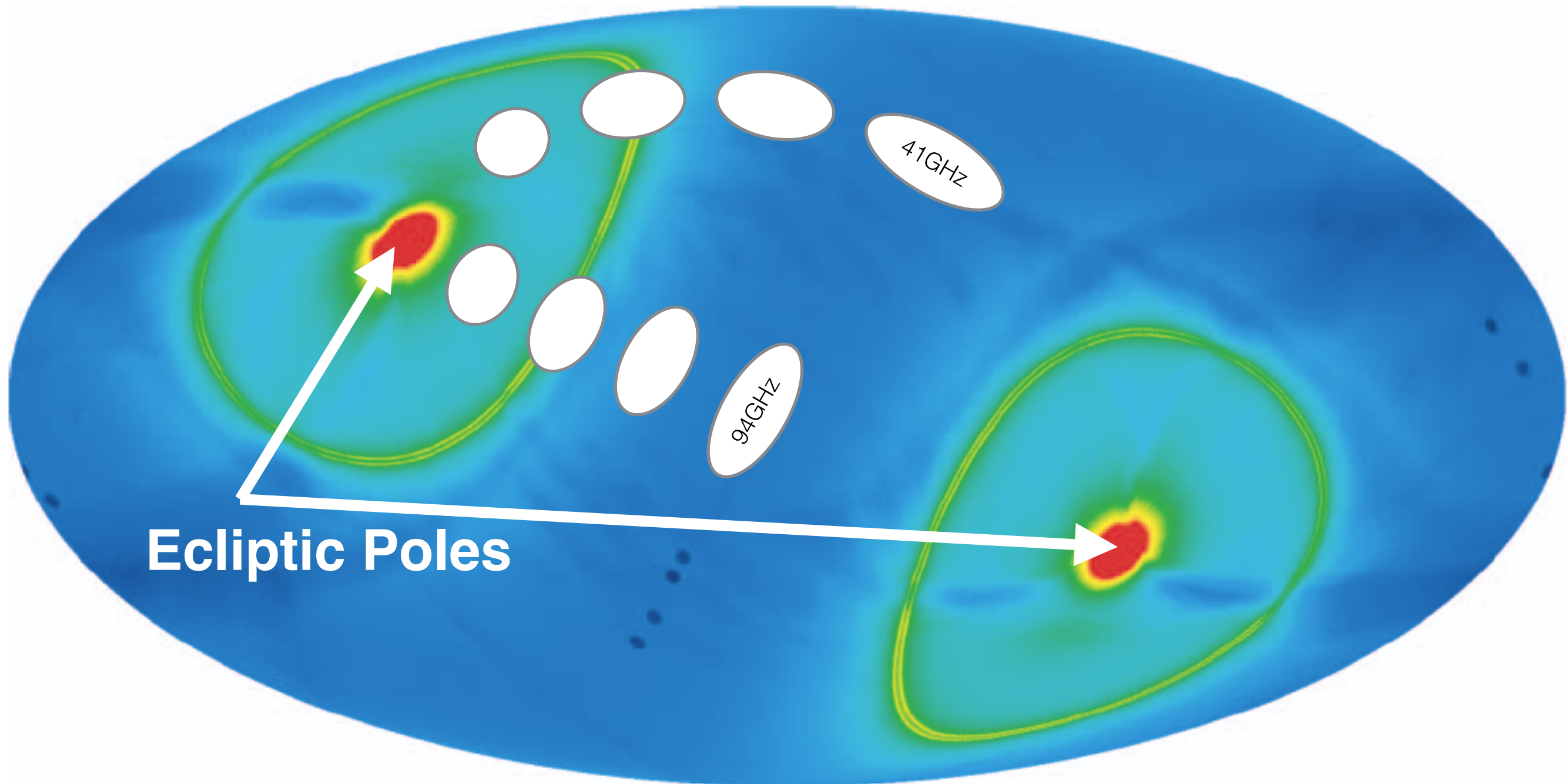
- In 2007, Ackerman, Carroll and Wise proposed g^* as a powerful probe of rotational symmetry
- In 2009, Groeneboom and Eriksen reported a significant detection, $g^*=0.15\pm0.04$, in the WMAP data at 94 GHz
- **Surprise!** And a beautiful story - a new observable proposed by theorists was looked for in the data, and was found

Subsequent Story

- In 2010, Groeneboom et al. reported that the WMAP data at 41 GHz gave the opposite sign: $g^* = -0.18 \pm 0.04$, suggesting instrumental systematics
- The best-fit preferred direction in the sky was the ecliptic pole
- Elliptical beam (point spread function) of WMAP was a culprit!

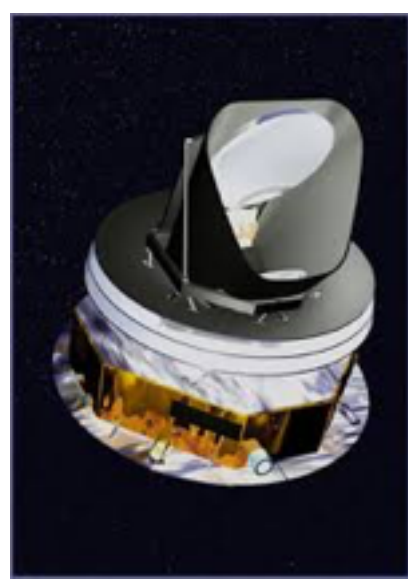


of observations in Galactic coordinates



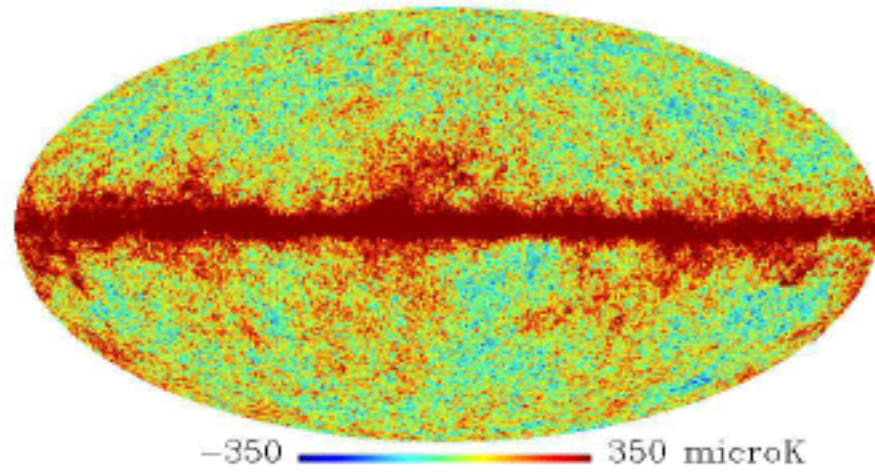
- WMAP visits ecliptic poles from many different directions, circularising beams
- WMAP visits ecliptic planes with 30% of possible angles

Planck 2013 Data

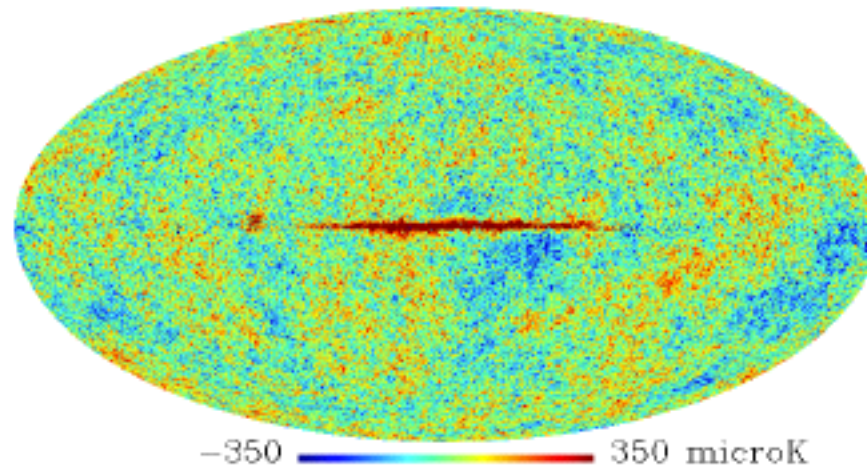


- We analysed the Planck 2013 temperature data at 143GHz, and found significant **$g^* = -0.111 \pm 0.013$** [after removing the foreground emission]
- This is consistent with what we expect from the beam ellipticity of the Planck data
- After subtracting the effect of beam ellipticities, no evidence for g^* was found

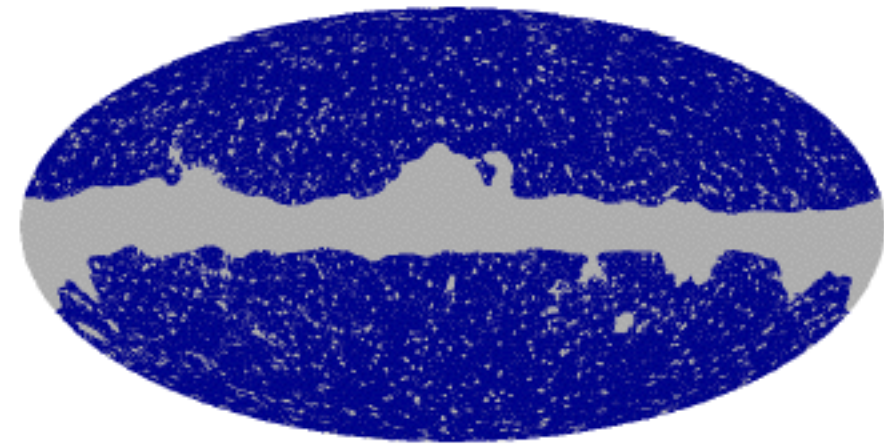
143Ghz data



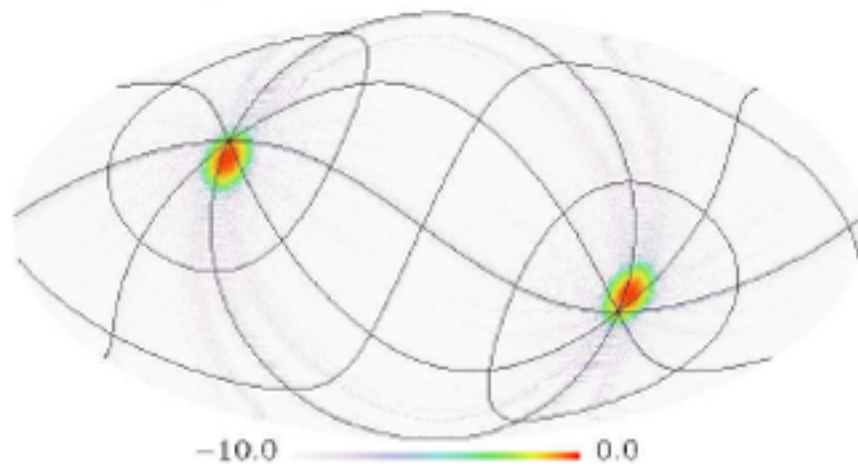
Foreground-cleaned 143Ghz data



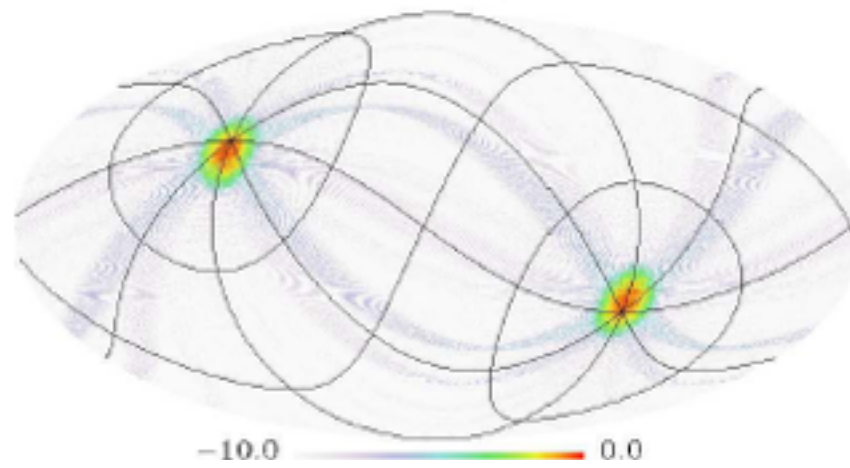
Foreground mask



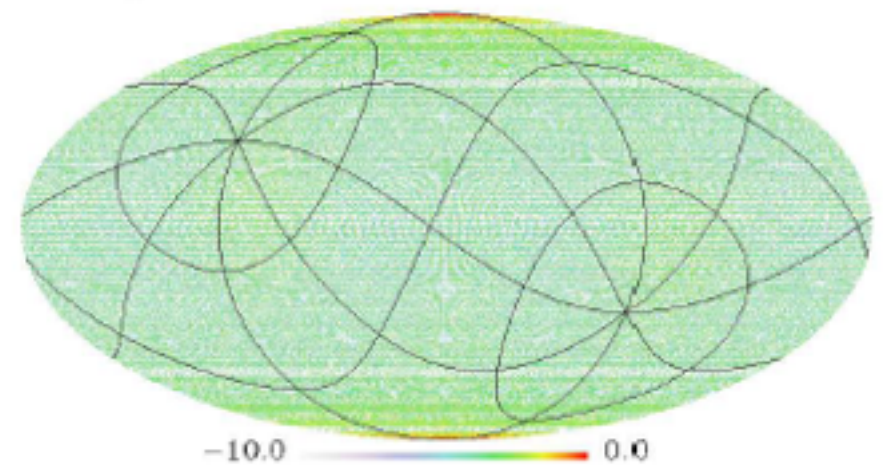
Foreground-cleaned 143Ghz data



Simulated with the asymmetric beam

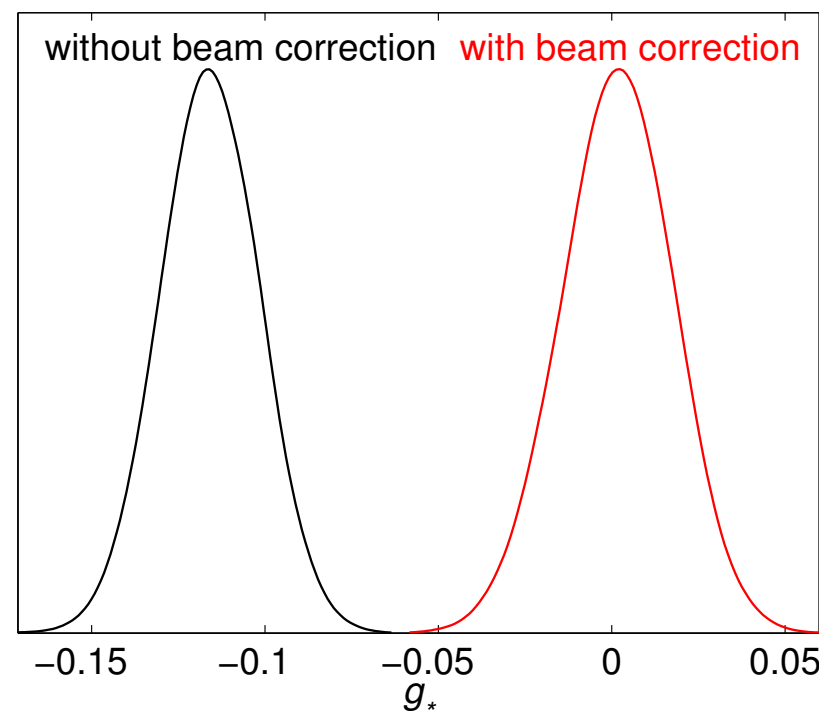


Foreground-cleaned and beam-corrected



$$g^*(\text{raw}) = -0.111 \pm 0.013$$

(68%CL)



$$g^* = 0.002 \pm 0.016$$

(68%CL)

Implication for Rotational Symmetry

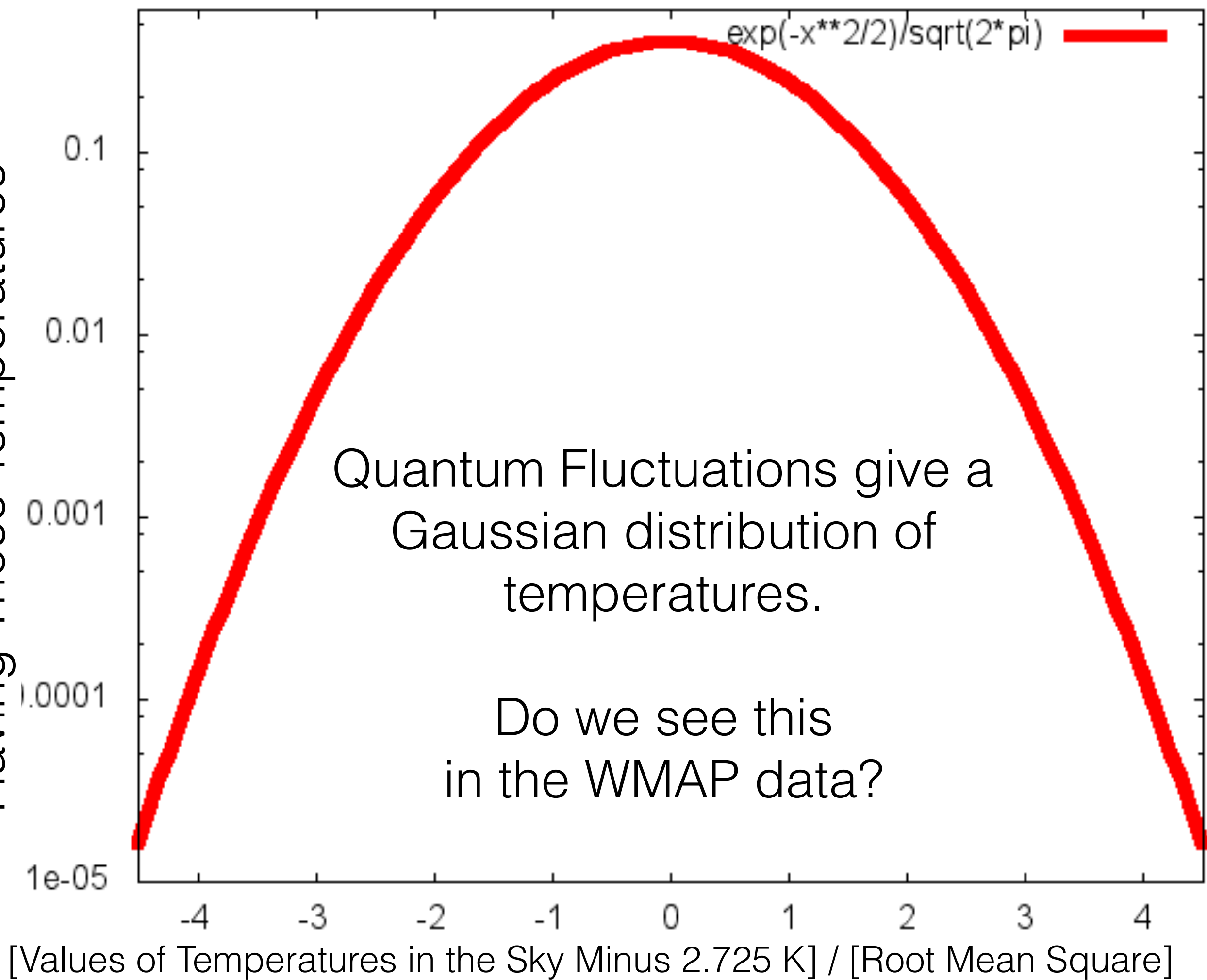
- g^* is consistent with zero, with 95%CL upper bound of $|g^*| < 0.03$
- Comparing this with the model prediction, $|g^*| \sim 24IN^2$, we conclude $I < 5 \times 10^{-7}$
- Thus, $\frac{\dot{\beta}}{H} \approx \frac{\mathcal{V}}{U} \approx \epsilon I < 5 \times 10^{-9}$

Breaking of rotational symmetry is tiny, if any!

[cf: “natural” value is either 10^{-2} or $e^{-3N} = e^{-150}!!$]

How do we know that
primordial fluctuations were of
quantum mechanical origin?

Fraction of the Number of Pixels
Having Those Temperatures



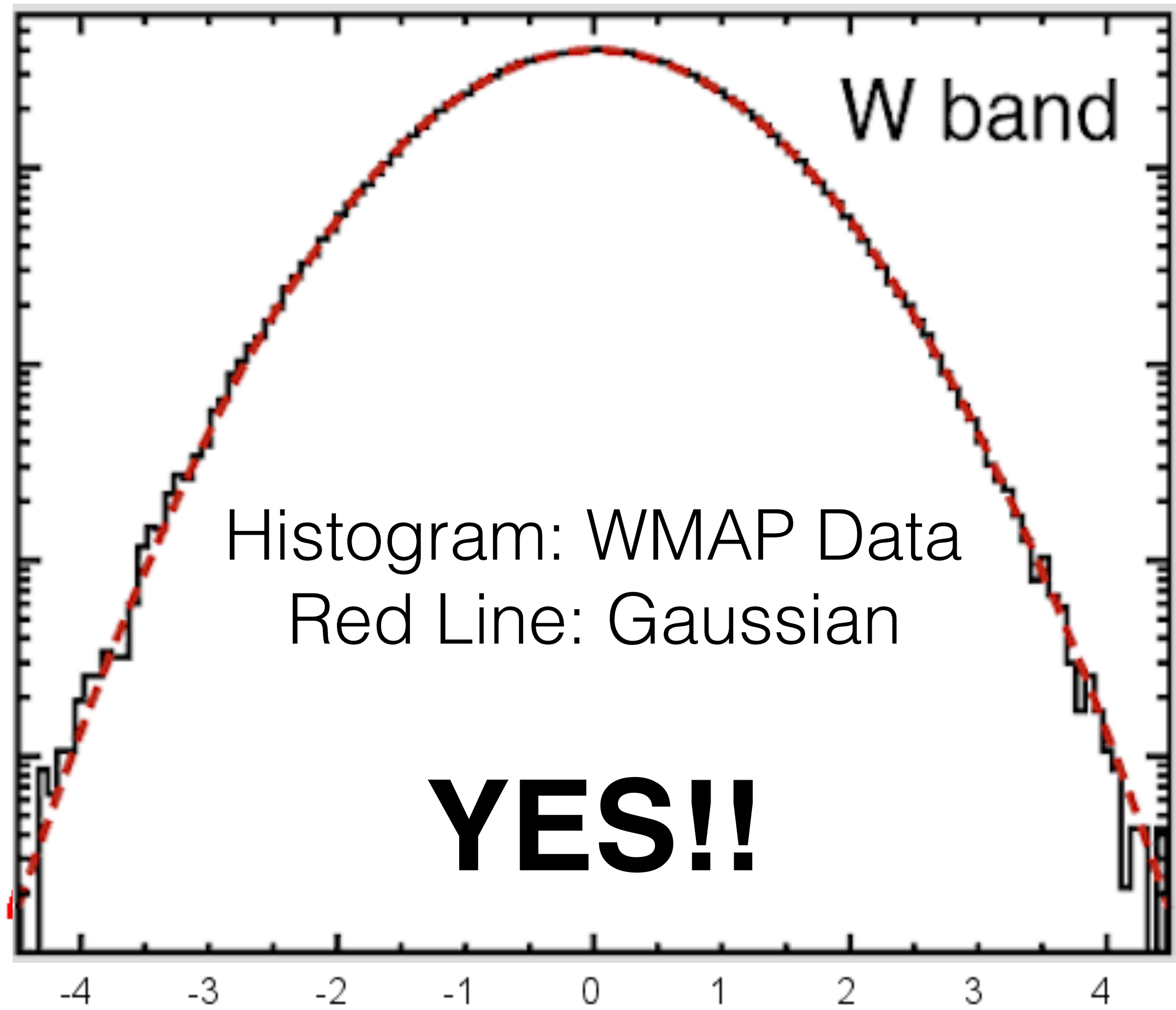
Fraction of the Number of Pixels
Having Those Temperatures

W band

Histogram: WMAP Data
Red Line: Gaussian

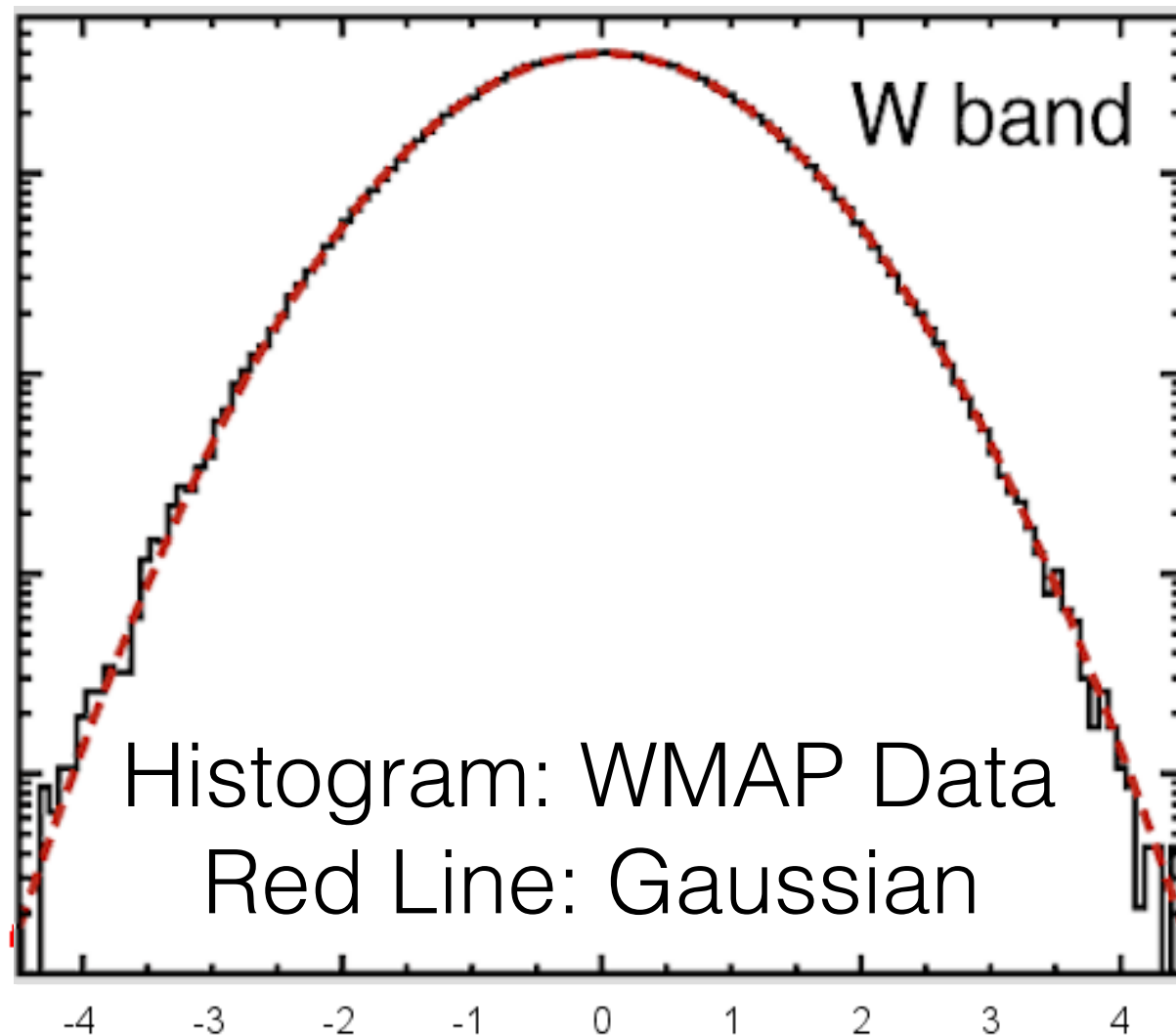
YES!!

$[\text{Values of Temperatures in the Sky Minus } 2.725 \text{ K}] / [\text{Root Mean Square}]$



Testing Gaussianity

Fraction of the Number of Pixels
Having Those Temperatures



[Values of Temperatures in the Sky Minus
2.725 K]/ [Root Mean Square]

Since a Gauss distribution is symmetric, it must yield a vanishing **3-point function**

$$\langle \delta T^3 \rangle \equiv \int_{-\infty}^{\infty} d\delta T \, P(\delta T) \delta T^3$$

More specifically, we measure this using temperatures at three different locations and average:

$$\langle \delta T(\hat{n}_1) \delta T(\hat{n}_2) \delta T(\hat{n}_3) \rangle$$

Non-Gaussianity:

A Powerful Test of Quantum Fluctuations

- The WMAP data show that the distribution of temperature fluctuations of CMB is **very precisely** Gaussian
 - with an upper bound on a deviation of **0.2%**
- With improved data provided by the Planck mission, the upper bound is now **0.03%**

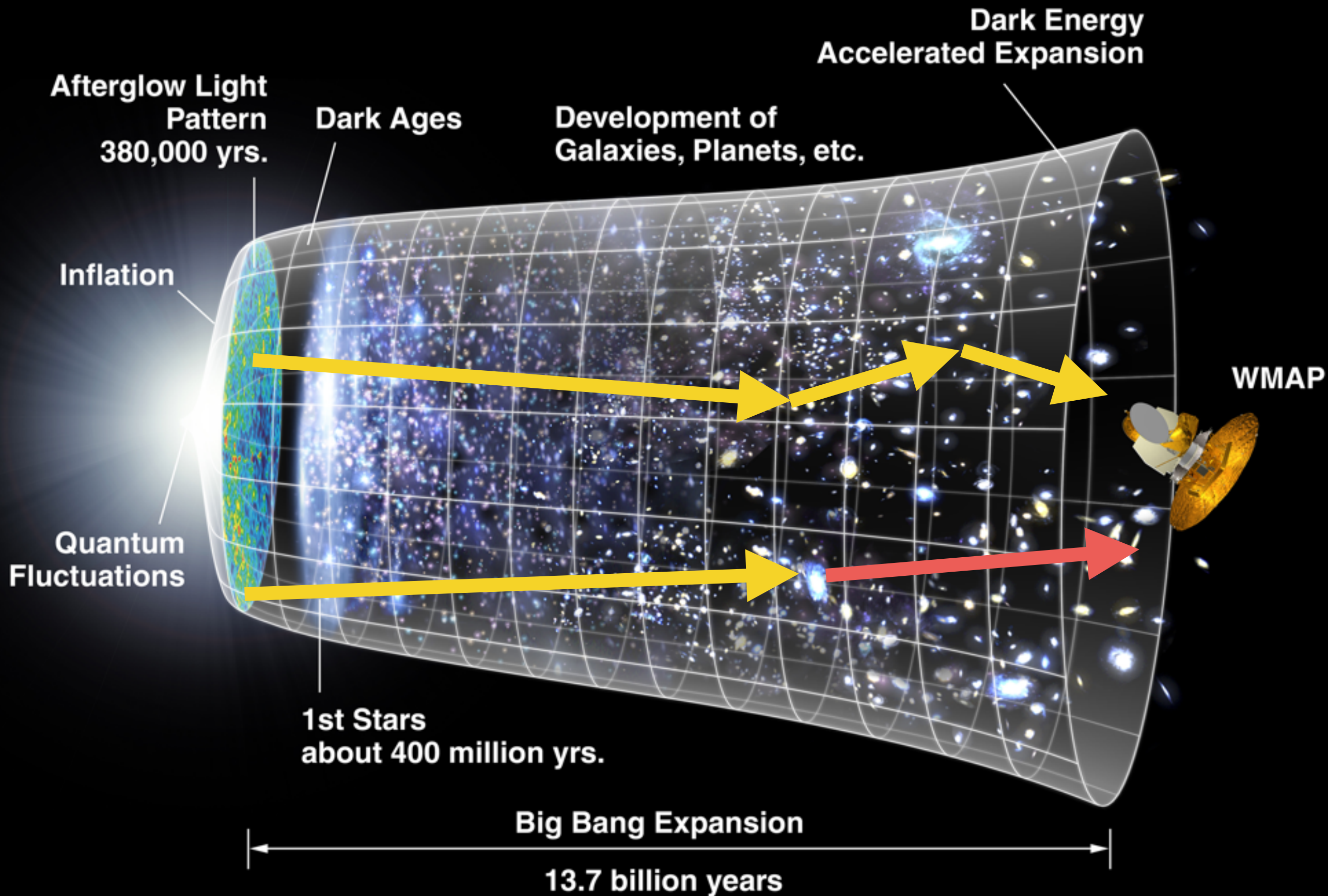
Secondary Anisotropies: Structure Formation seen in the CMB

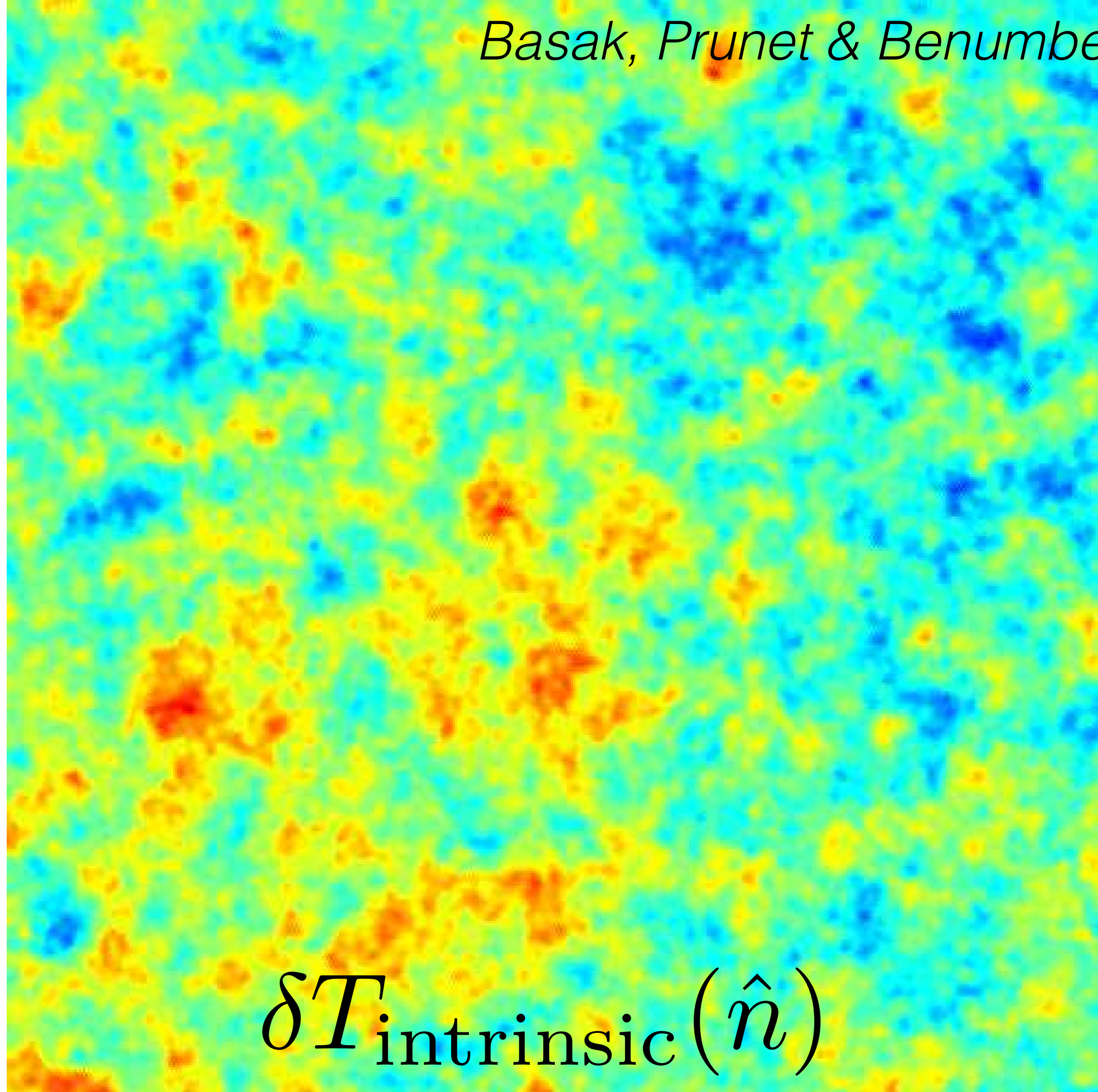
Gravitational Lensing

- Matter bends light of the CMB

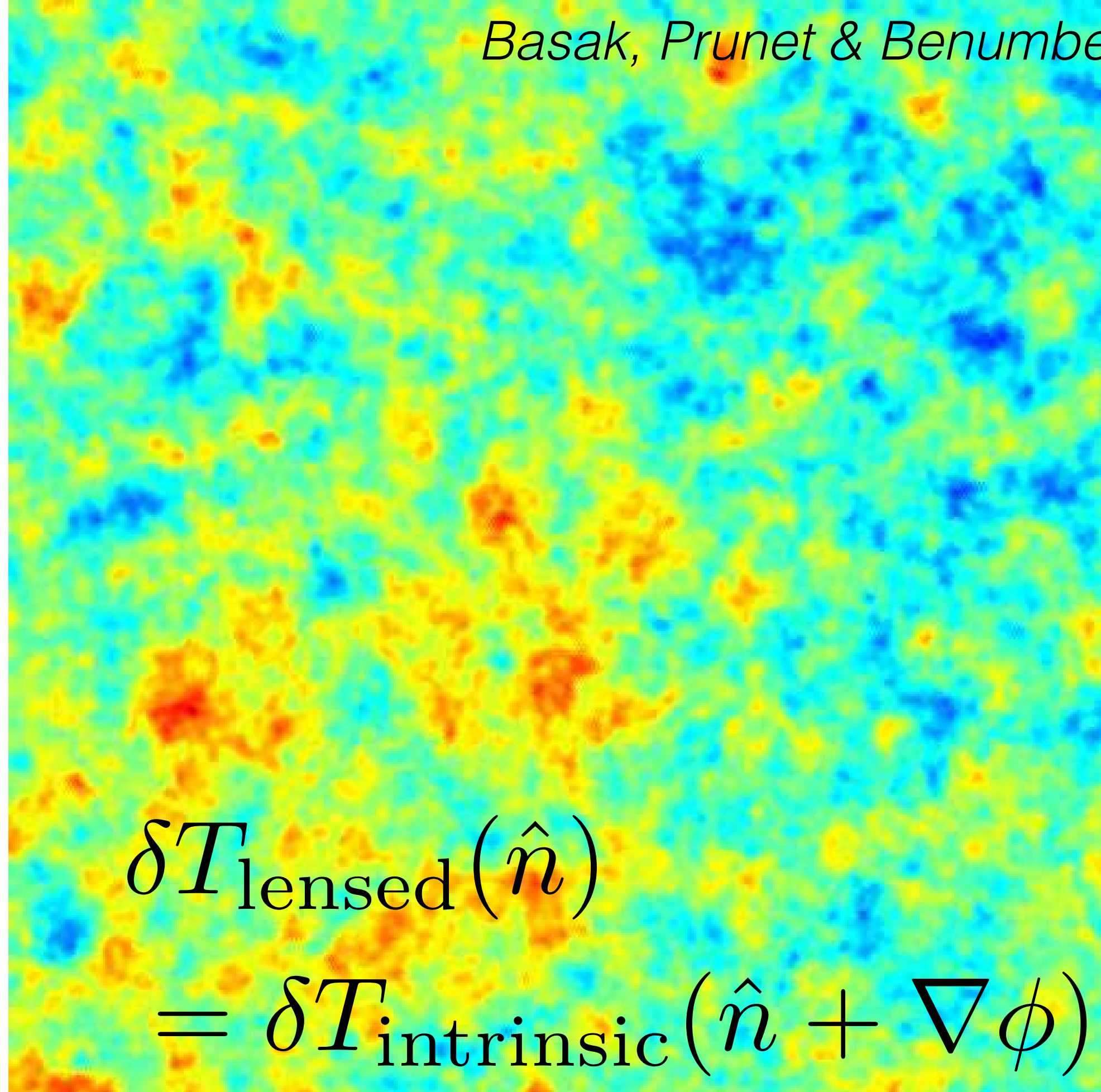
Sunyaev-Zel'dovich Effect

- Electrons in hot, collapsed gas up-scatter low-energy CMB photons, distorting the black-body spectrum of the CMB
- Both have been measured, providing the key insights into how the structures grew out of initial conditions.
Initial conditions to structure formation, using the CMB data only!

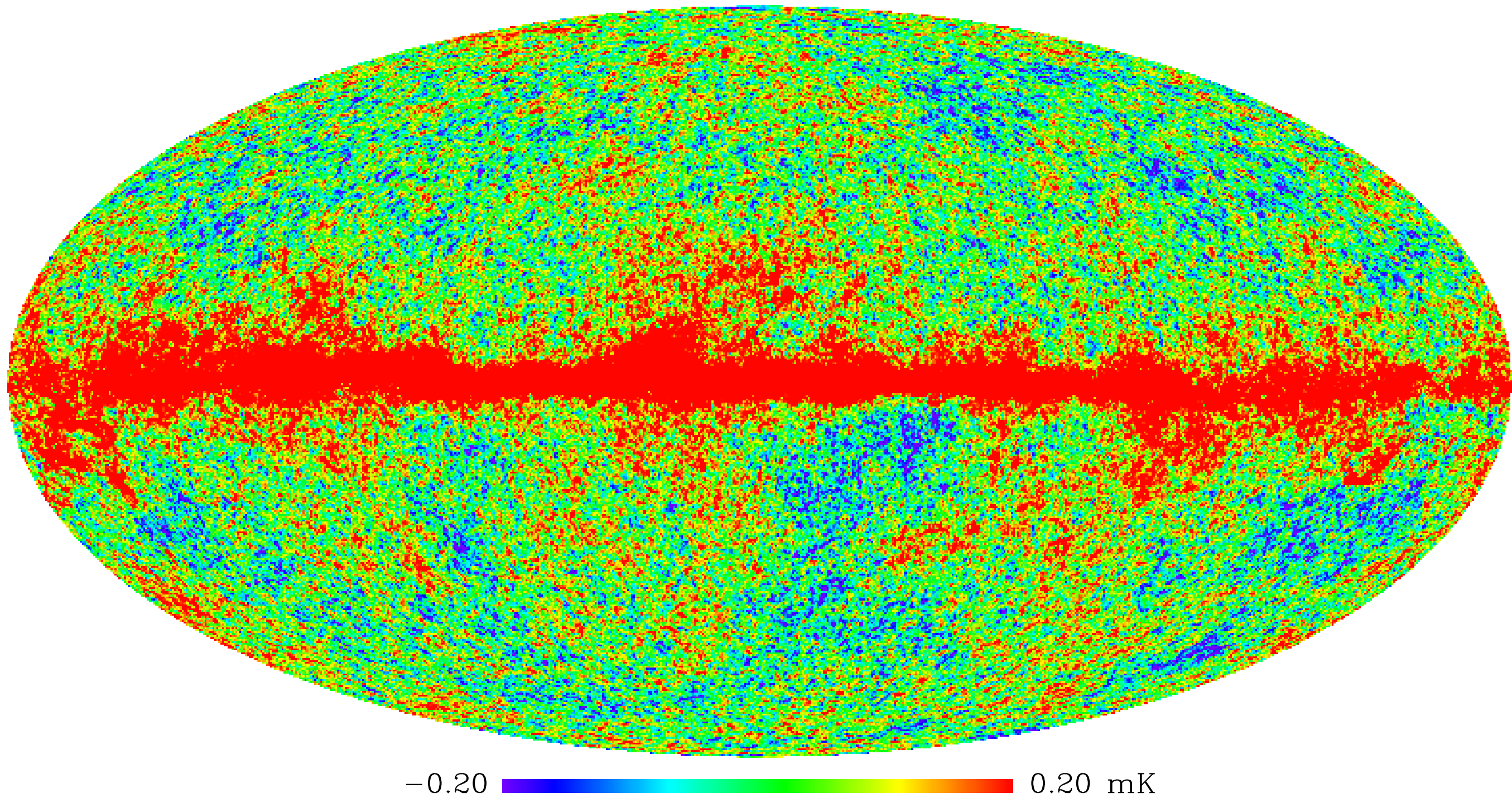


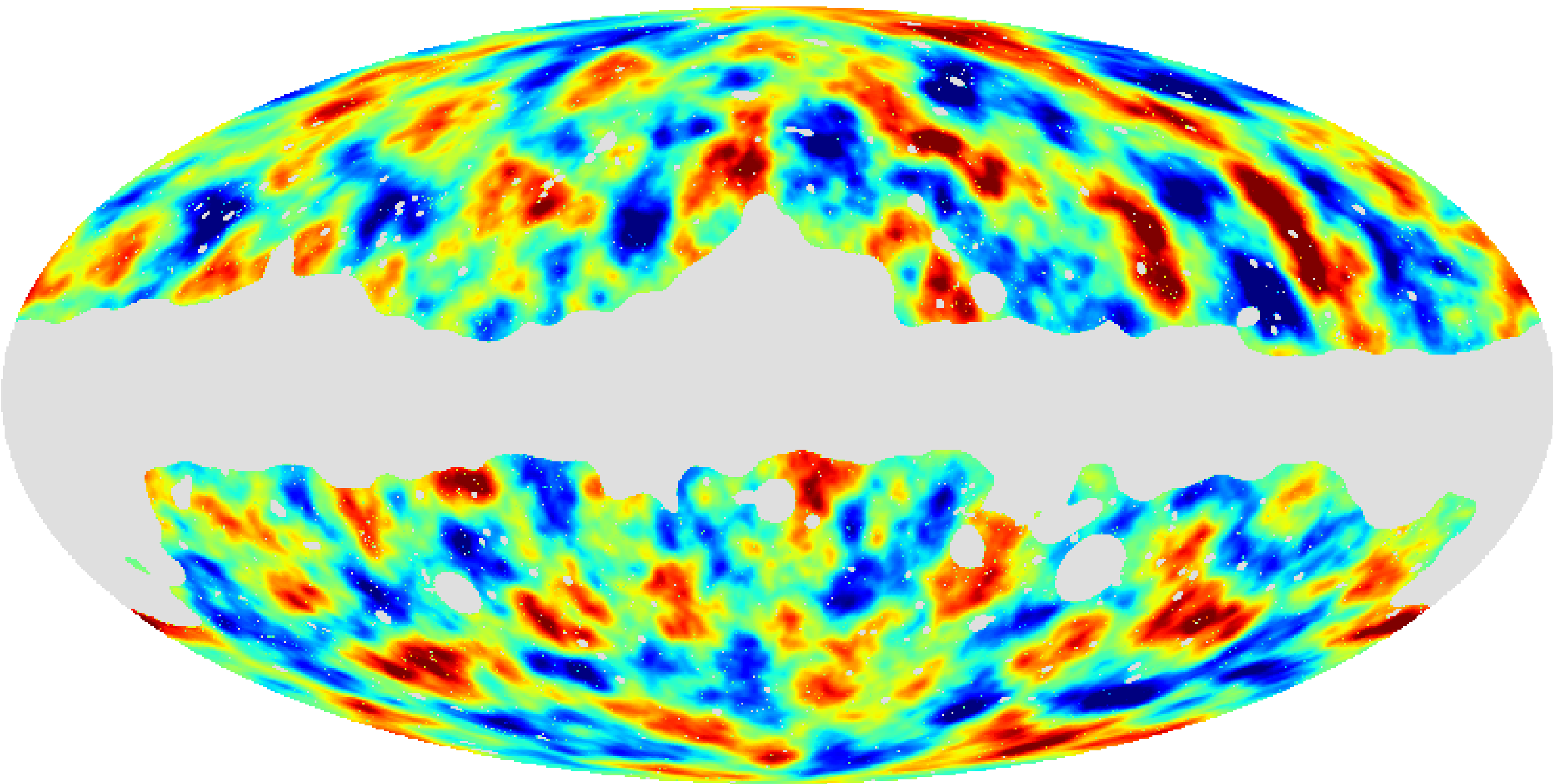


−605 605 μK

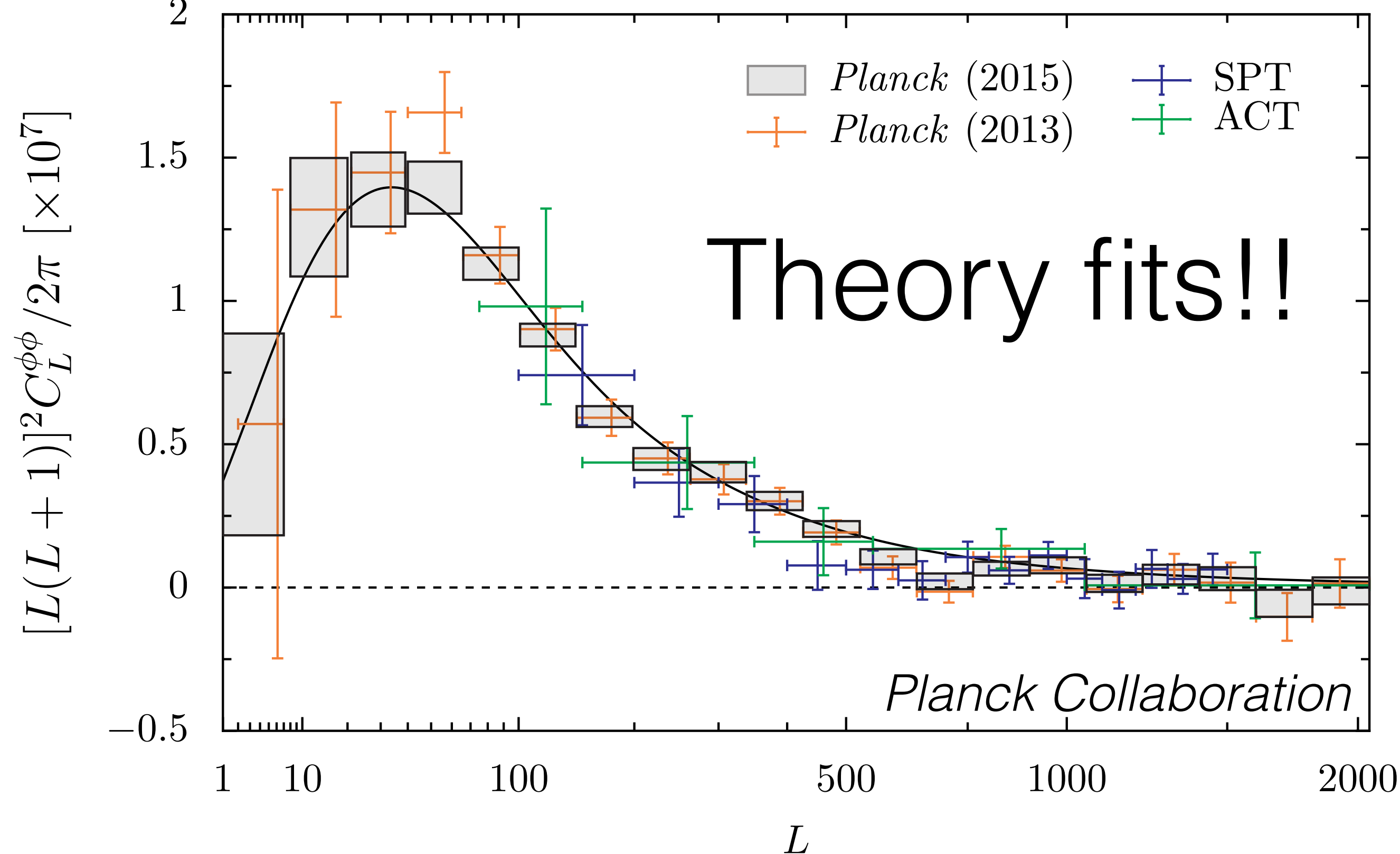


Planck 29–Month Map [100 GHz]





A full-sky lensing potential map!



...and our knowledge of the matter distribution improves

$$\Omega_m = 0.315 \pm 0.013$$

$$\sigma_8 = 0.829 \pm 0.014$$

Adding the
lensing info

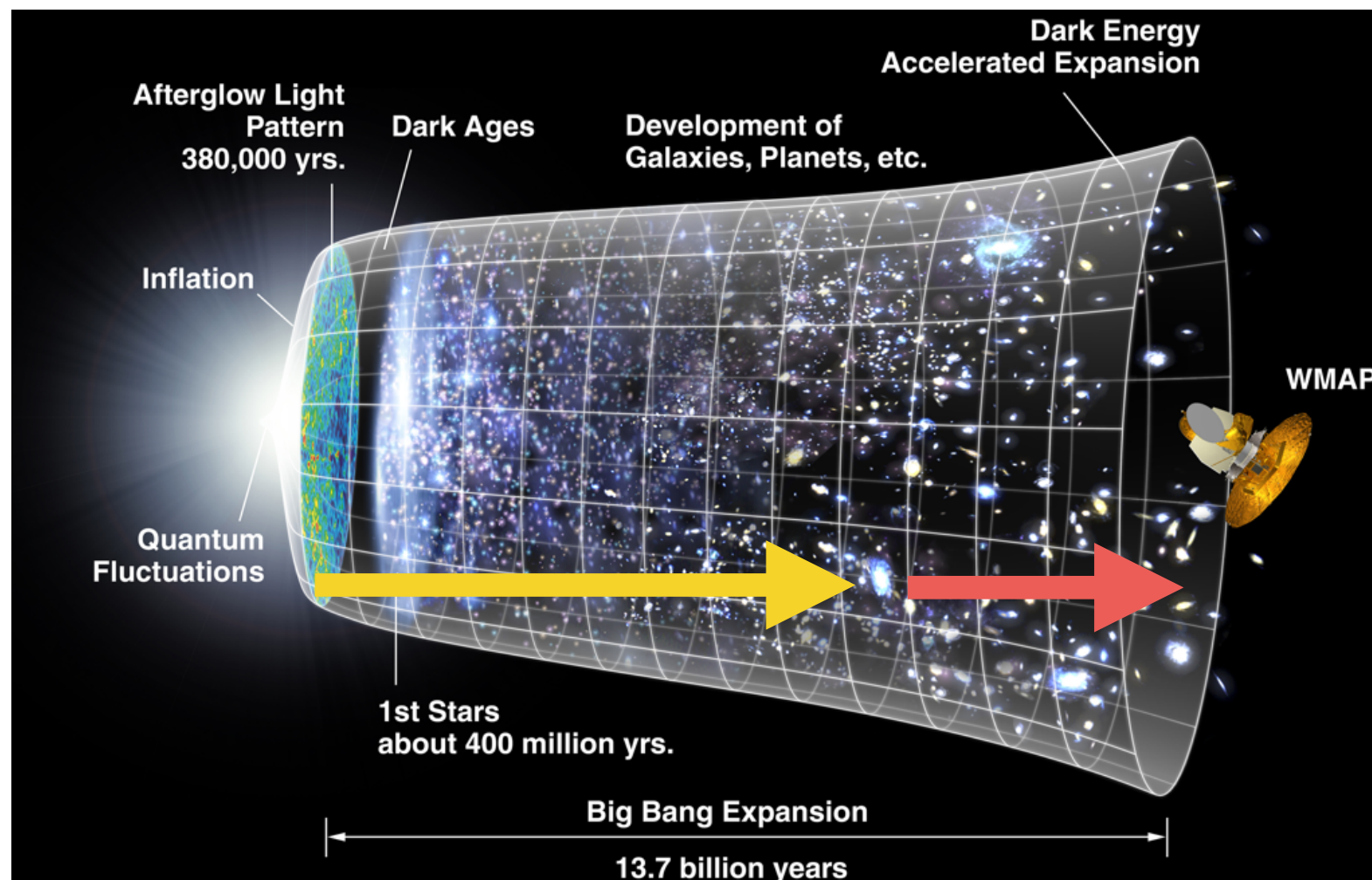
$$\Omega_m = 0.308 \pm 0.012$$

$$\sigma_8 = 0.8149 \pm 0.0093$$

The Sunyaev-Zel'dovich Effect

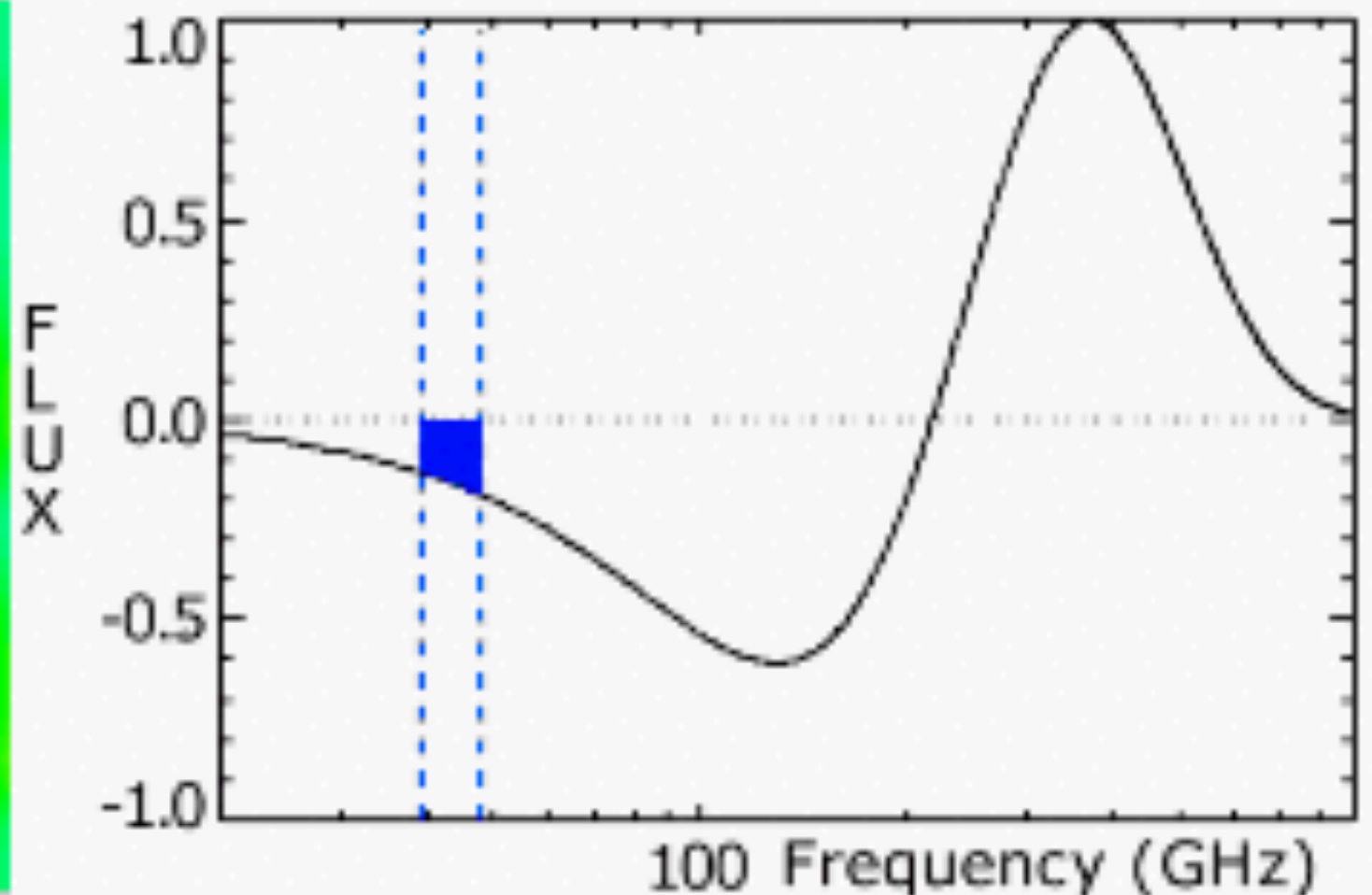
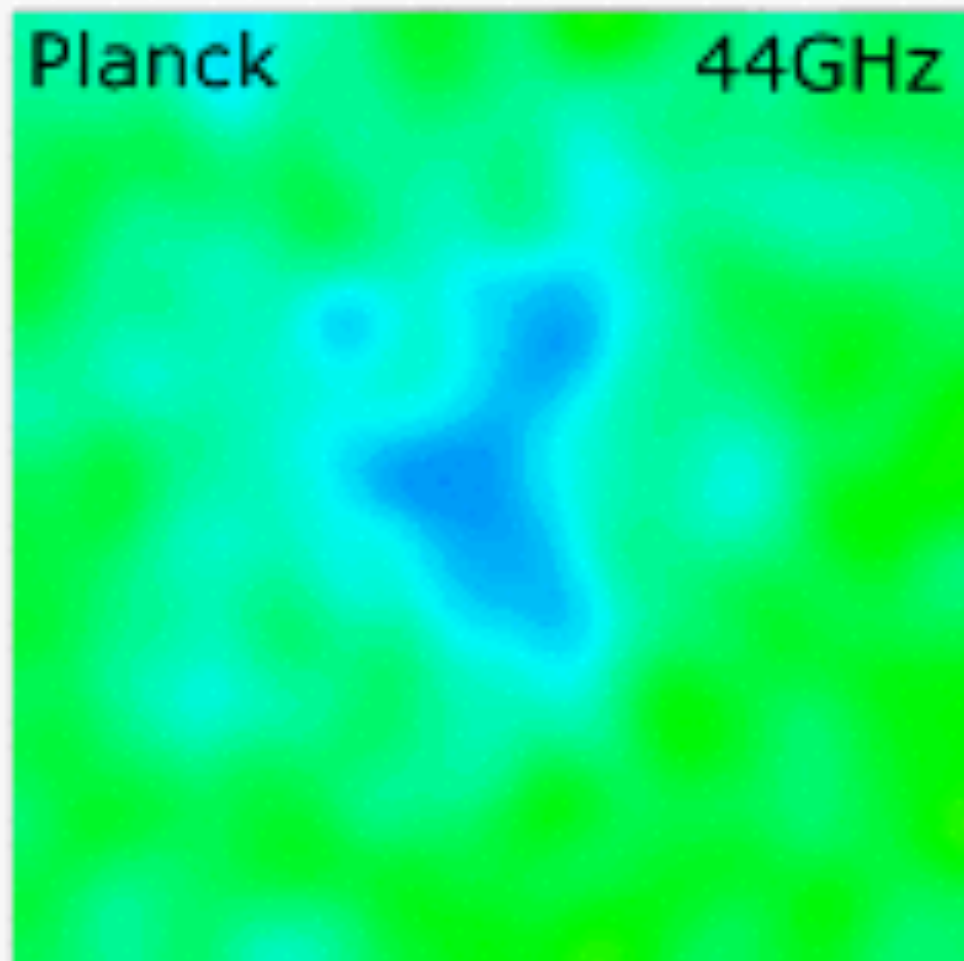
- The “thermal” SZ effect (in the non-relativistic limit) enables us to map thermal pressure in the universe

$$Y_{\text{tSZ}}(\boldsymbol{\theta}) = \frac{k_B \sigma_T}{m_e c^2} \int dl n_e(\boldsymbol{\theta}, l) T(\boldsymbol{\theta}, l)$$



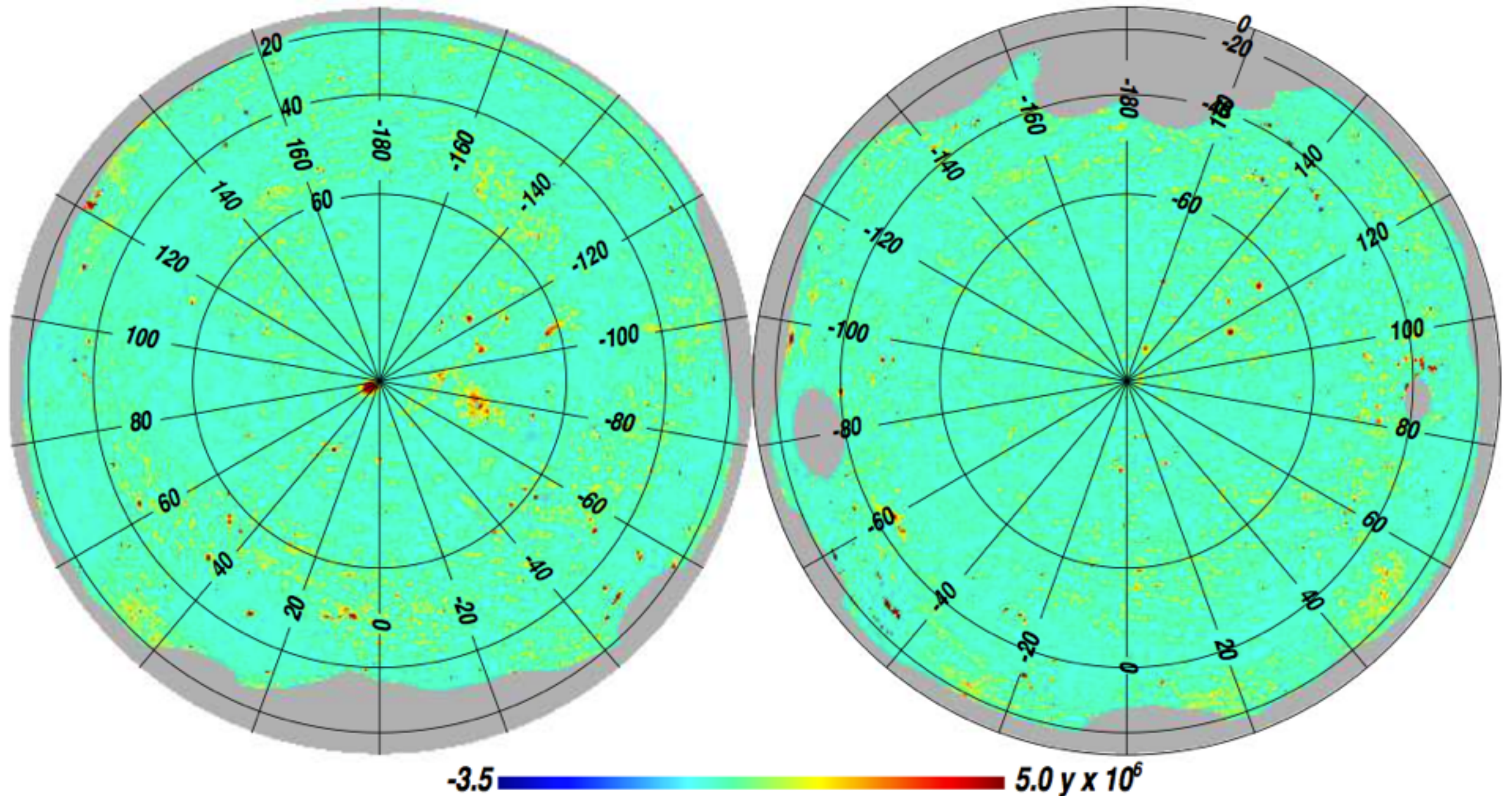
Thermal SZ (tSZ) Effect

- The unique frequency dependence of tSZ: we can make a map of Y_{tSZ}

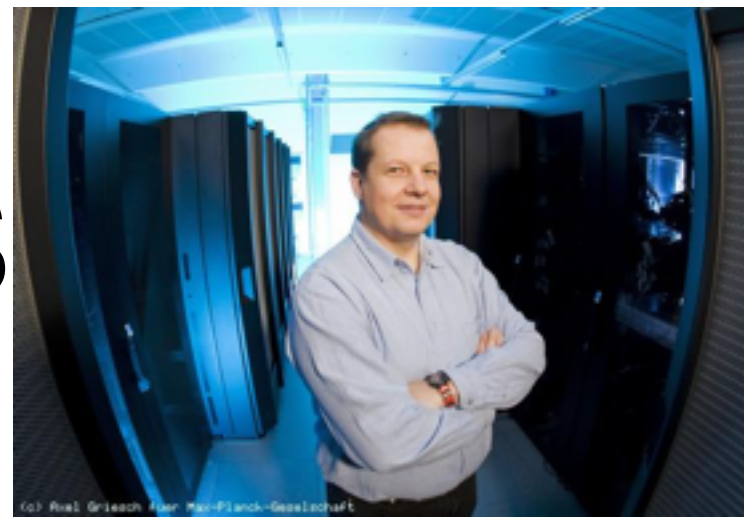


Full-sky Thermal Pressure Map

North Galactic Pole *MILCA tSZ map* South Galactic Pole



We can simulate this



Klaus Dolag (MPA/LMU)

arXiv:1509.05134 [accepted for publication in MNRAS]

SZ effects in the Magneticum Pathfinder Simulation: Comparison with the Planck, SPT, and ACT results

K. Dolag^{1,2*}, E. Komatsu^{2,3} and R. Sunyaev^{2,4}

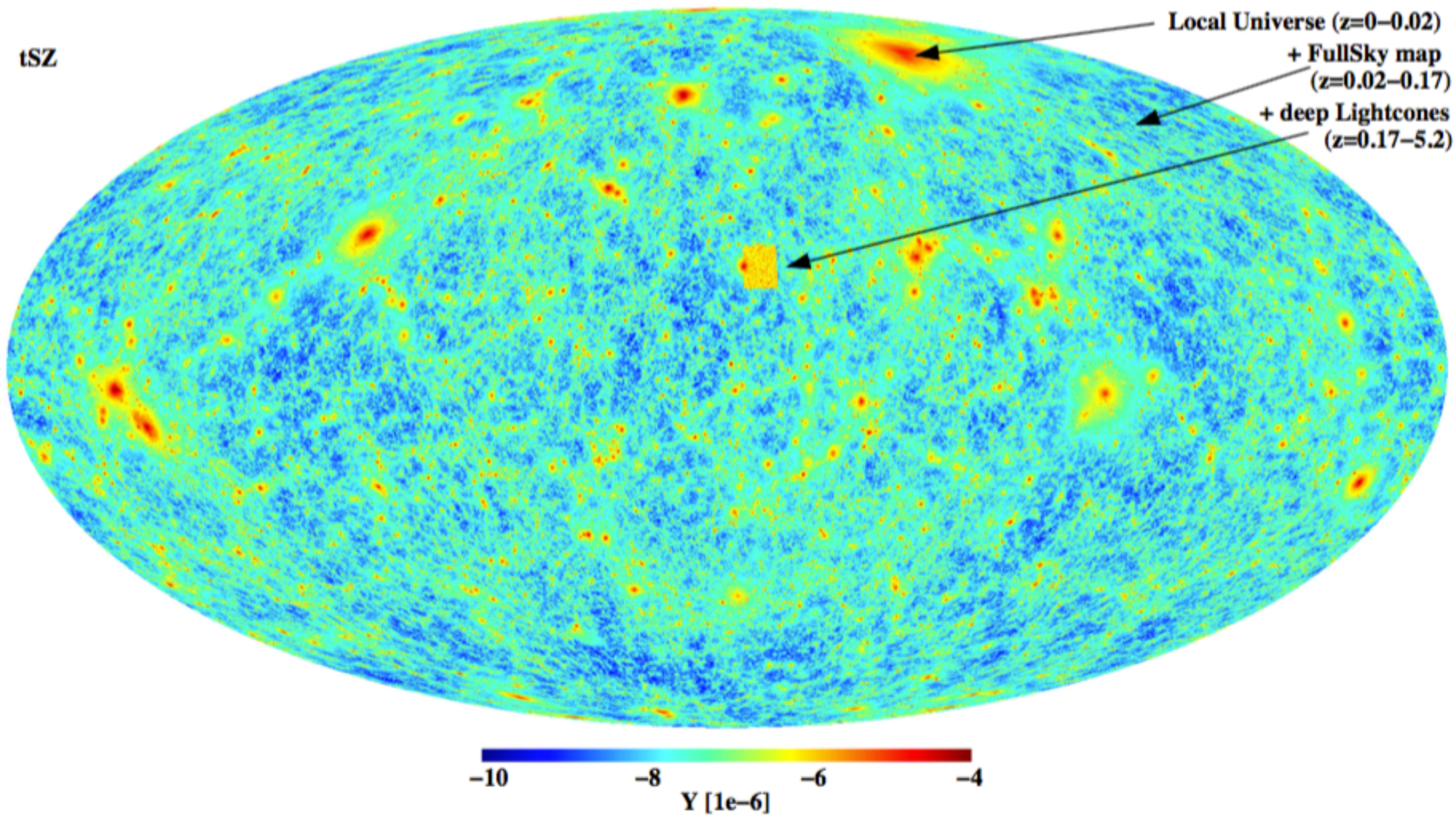
¹ *University Observatory Munich, Scheinerstr. 1, 81679 Munich, Germany*

² *Max-Planck-Institut für Astrophysik, Karl-Schwarzschild Strasse 1, 85748 Garching, Germany*

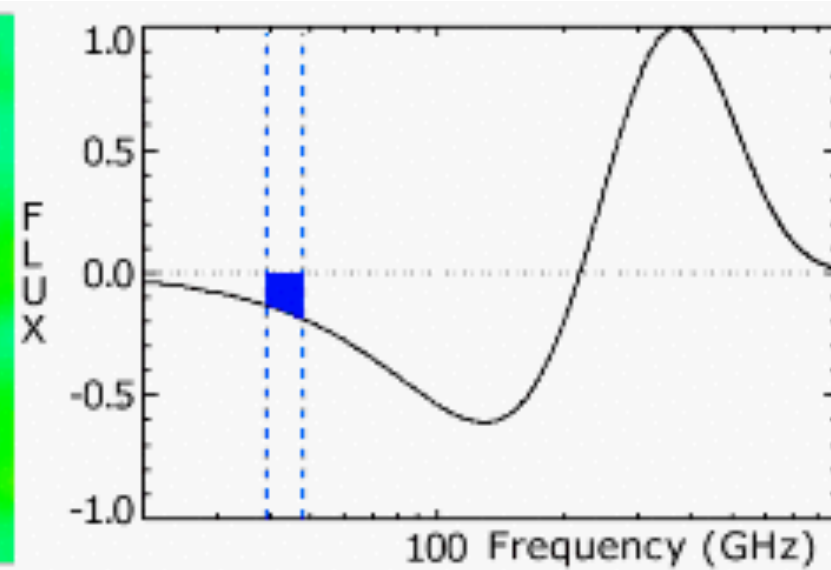
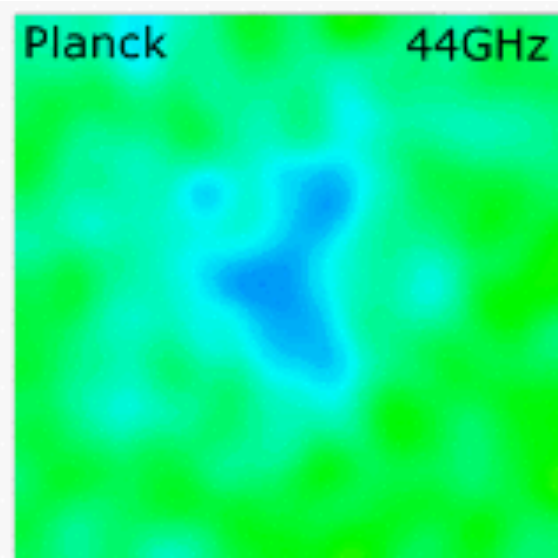
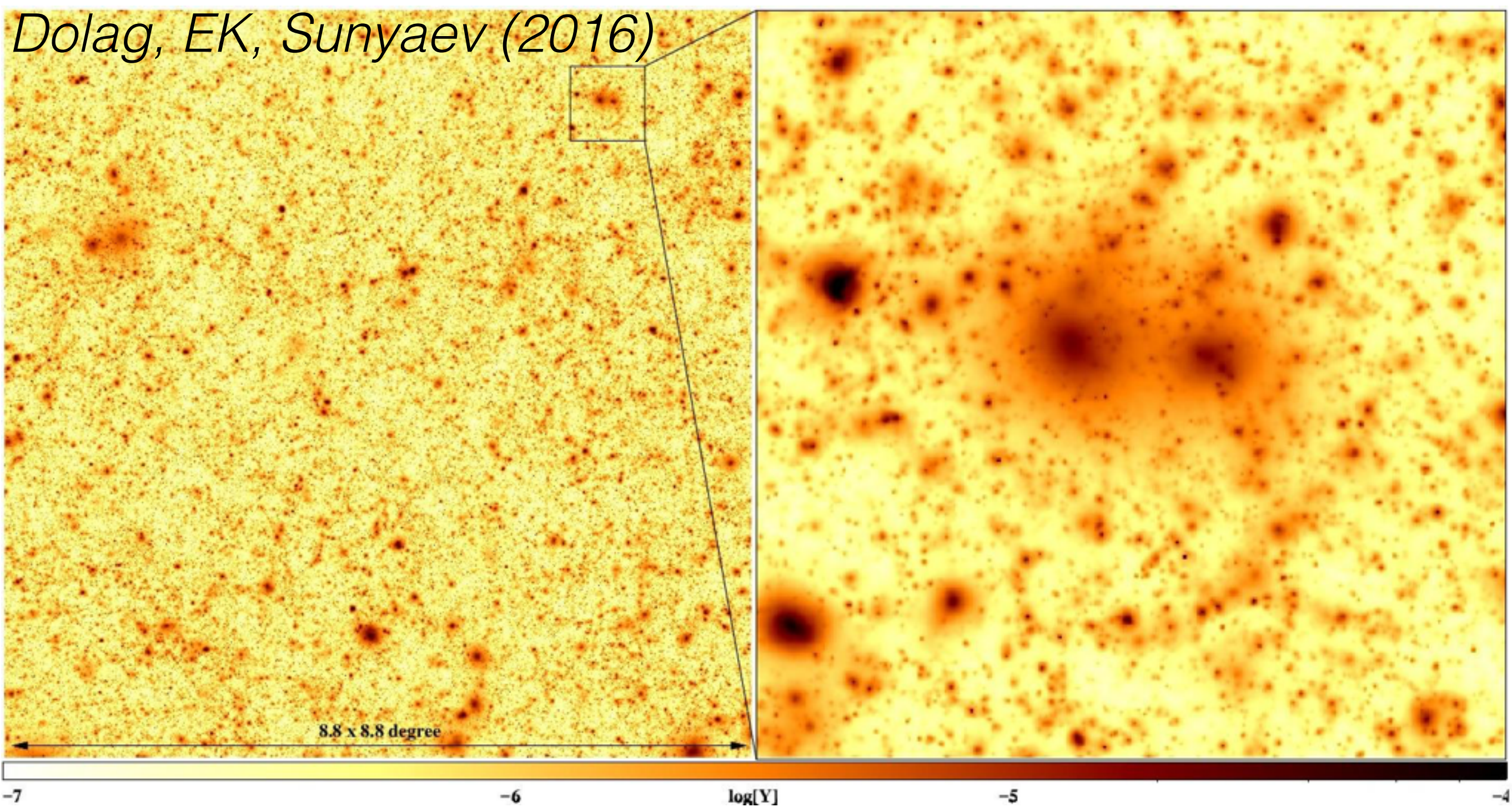
³ *Kavli Institute for the Physics and Mathematics of the Universe (Kavli IPMU, WPI), Todai Institutes for Advanced Study, the University of Tokyo, Kashiwa 277-8583, Japan*

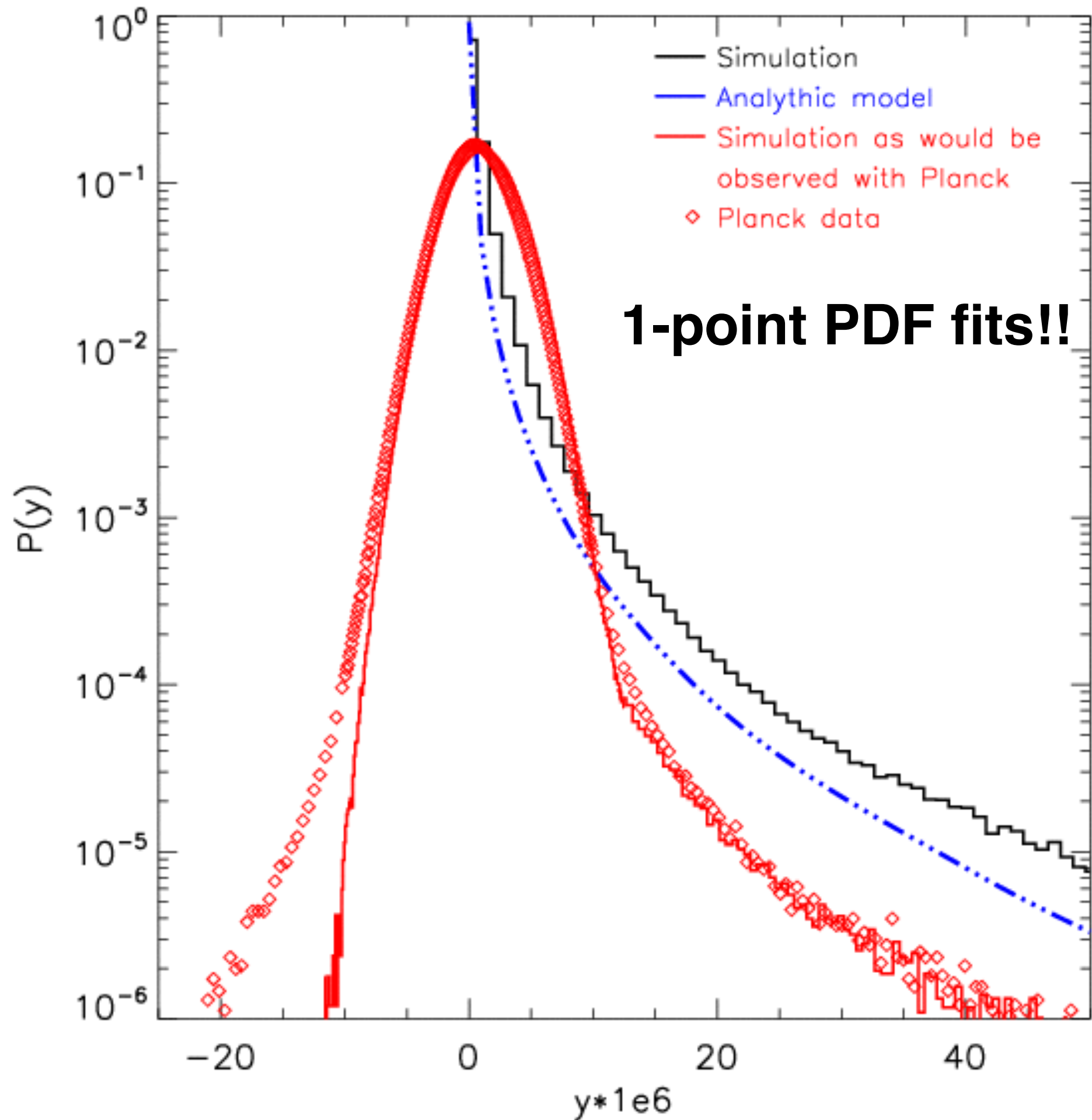
⁴ *Space Research Institute (IKI), Russian Academy of Sciences, Profsoyuznaya str. 84/32, Moscow, 117997 Russia*

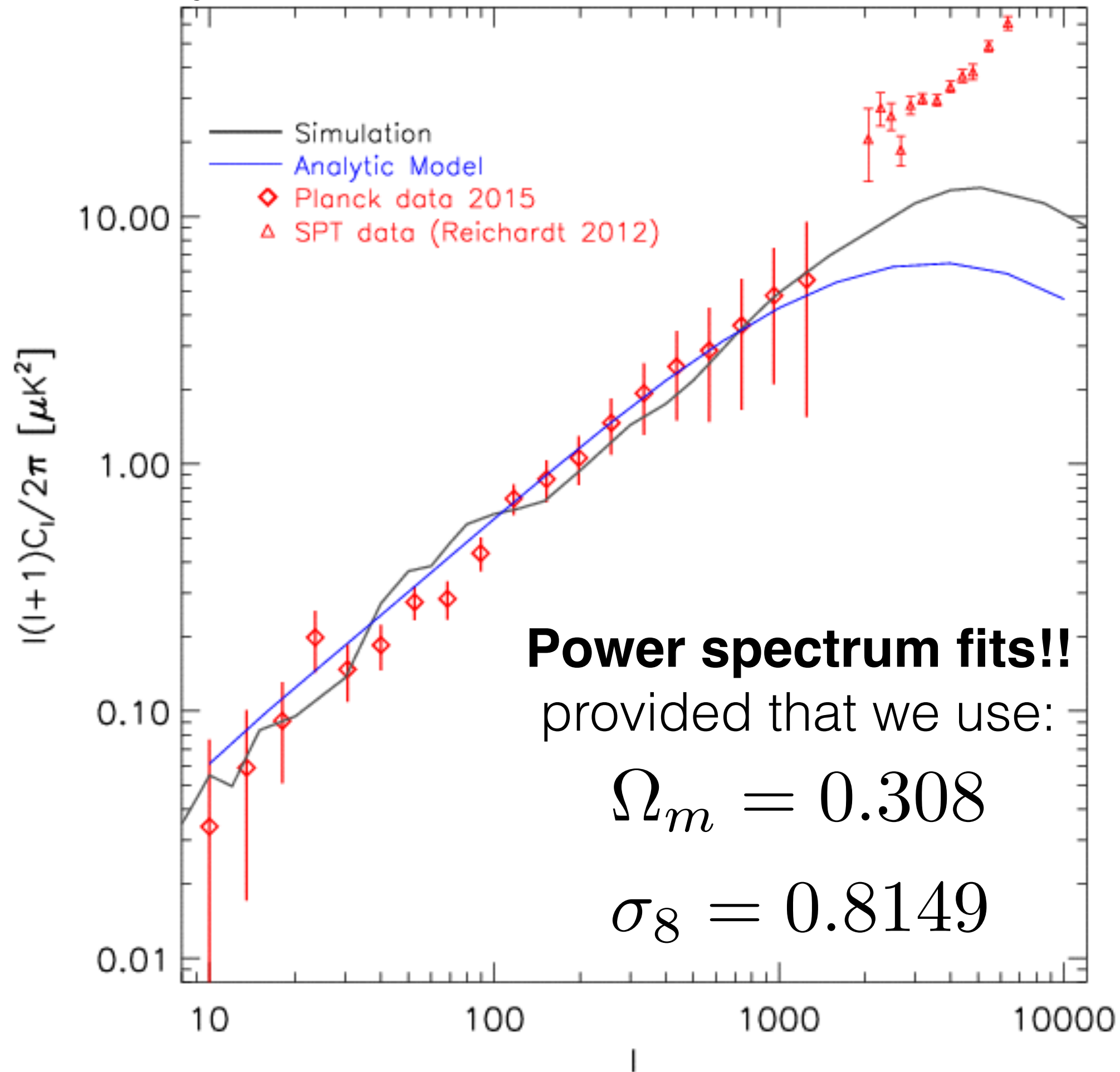
- Volume: $(896 \text{ Mpc}/h)^3$
- Cosmological hydro (P-GADGET3) with star formation and AGN feed back
- 2×1526^3 particles ($m_{\text{DM}}=7.5 \times 10^8 \text{ M}_{\text{sun}}/h$)

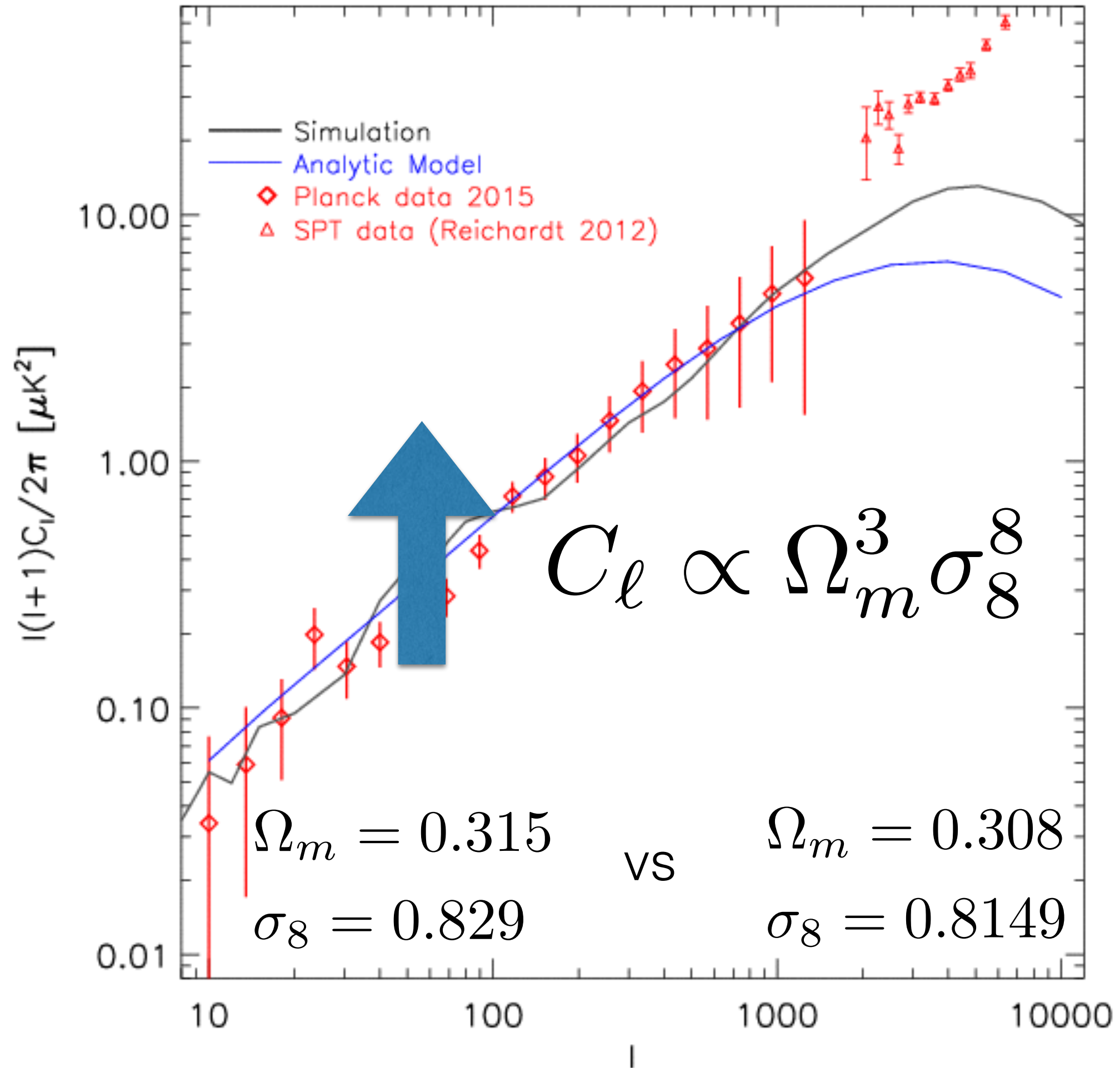


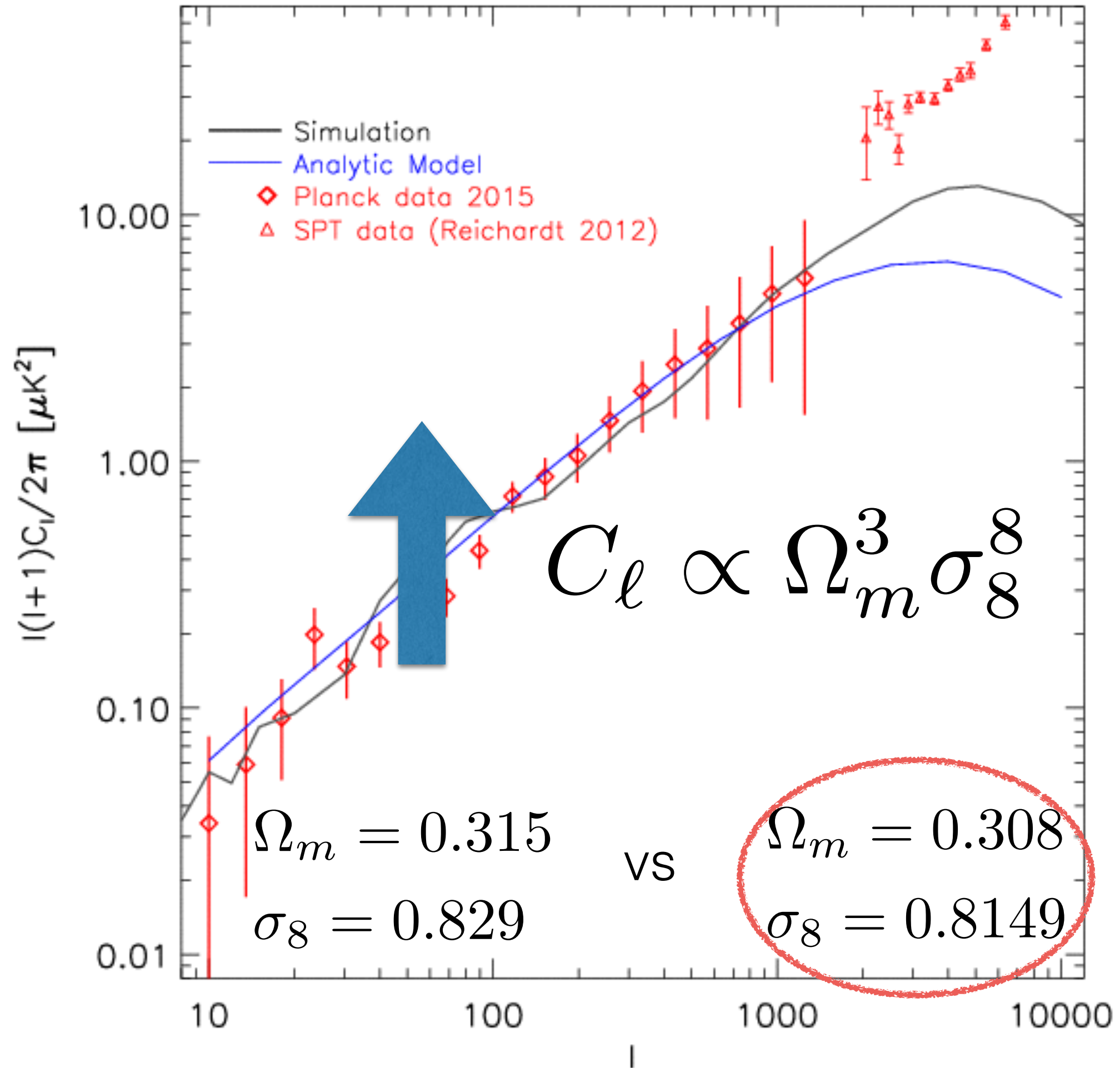
$$Y_{\text{tSZ}}(\boldsymbol{\theta}) = \frac{k_B \sigma_T}{m_e c^2} \int dl n_e(\boldsymbol{\theta}, l) T(\boldsymbol{\theta}, l)$$

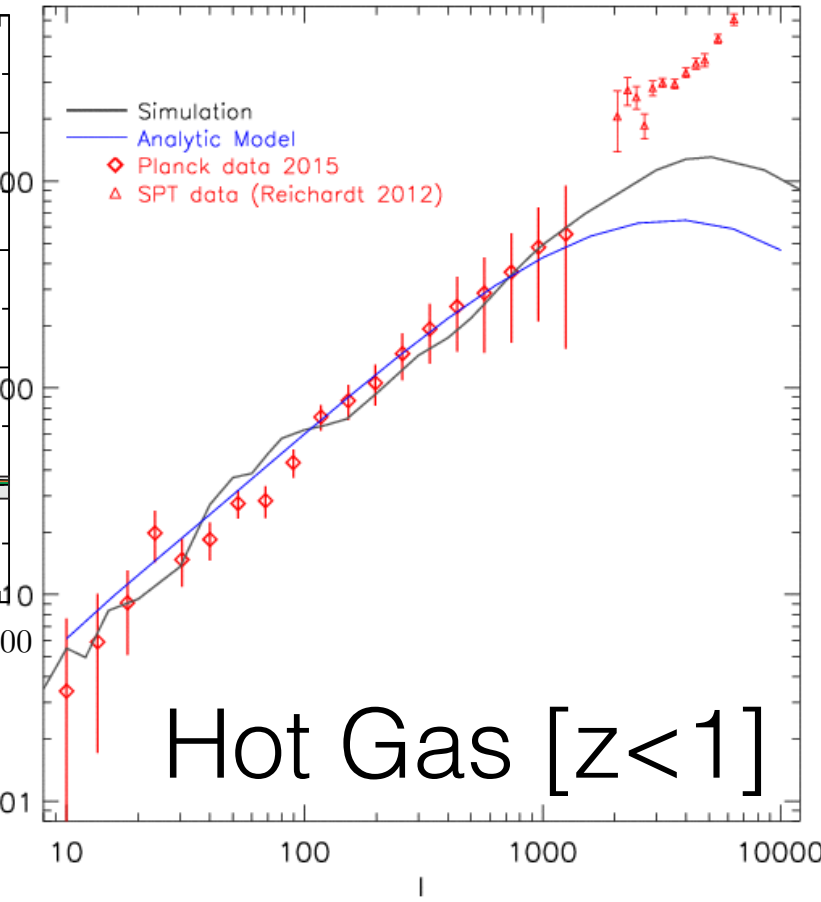
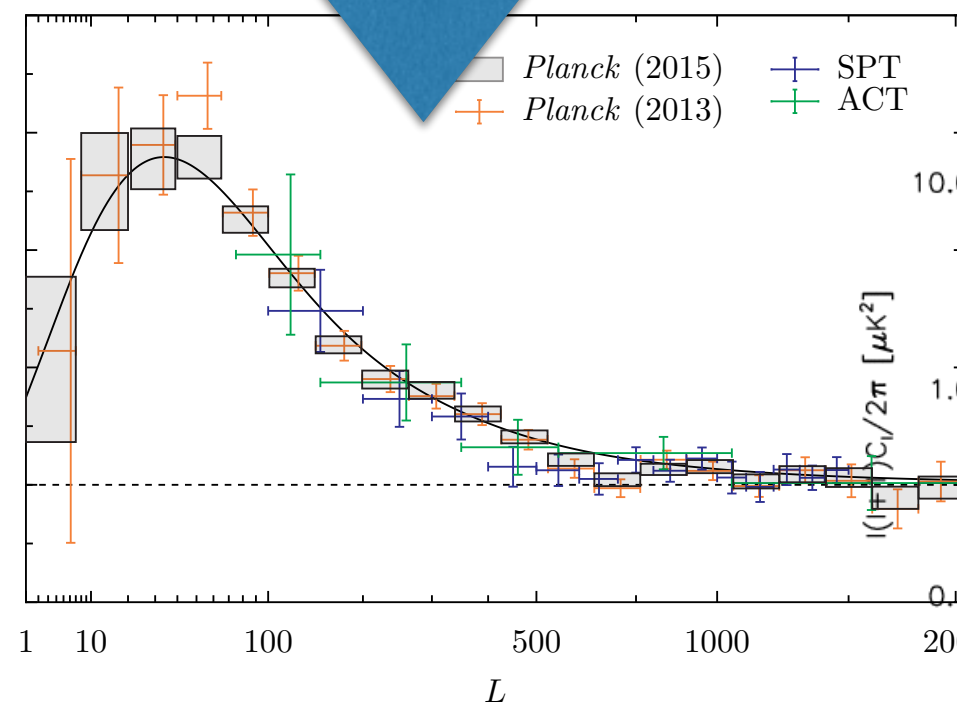
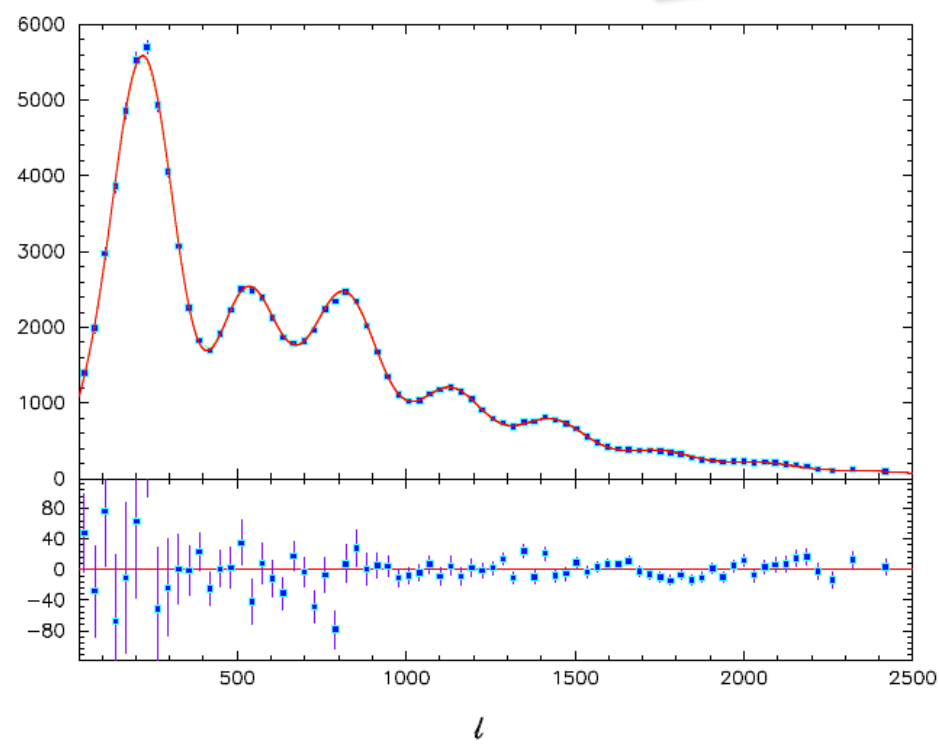
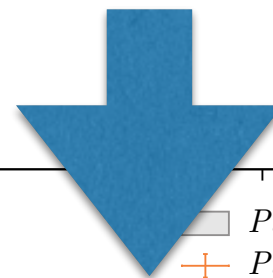
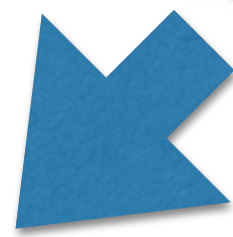
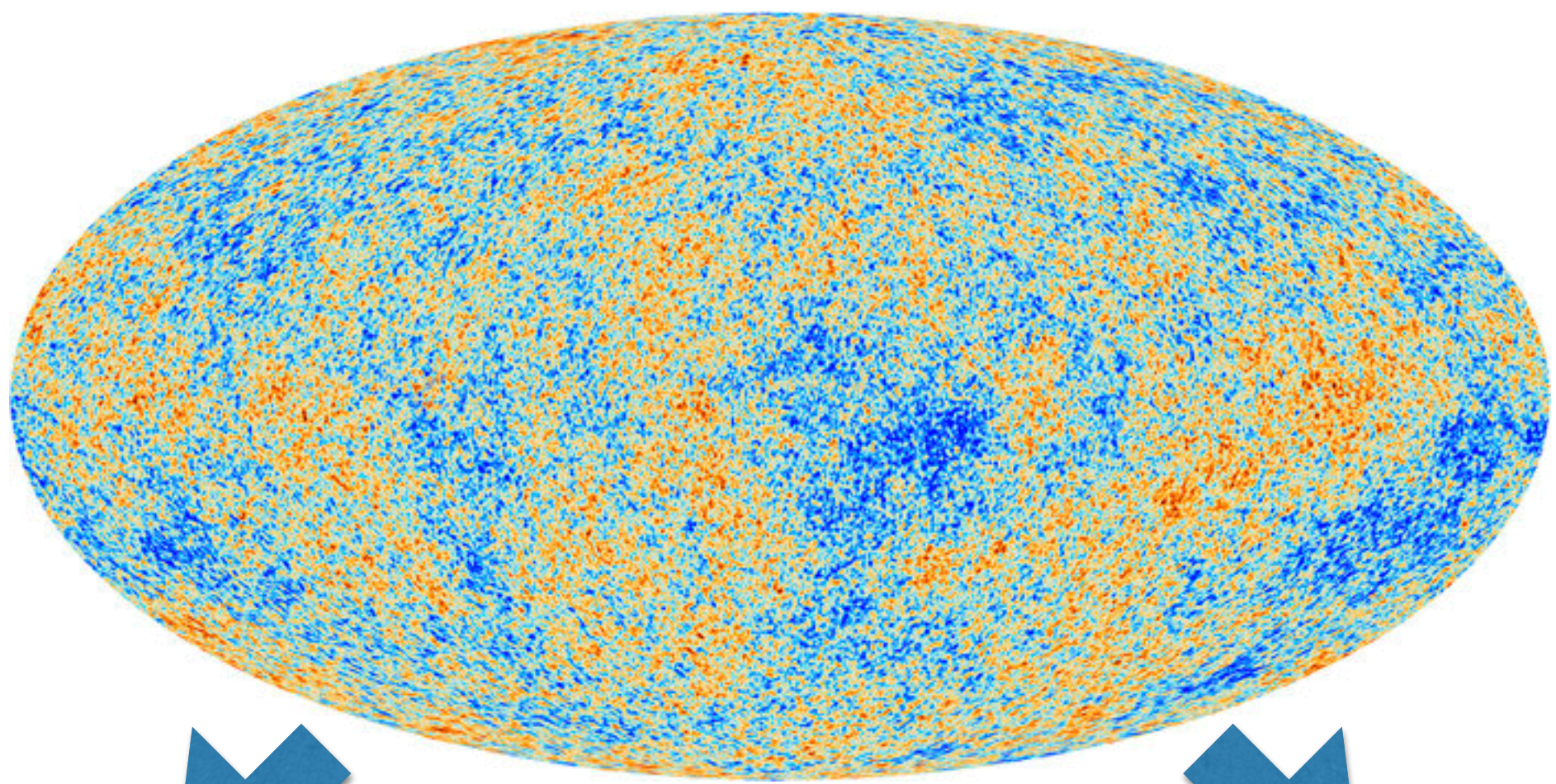








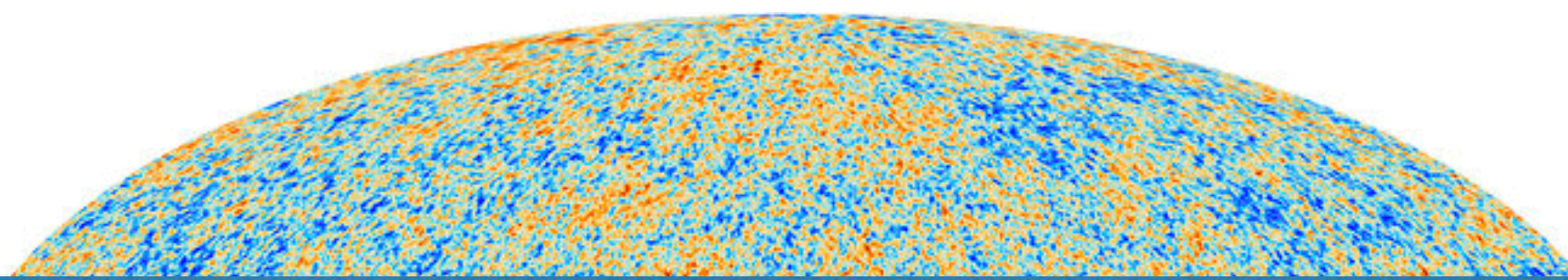




Initial Condition

Dark Matter [$z \sim 2$]

Hot Gas [$z < 1$]

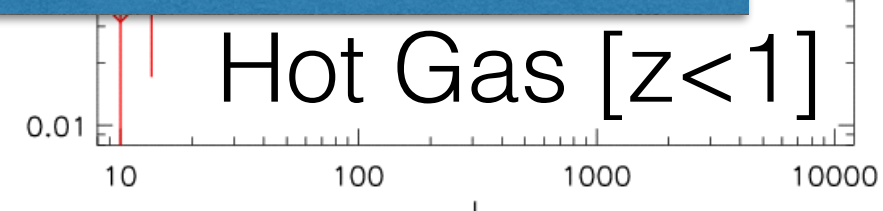


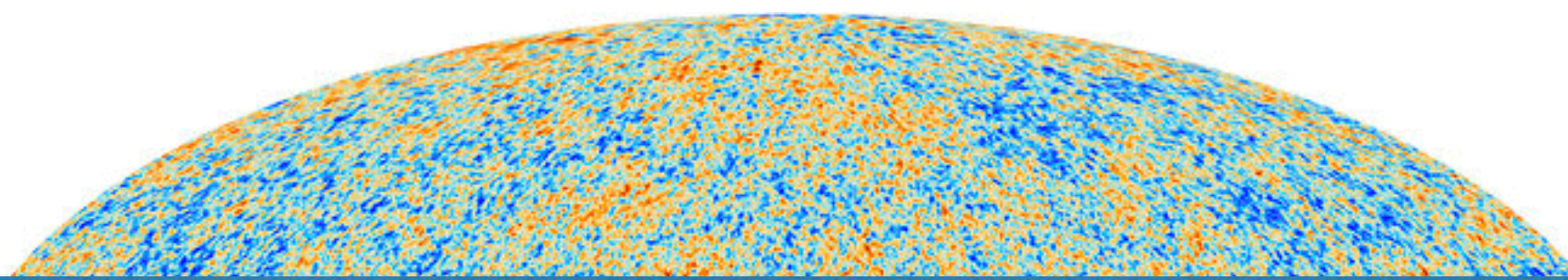
Standard Λ CDM Model, starting with
inflation producing adiabatic,
Gaussian, isotropic, $n_s < 1$ primordial
fluctuations fit all the data from the
**initial condition to structure
formation!**

These results are all solely based on the
microwave background data

Initial Condition

Dark Matter [$z \sim 2$]

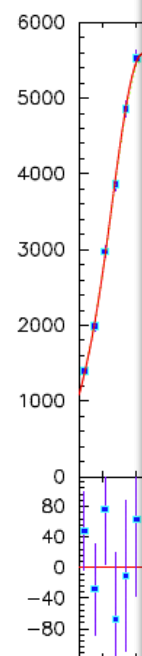




Standard Λ CDM Model, starting with
inflation producing adiabatic,

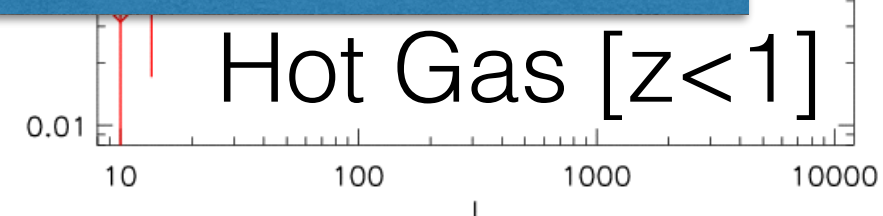
But, why not looking at
the large-scale
structure data directly?

These results are all solely based on the
microwave background data

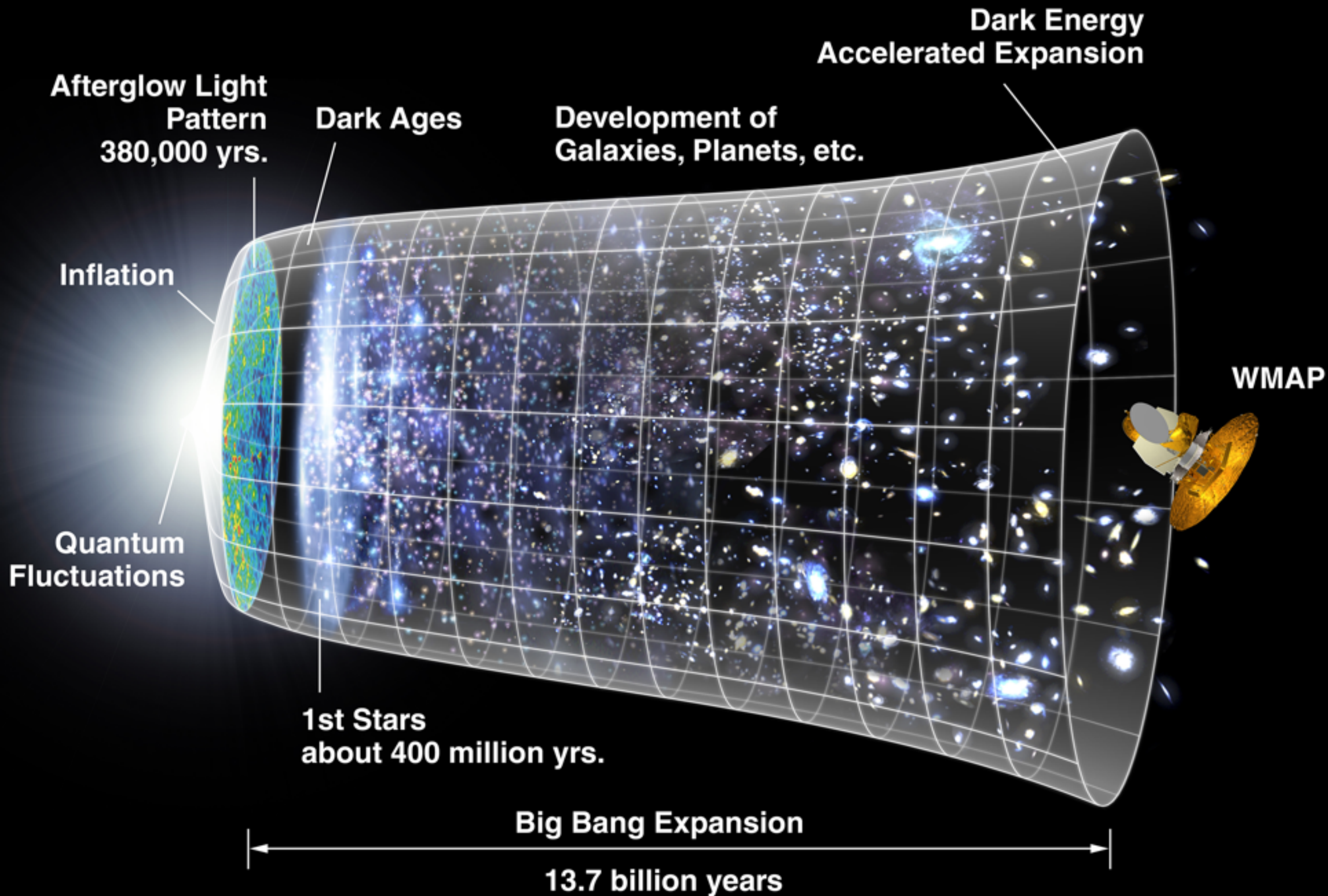


Initial Condition

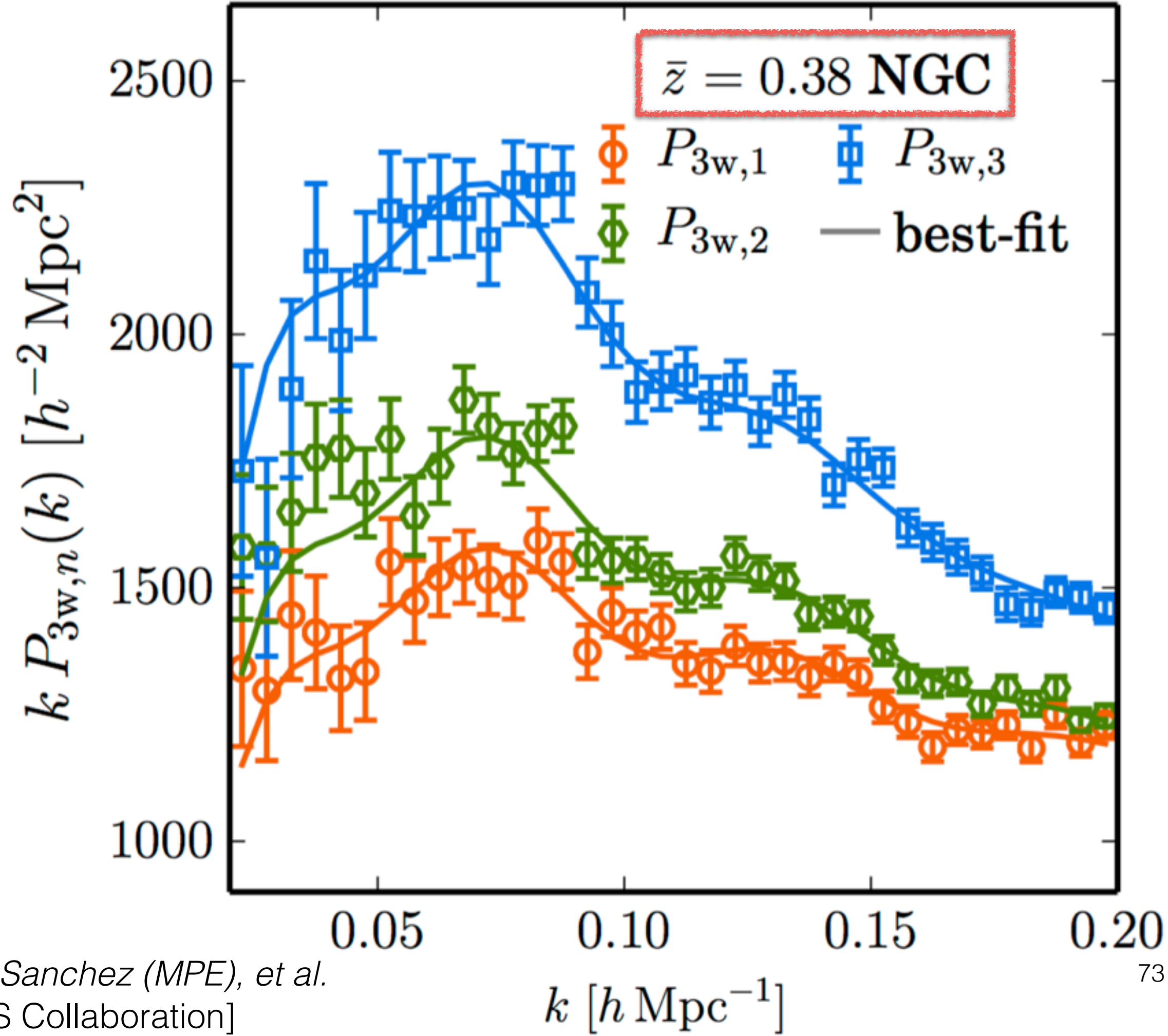
Dark Matter [$z \sim 2$]

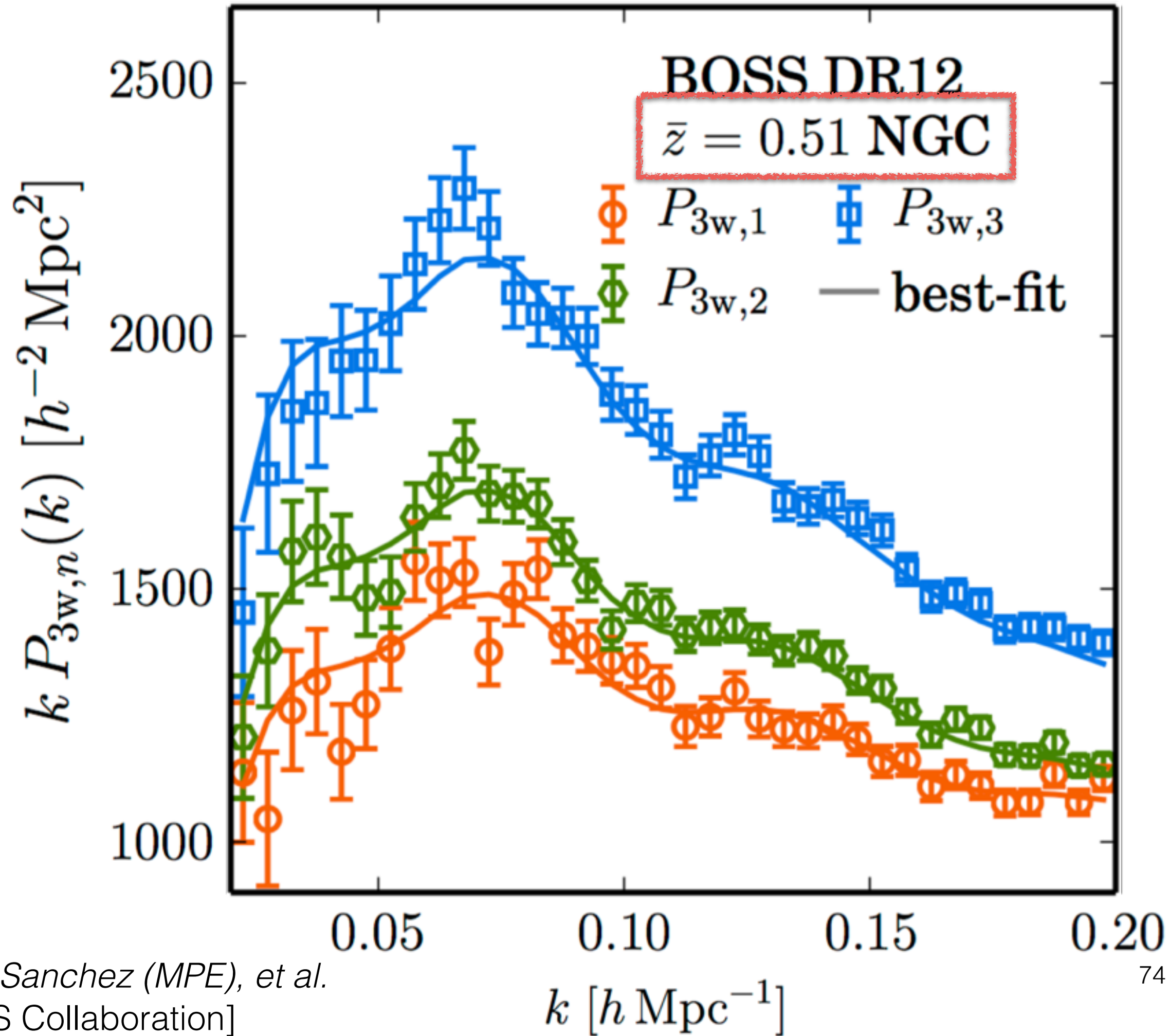


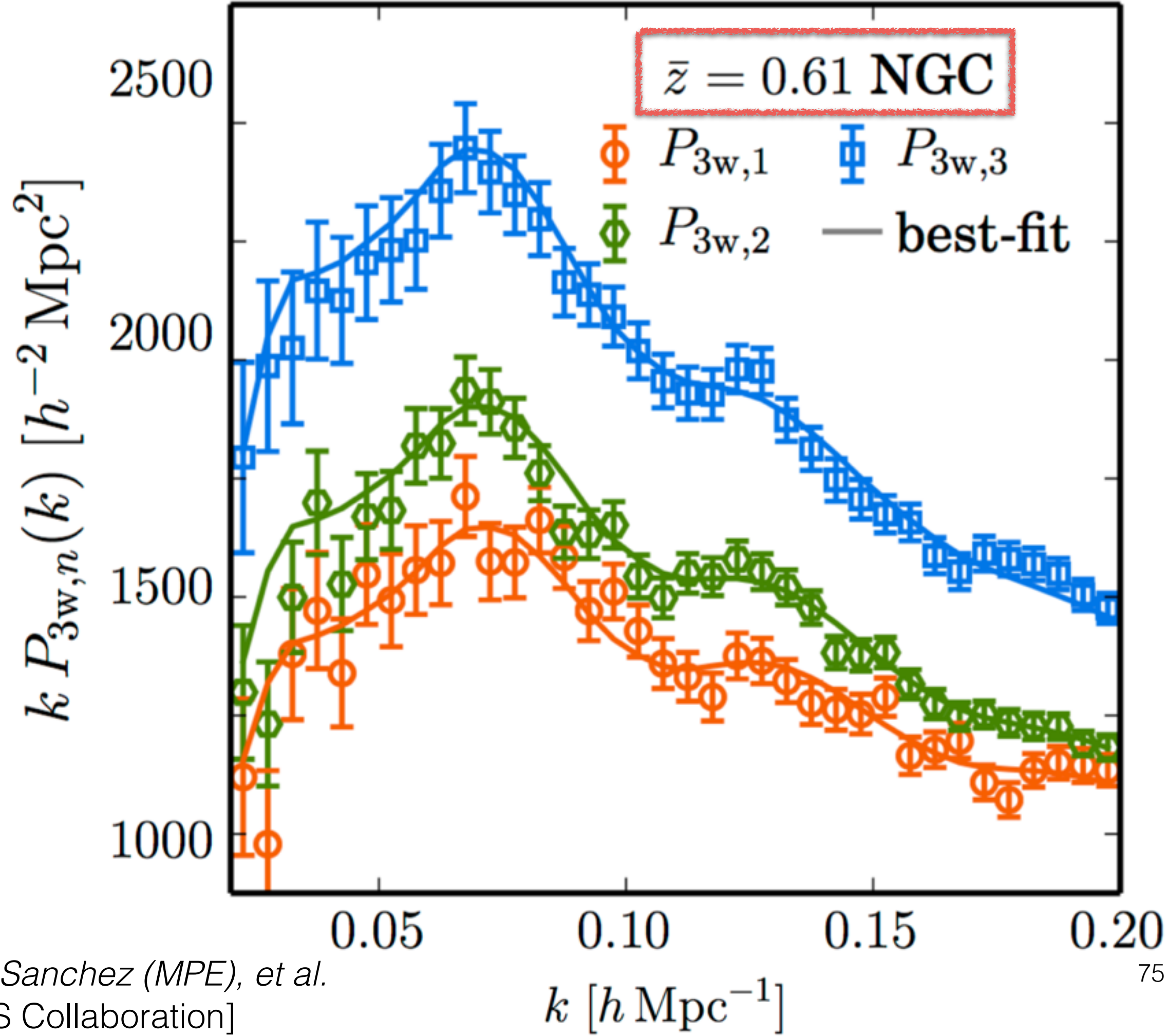
Hot Gas [$z < 1$]



**3-dimensional positions of
a million galaxies have been mapped over
 $\sim 10,000 \text{ deg}^2$ of the sky**







What have we learned?

- **Theory continues to fit!** with the cosmological parameters consistent with CMB, lensing, and SZ effect
- What else can we learn?

A Simple Question

- How do the cosmic structures evolve in an over-dense region?

Simple Statistics

500 Mpc/h



- Divide the survey volume into many sub-volumes V_L , and compare locally-measured power spectra with the corresponding local over-densities

Simple Statistics

500 Mpc/h



V_L

- Divide the survey volume into many sub-volumes V_L , and compare locally-measured power spectra with the corresponding local over-densities

Simple Statistics

$$\bar{\delta}(\mathbf{r}_L) = \frac{1}{V_L} \int_{V_L} d^3r \delta(\mathbf{r})$$

500 Mpc/h



$\bar{\delta}(\mathbf{r}_L)$ V_L

- Divide the survey volume into many sub-volumes V_L , and compare locally-measured power spectra with the corresponding local over-densities

Simple Statistics

$$\bar{\delta}(\mathbf{r}_L) = \frac{1}{V_L} \int_{V_L} d^3r \delta(\mathbf{r})$$

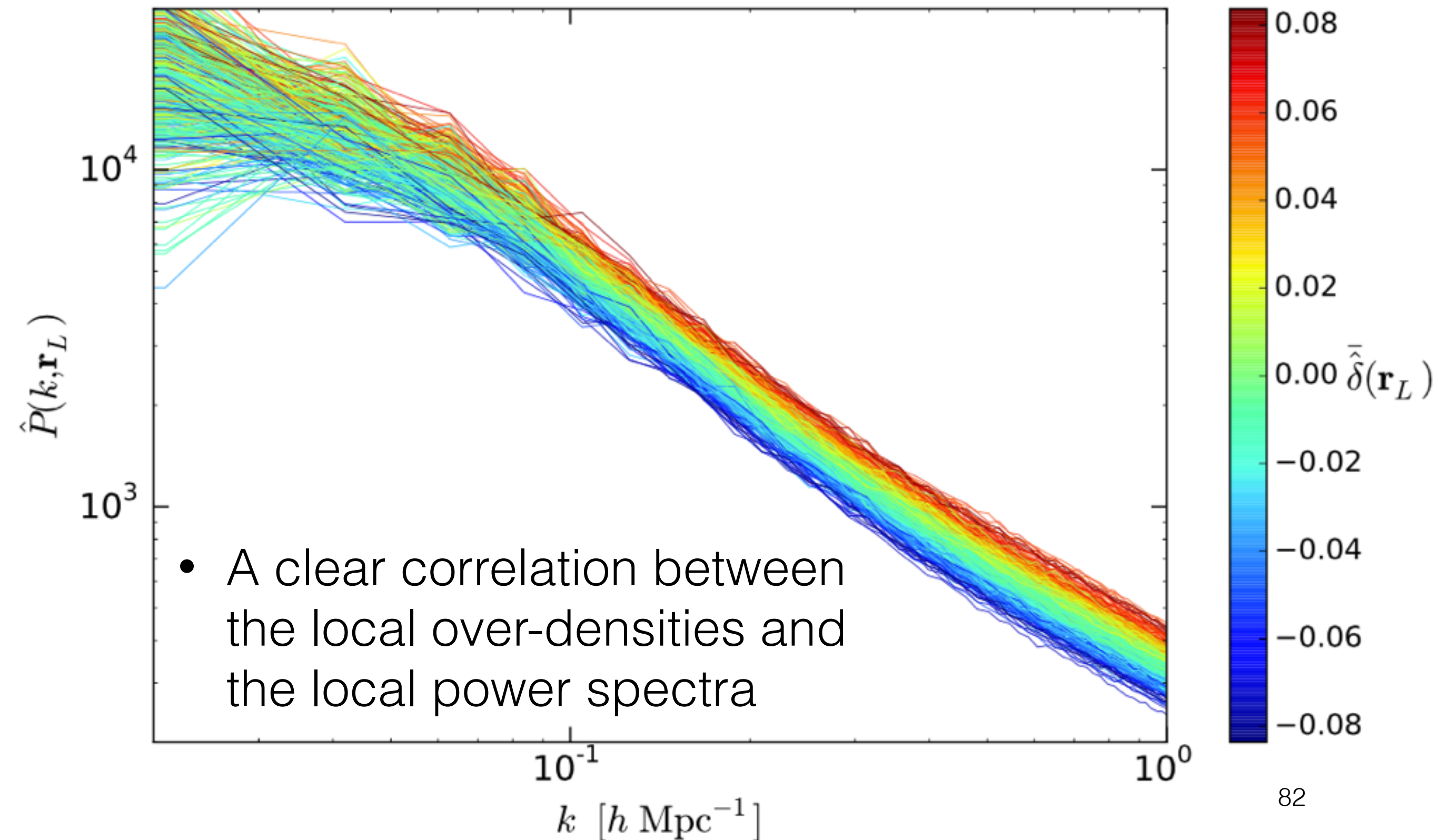
500 Mpc/h

$$\hat{P}(\mathbf{k}, \mathbf{r}_L) \equiv \frac{1}{V_L} |\delta(\mathbf{k}, \mathbf{r}_L)|^2$$

V_L
 $\bar{\delta}(\mathbf{r}_L)$
 $\hat{P}(\mathbf{k}, \mathbf{r}_L)$

- Divide the survey volume into many sub-volumes V_L , and compare locally-measured power spectra with the corresponding local over-densities

Position-dependent $\hat{P}(k)$



Integrated Bispectrum, $iB(k)$

- Correlating the local over-densities and power spectra, we obtain the “integrated bispectrum”:

$$i\hat{B}_L(k) = \frac{1}{N_{\text{cut}}^3} \sum_{i=1}^{N_{\text{cut}}^3} \hat{P}(k, \mathbf{r}_{L,i}) \hat{\delta}(\mathbf{r}_{L,i})$$

- This is a (particular configuration of) **three-point function**! The three-point function in Fourier space is the bispectrum, and is defined as

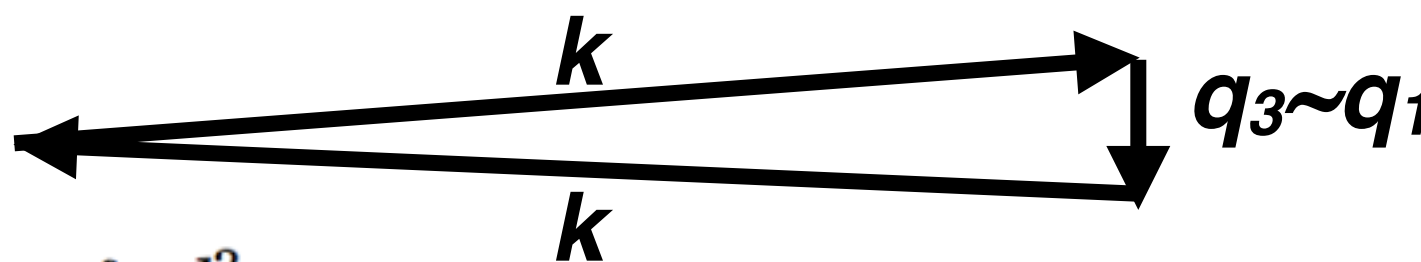
$$\langle \delta(\mathbf{q}_1) \delta(\mathbf{q}_2) \delta(\mathbf{q}_3) \rangle = B(\mathbf{q}_1, \mathbf{q}_2, \mathbf{q}_3) (2\pi)^3 \delta_D(\mathbf{q}_1 + \mathbf{q}_2 + \mathbf{q}_3)$$

Integrated Bispectrum, $iB(k)$

- Correlating the local over-densities and power spectra, we obtain the “integrated bispectrum”:

$$i\hat{B}_L(k) = \frac{1}{N_{\text{cut}}^3} \sum_{i=1}^{N_{\text{cut}}^3} \hat{P}(k, \mathbf{r}_{L,i}) \hat{\delta}(\mathbf{r}_{L,i})$$

- The expectation value of this quantity is an integral of the bispectrum that picks up the contributions **mostly from the squeezed limit**:



$$iB_L(k) = \langle \hat{P}(k, \mathbf{r}_L) \bar{\delta}(\mathbf{r}_L) \rangle$$

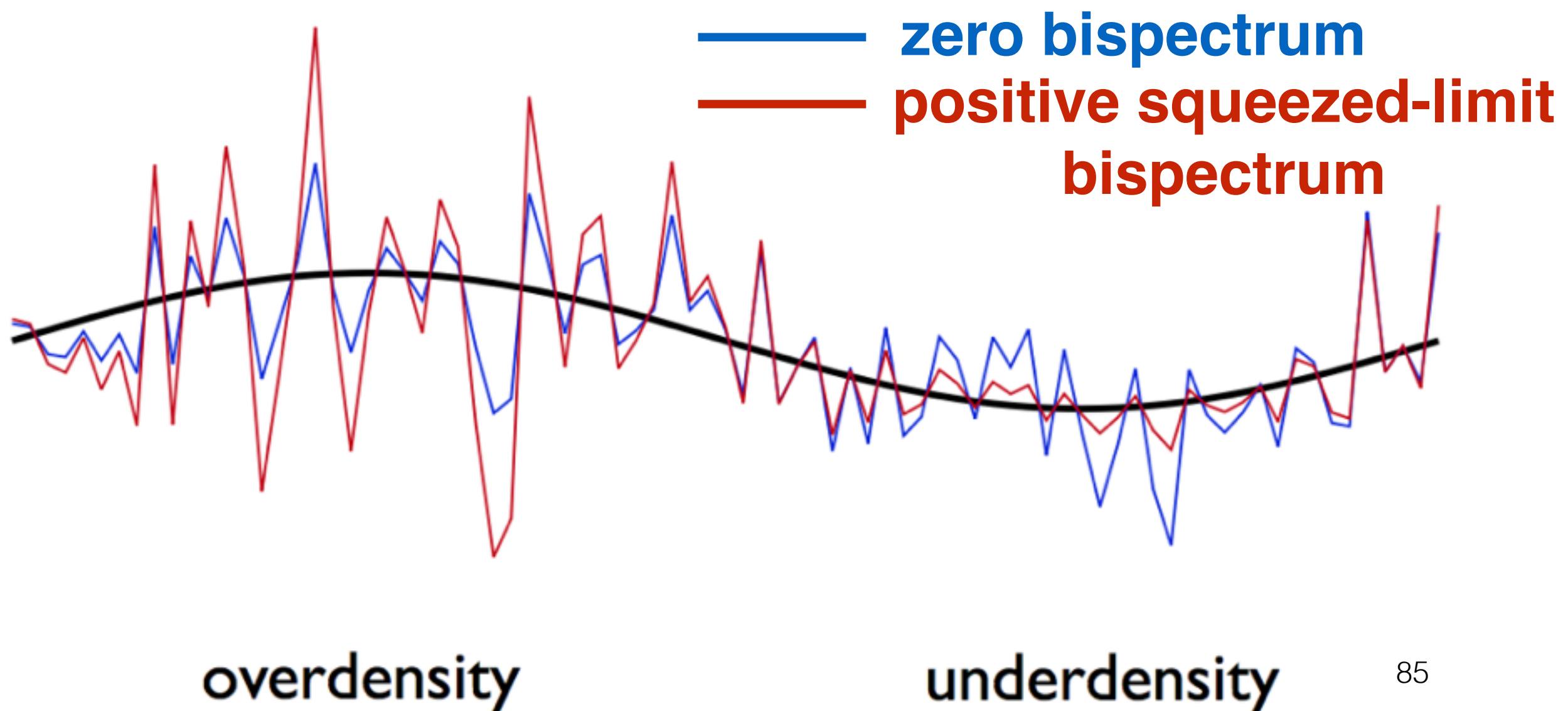
$$= \frac{1}{V_L^2} \int \frac{d^2 \hat{k}}{4\pi} \int \frac{d^3 q_1}{(2\pi)^3} \int \frac{d^3 q_3}{(2\pi)^3} B(\mathbf{k} - \mathbf{q}_1, -\mathbf{k} + \mathbf{q}_1 + \mathbf{q}_3, -\mathbf{q}_3)$$

“taking the squeezed limit and then angular averaging”

$$\times W_L(\mathbf{q}_1) W_L(-\mathbf{q}_1 - \mathbf{q}_3) W_L(\mathbf{q}_3)$$

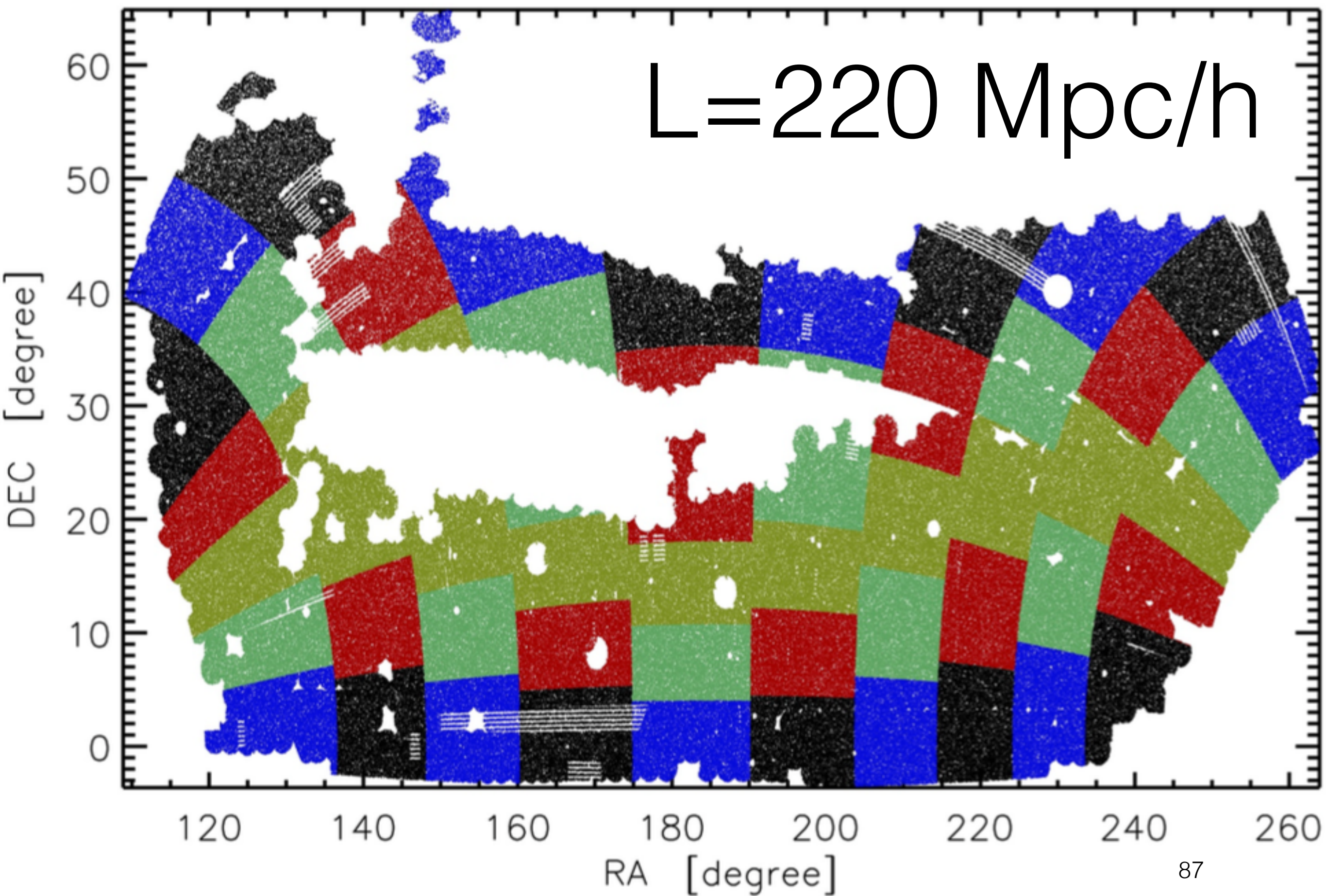
Power Spectrum Response

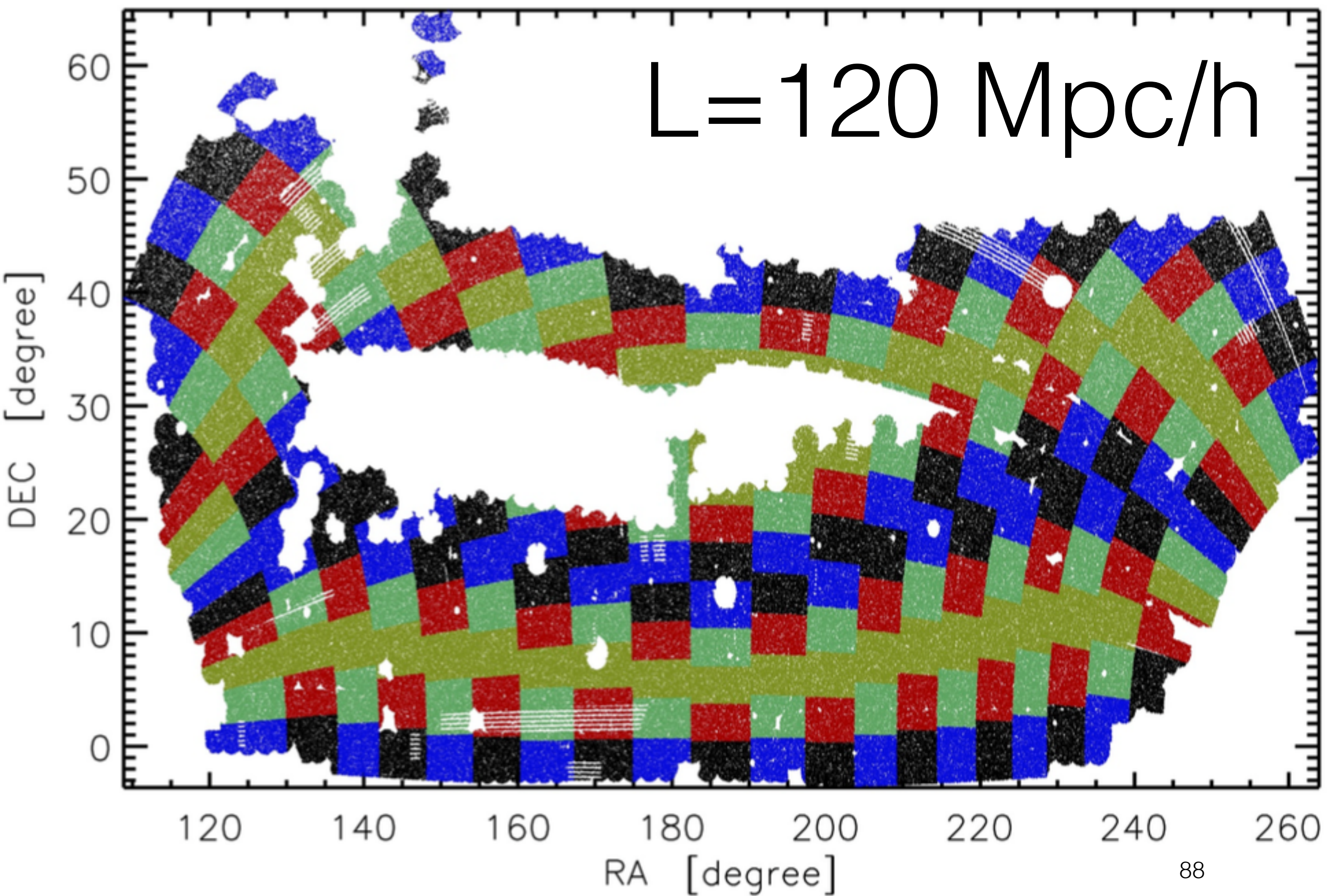
- The integrated bispectrum measures how the local power spectrum responds to its environment, i.e., a long-wavelength density fluctuation



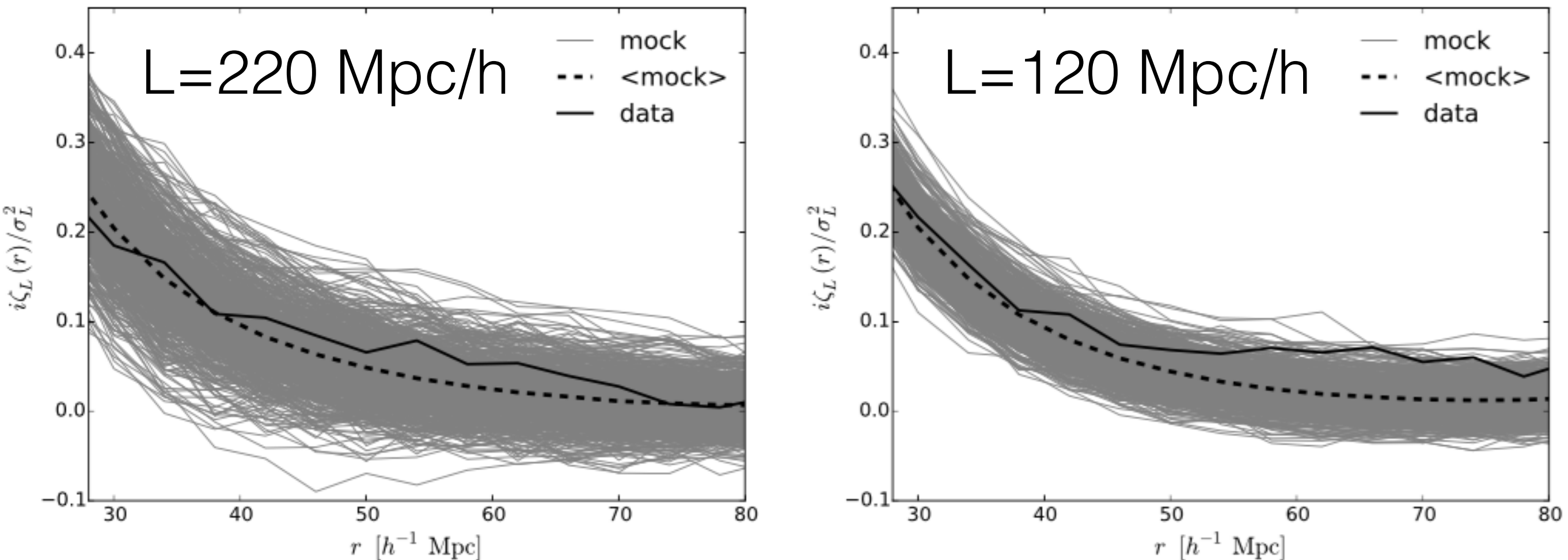
SDSS-III/BOSS DR10

- OK, now, let's look at the real data (BOSS DR10) to see if we can detect the expected influence of environments on the small-scale structure growth





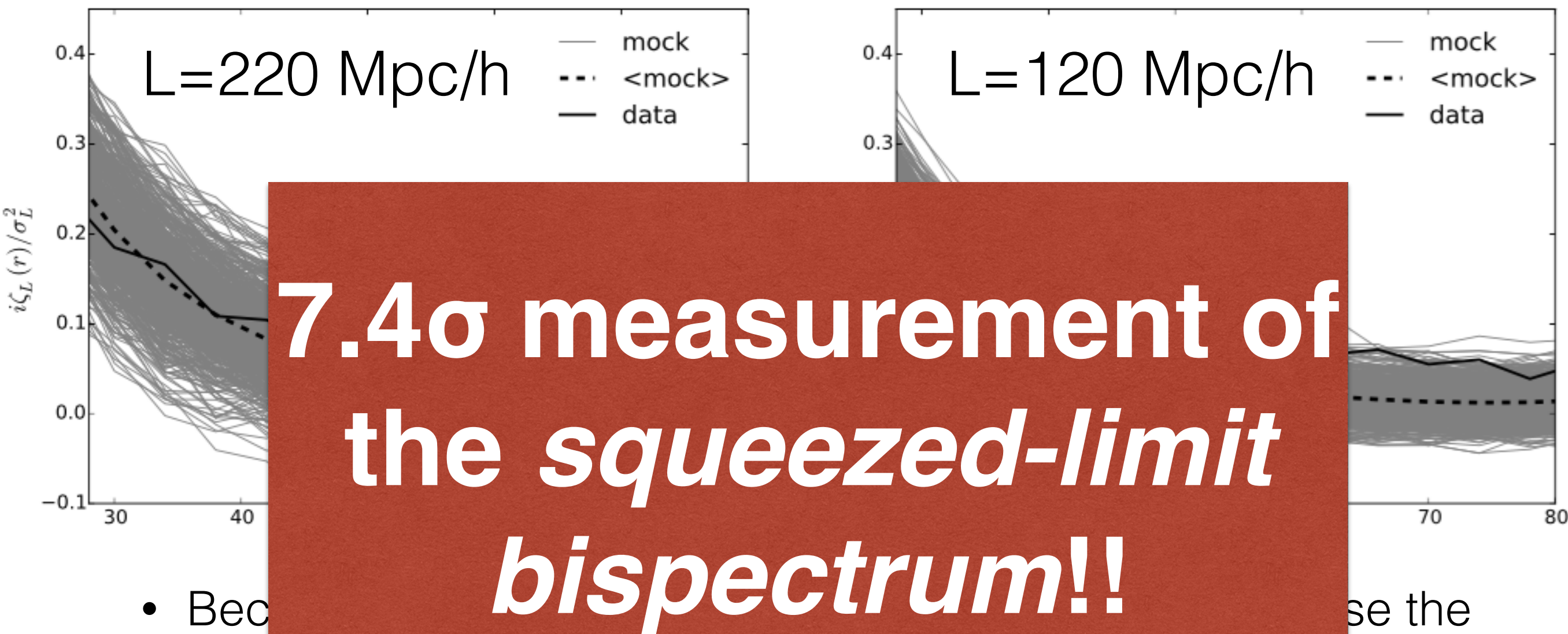
Results: $\chi^2/\text{DOF} = 46.4/38$



- Because of complex geometry of DR10 footprint, we use the local correlation function, instead of the power spectrum. Power spectrum will be presented using DR12 in the future

- Integrated three-point function, $i\zeta(r)$, is just Fourier transform of $iB(k)$:
$$i\zeta_L(\mathbf{r}) = \int \frac{d^3k}{(2\pi)^3} iB_L(\mathbf{k}) e^{i\mathbf{r}\cdot\mathbf{k}}$$

Results: $\chi^2/\text{DOF} = 46.4/38$



- Because of the local nature of the measurement, the power spectrum will be presented using DR12 in the future

- Integrated three-point function, $i\zeta(\mathbf{r})$, is just Fourier transform of $iB(\mathbf{k})$:

$$i\zeta_L(\mathbf{r}) = \int \frac{d^3k}{(2\pi)^3} iB_L(\mathbf{k}) e^{i\mathbf{r}\cdot\mathbf{k}}$$

Nice, but what is this good for?

- Primordial non-Gaussianity (“local-type f_{NL} ”)
- The constraint from BOSS is work in progress, but the Fisher matrix analysis suggests that the integrated bispectrum is a **nearly optimal estimator for the local-type f_{NL}**
- We no longer need to measure the full bispectrum, if we are just interested in $f_{\text{NL}}^{\text{local}}$!

Nice, but what is this good for?

- Primordial non-Gaussianity (“local-type f_{NL} ”)

- The bispectrum is insensitive to the initial condition!

- We no longer need to measure the full bispectrum, if we are just interested in $f_{\text{NL}}^{\text{local}}$!

Nice, but what is this good for?

- We can also learn about galaxy bias
 - Local bias model:
 - $\delta_g(x) = b_1 \delta_m(x) + (\mathbf{b}_2/2) [\delta_m(x)]^2 + \dots$
- **The bispectrum can give us b_2 at the leading (tree-level) order**, unlike for the power spectrum that has b_2 at the next-to-leading order

Result on b_2

- We use the simplest, tree-level SPT bispectrum in redshift space with the local bias model to interpret our measurements
- [We also use information from BOSS's 2-point correlation function on σ_8 and BOSS's weak lensing data on σ_8]
- We find: **$b_2 = 0.41 \pm 0.41$**

CMB Research: Next Frontier

Primordial Gravitational Waves

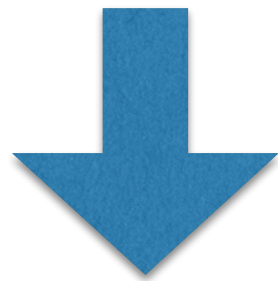
Extraordinary claims require extraordinary evidence.

The same quantum fluctuations could also generate gravitational waves, and we wish to find them

Measuring GW

- GW changes the distances between two points

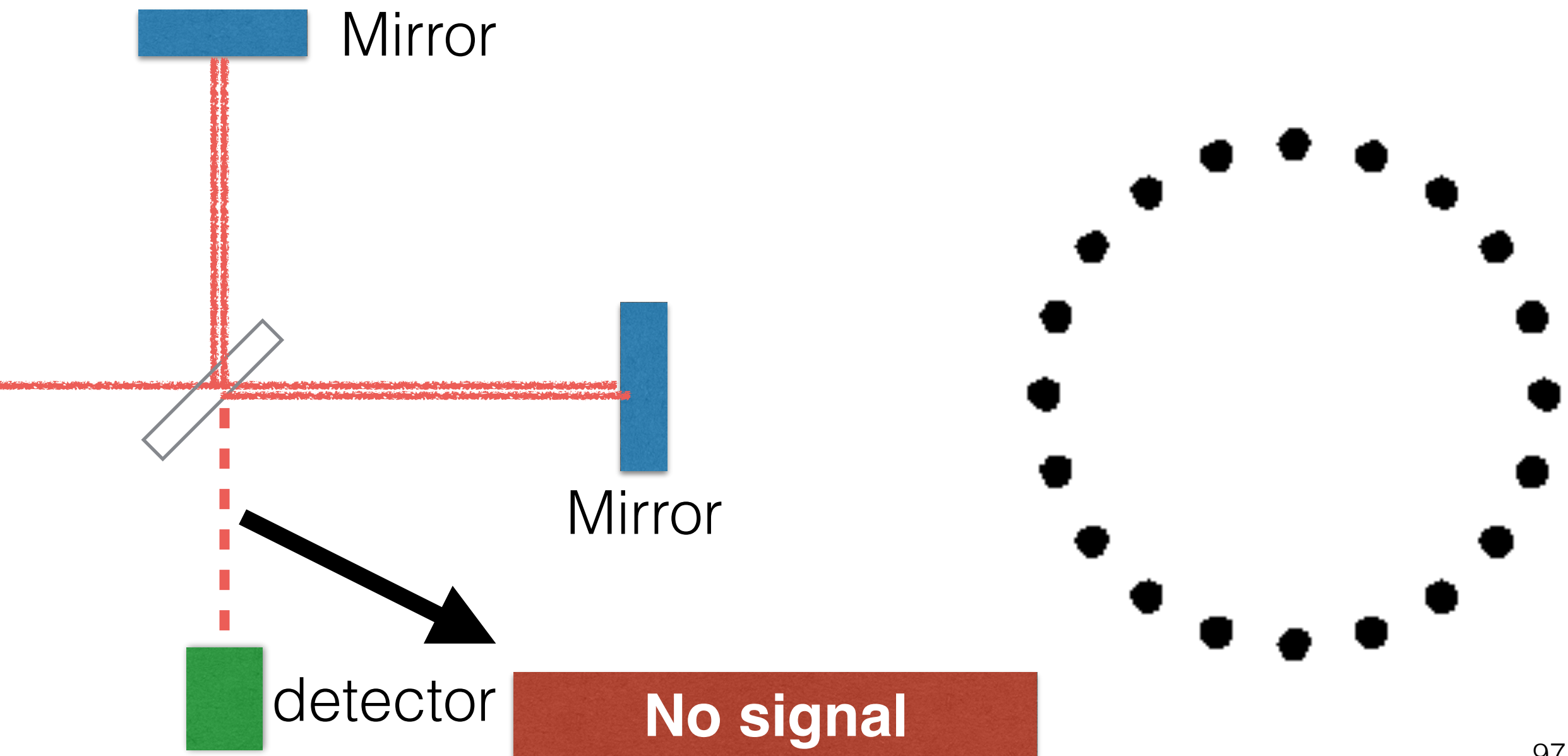
$$d\ell^2 = d\mathbf{x}^2 = \sum_{ij} \delta_{ij} dx^i dx^j$$



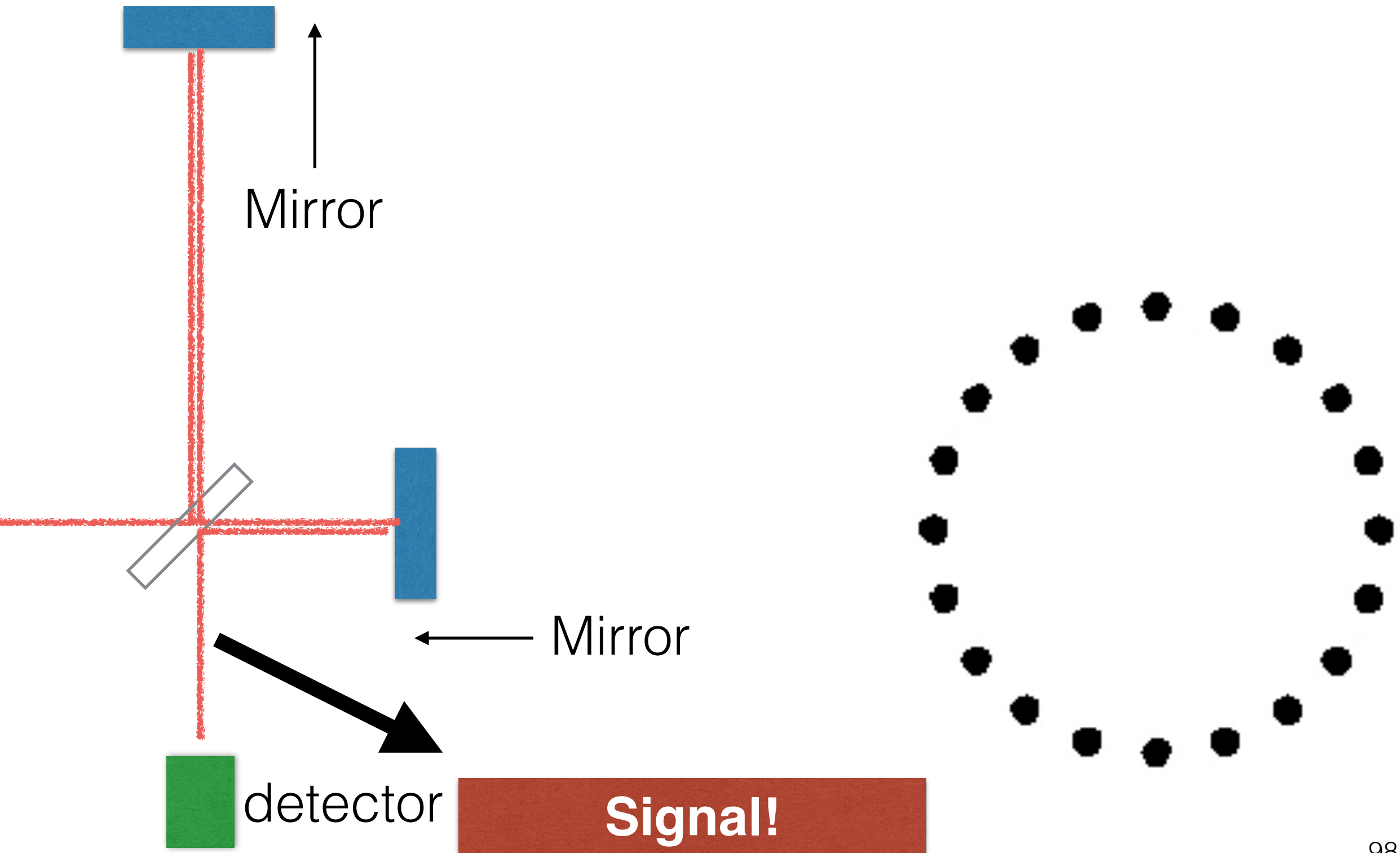
$$d\ell^2 = \sum_{ij} (\delta_{ij} + \textcolor{red}{h_{ij}}) dx^i dx^j$$



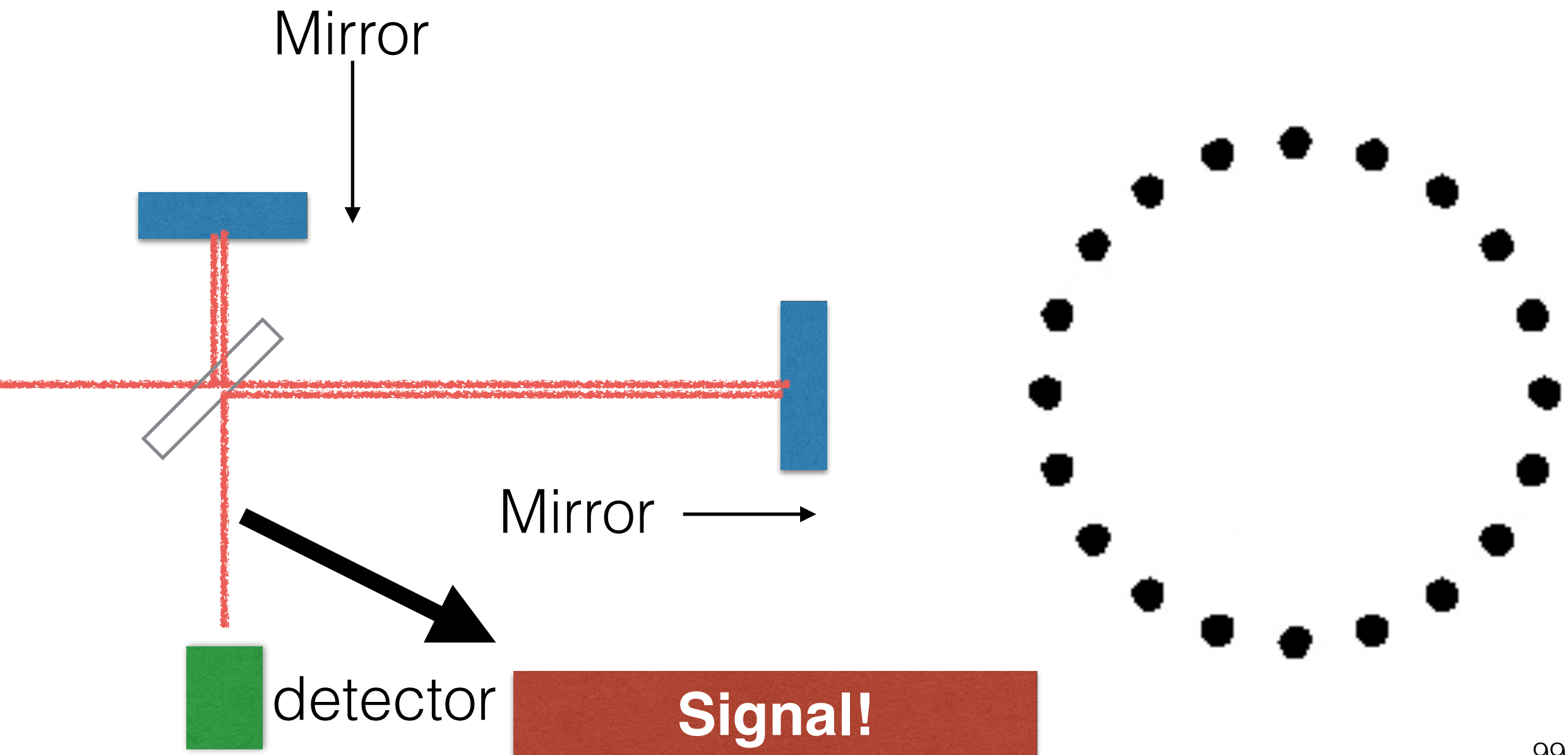
Laser Interferometer



Laser Interferometer



Laser Interferometer



LIGO detected GW from a binary blackholes, with the wavelength of thousands of kilometres

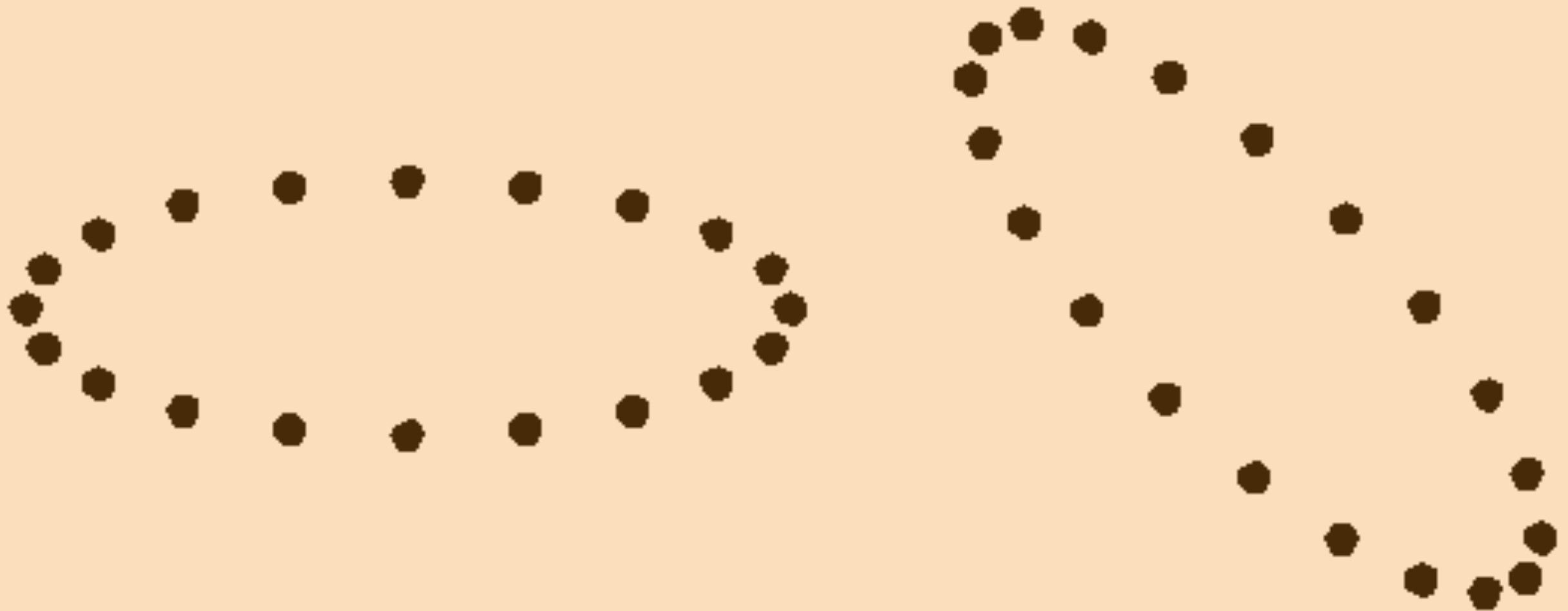
But, the primordial GW affecting the CMB has a wavelength of **billions of light-years!!** How do we find it?

Detecting GW by CMB

Isotropic electro-magnetic fields

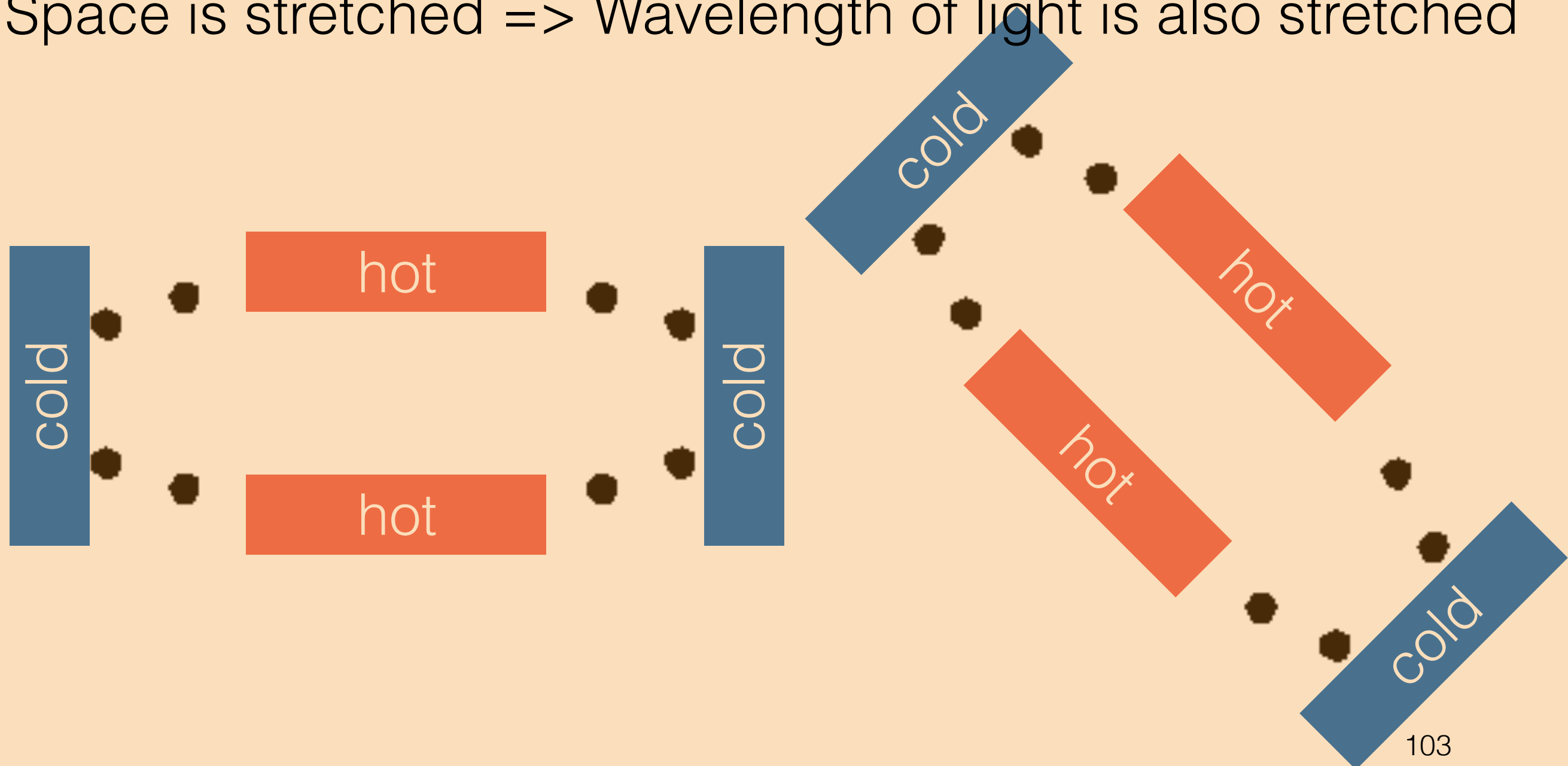
Detecting GW by CMB

GW propagating in isotropic electro-magnetic fields



Detecting GW by CMB

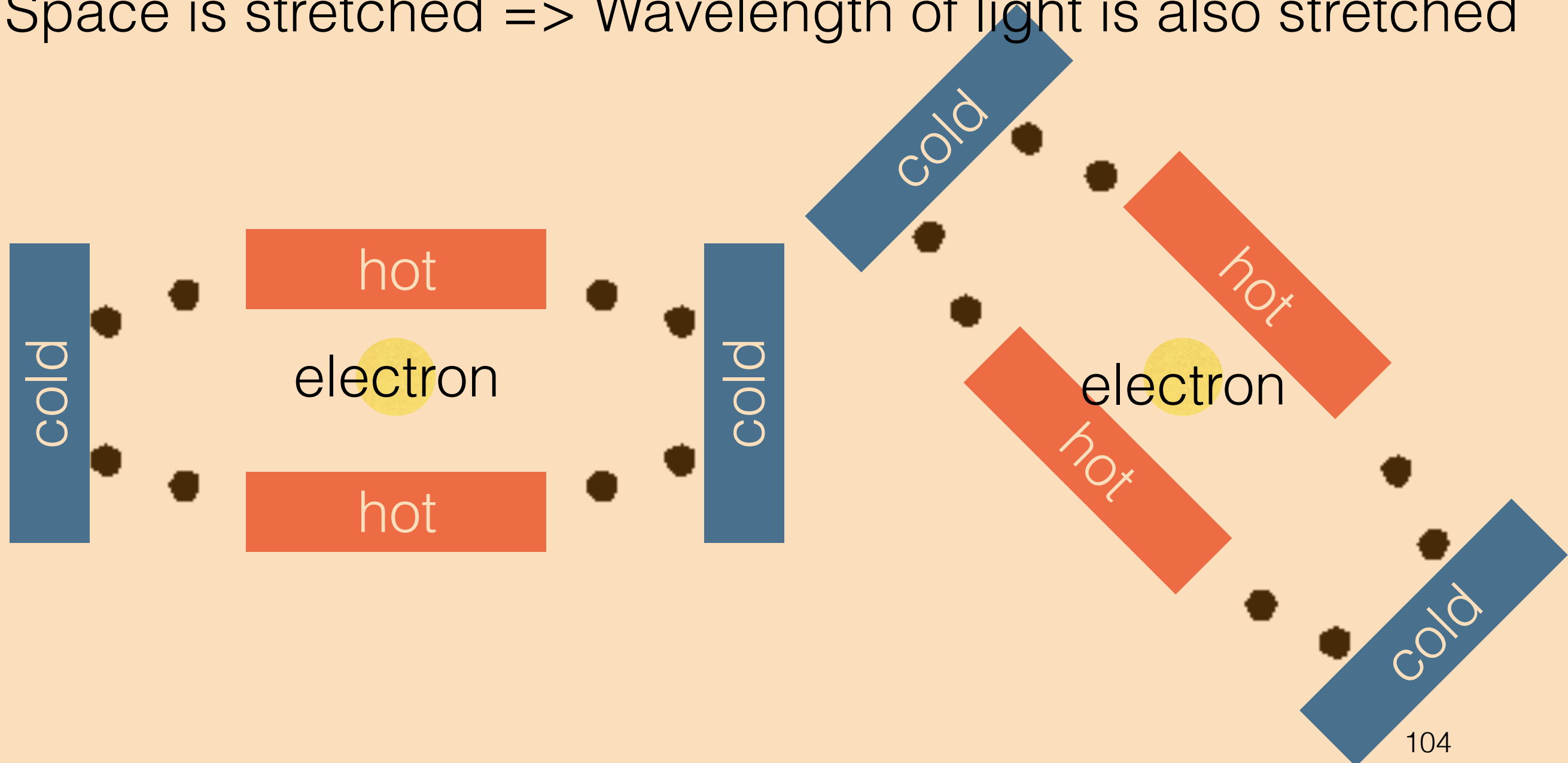
Space is stretched $\Rightarrow$ Wavelength of light is also stretched



Detecting GW by CMB

Polarisation

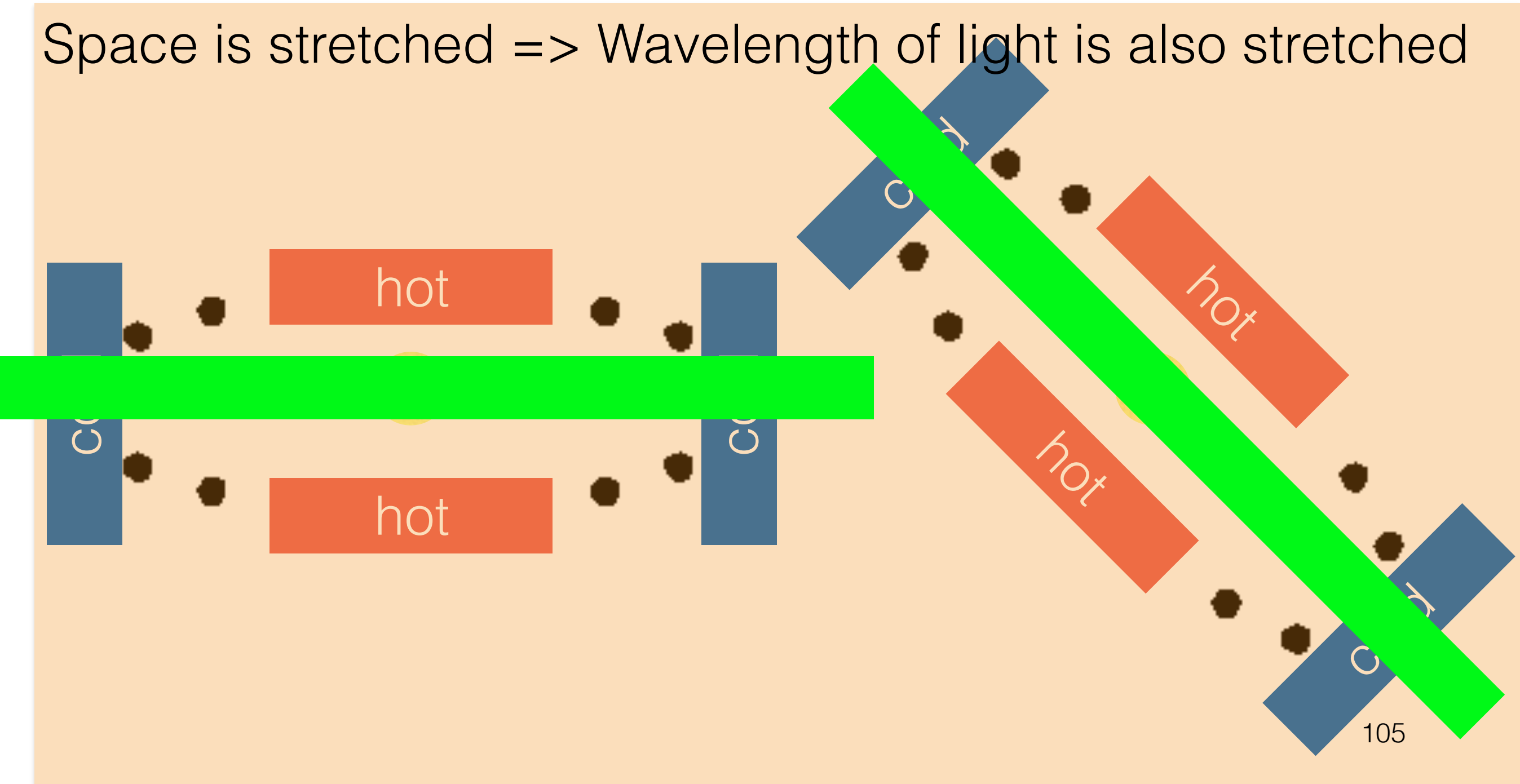
Space is stretched => Wavelength of light is also stretched



Detecting GW by CMB

Polarisation

Space is stretched => Wavelength of light is also stretched



Current Situation

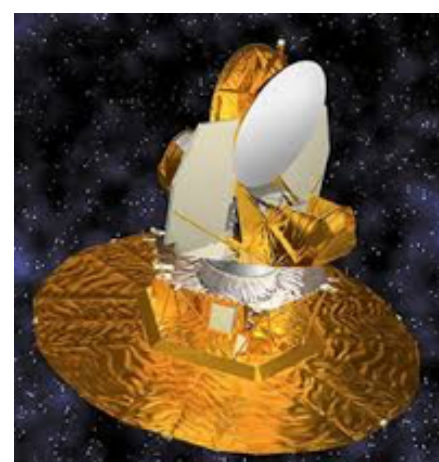
- No detection of polarisation from primordial GW yet
- Many ground-based and balloon-borne experiments are taking data now



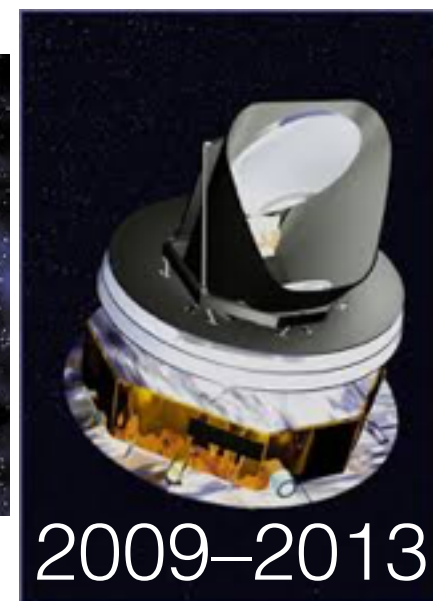
The search continues!!



1989–1993



2001–2010



2009–2013



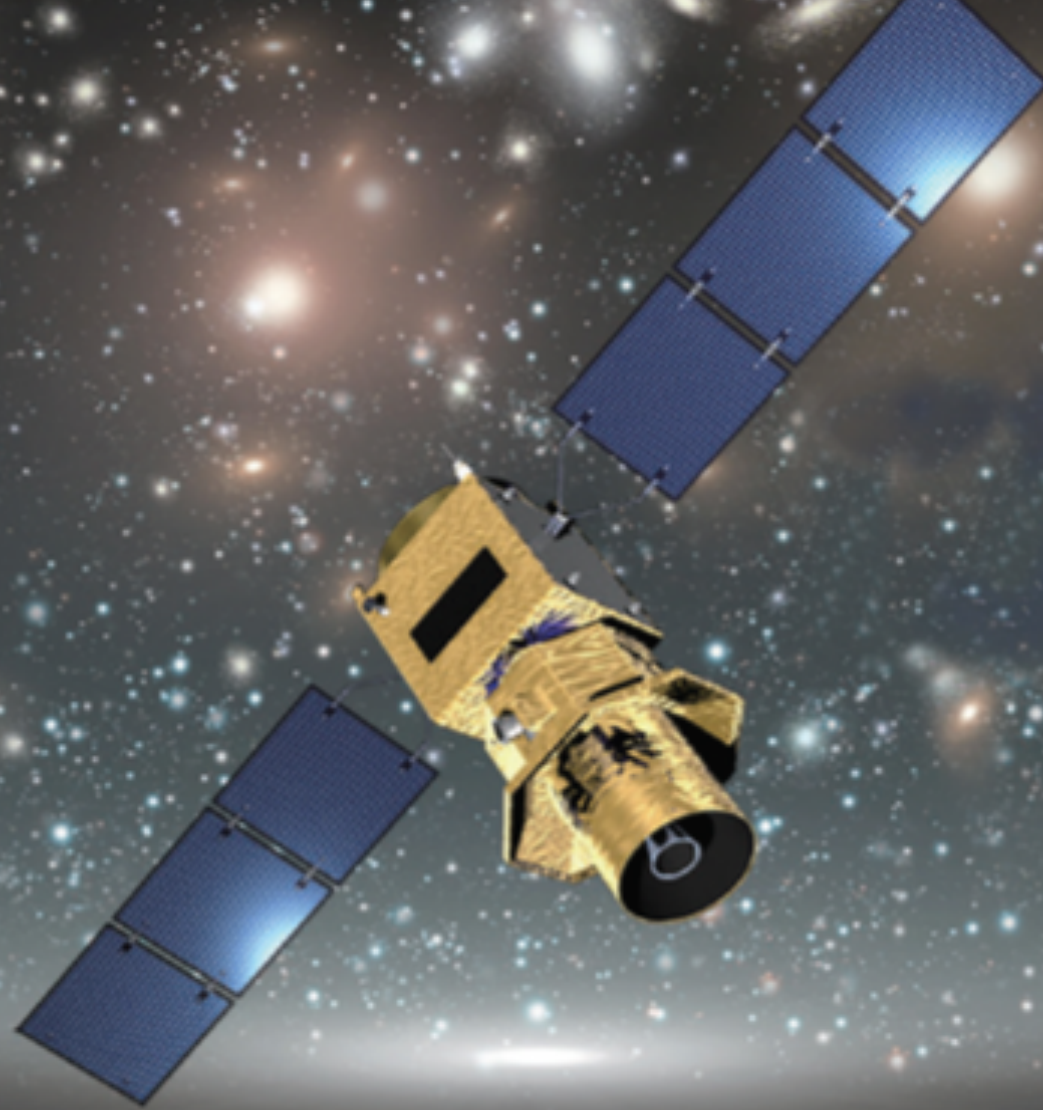
202X–

JAXA

+ possibly NASA

LiteBIRD

2025– [proposed]

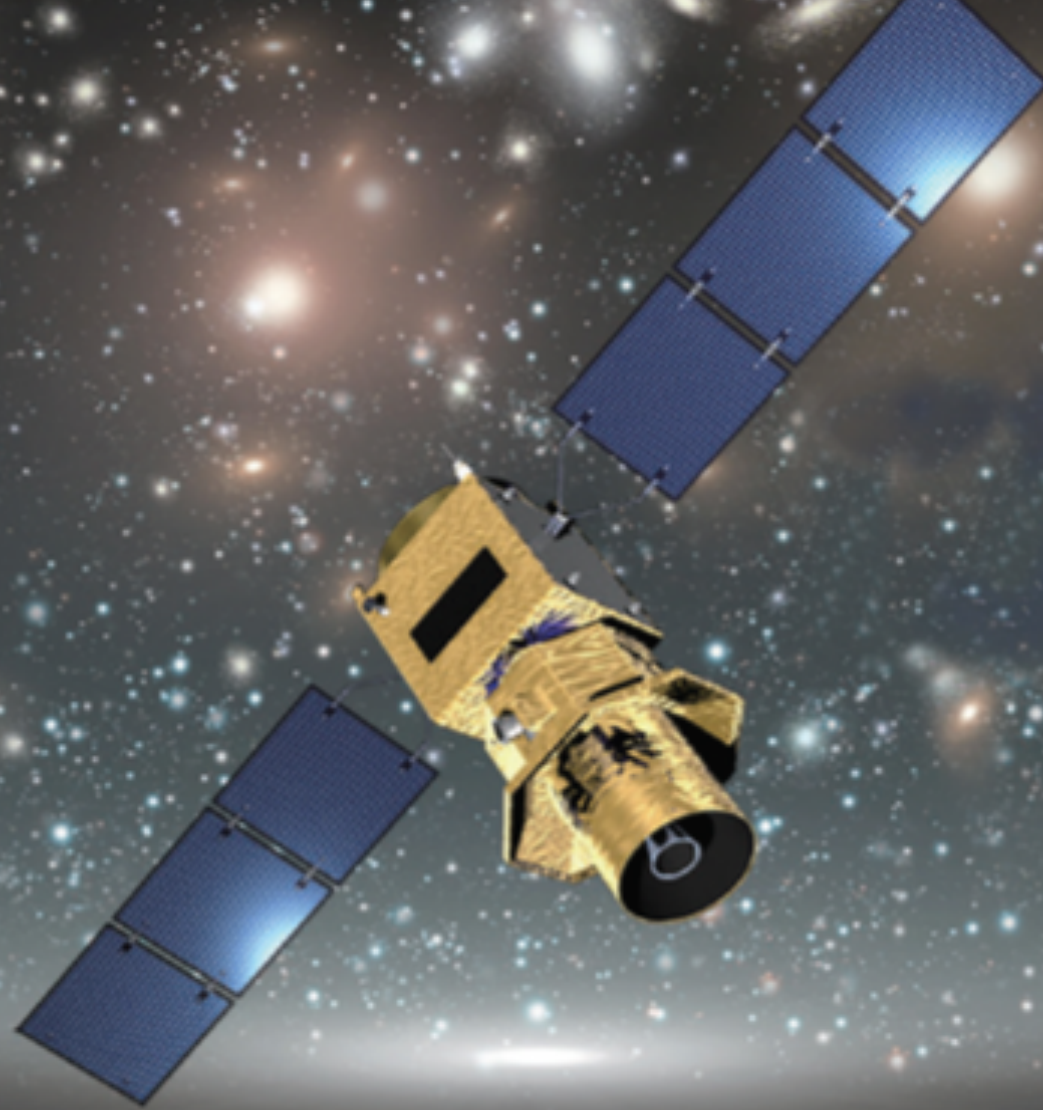


JAXA

+ possibly NASA

LiteBIRD

2025– [proposed]



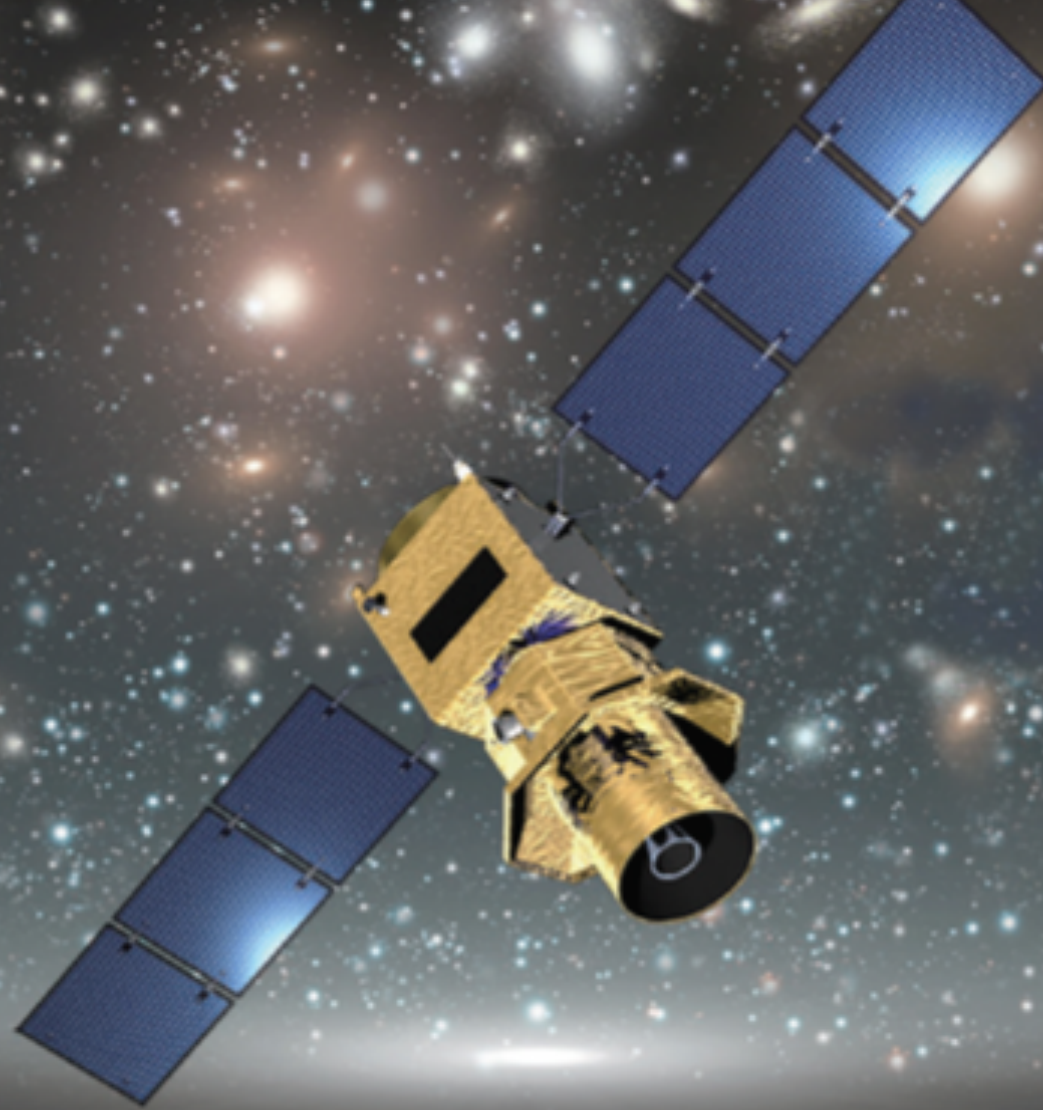
**Polarisation satellite dedicated to
measure CMB polarisation from
primordial GW, with a couple thousand
super-conducting detectors in space**

JAXA

+ possibly NASA

LiteBIRD

2025– [proposed]



**Down-selected by JAXA as
one of the two missions
competing for a launch in 2025**

Conclusions:

From Initial Conditions...

- **$n_s < 1$ finally discovered**. Time-translation symmetry of space-time is broken during inflation, as predicted
- **No evidence for breaking of rotational symmetry**, consistent with the simplest model
- Strong evidence for inflation, but not a final word
- Vigorous searches for the **primordial GW from inflation** are underway. Fingers crossed for the final selection of LiteBIRD by JAXA

Conclusions: to Structure Formation...

- **Secondary anisotropy of the CMB** tells us about the subsequent gravitational and baryonic evolution of initial fluctuations
 - Theory fits!
- The large-scale structure of the universe traced by the distribution of galaxies also tests our theory for the subsequent evolution
 - Theory fits!

Conclusions: ...and Back!!

- Measurements of the three-point function (**bispectrum**) of the **large-scale structure** are not as scary as before
- They can finally be used to test properties of the primordial fluctuations, potentially better than CMB studies could
- **The position-dependent power spectrum is a powerful, easy, and intuitive way to get this job done**

[Review: Chi-Ting Chiang, arXiv:1508.03256]

