

# What does cosmology tell us about physics beyond SM?

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# Do we even need physics beyond SM?

- Let us remind ourselves that the answer to this question is a definite **yes**, despite null results from LHC (Jinnouchi's talk) because:
  - We have dark matter,
  - We have dark energy, and
  - We (probably) have inflation,
- all of which require new physics.

# “Standard Model” of Our Universe

## ● Standard Model

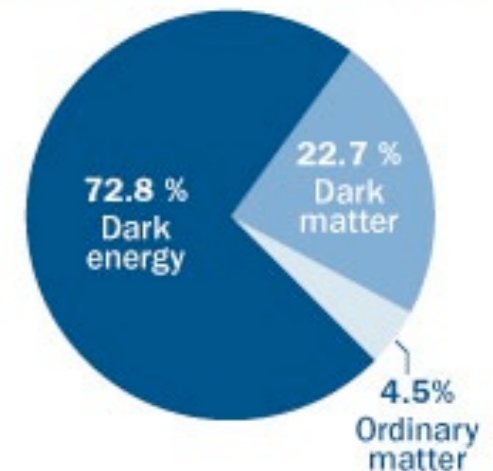
- H&He = 4.56% ( $\pm 0.16\%$ )
- Dark Matter = 27.2% ( $\pm 1.6\%$ )
- Dark Energy = 72.8% ( $\pm 1.6\%$ )
- $H_0 = 70.4 \pm 1.4$  km/s/Mpc
- Age of the Universe = 13.75 billion years ( $\pm 0.11$  billion years)

## Universal Stats

Age of the universe today  
**13.75 billion years**

Age of the cosmos at  
time of reionization  
**457 million years**

Universe composition



*“ScienceNews” article on the WMAP 7-year results*

# *How does cosmology tell us about physics beyond SM?*

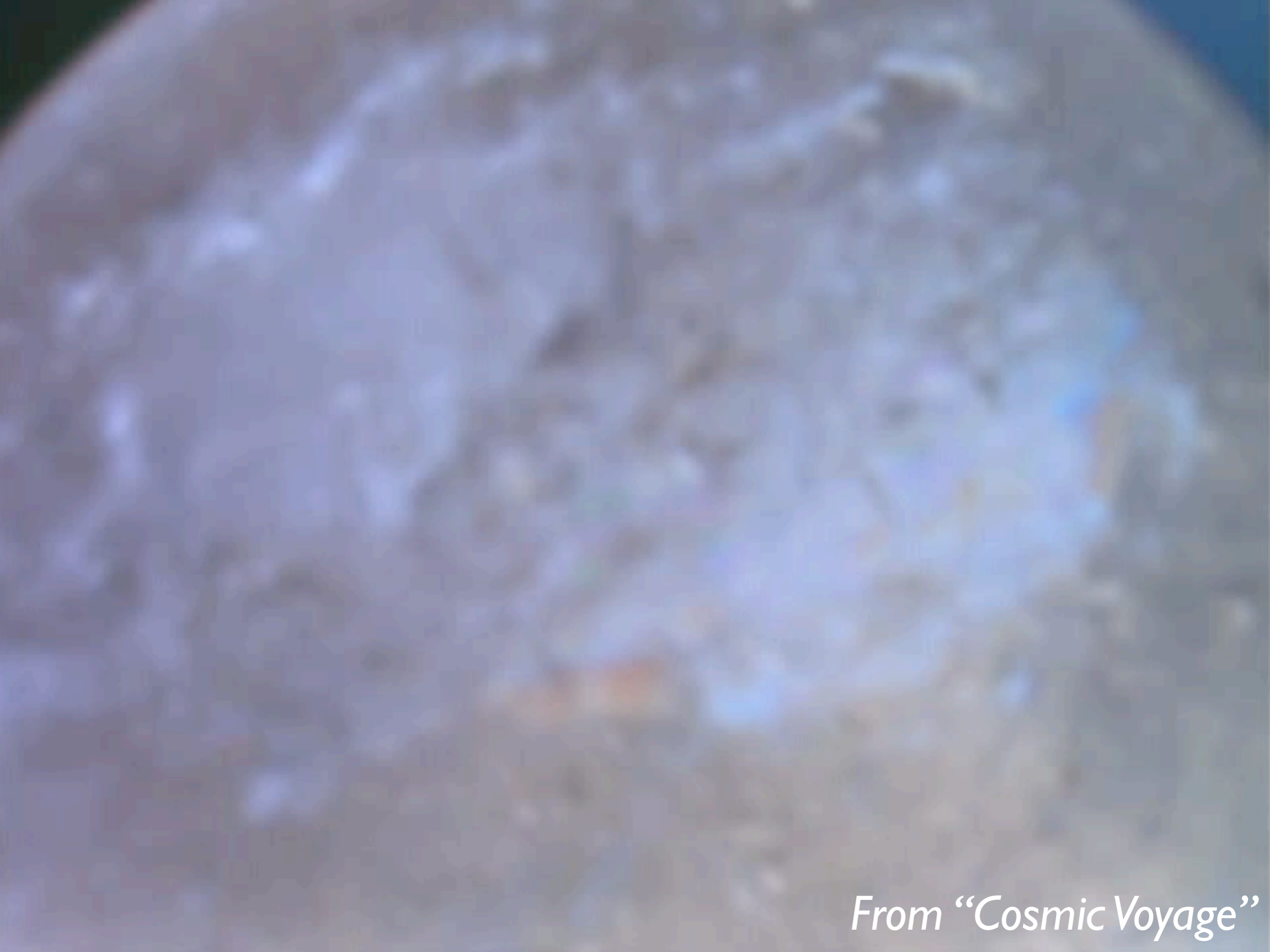
Let me focus on the  
cosmic microwave background (CMB)

# The Breakthrough

- Now we can **observe** the physical condition of the Universe when it was very young.

# CMB

- Fossil light of the Big Bang!



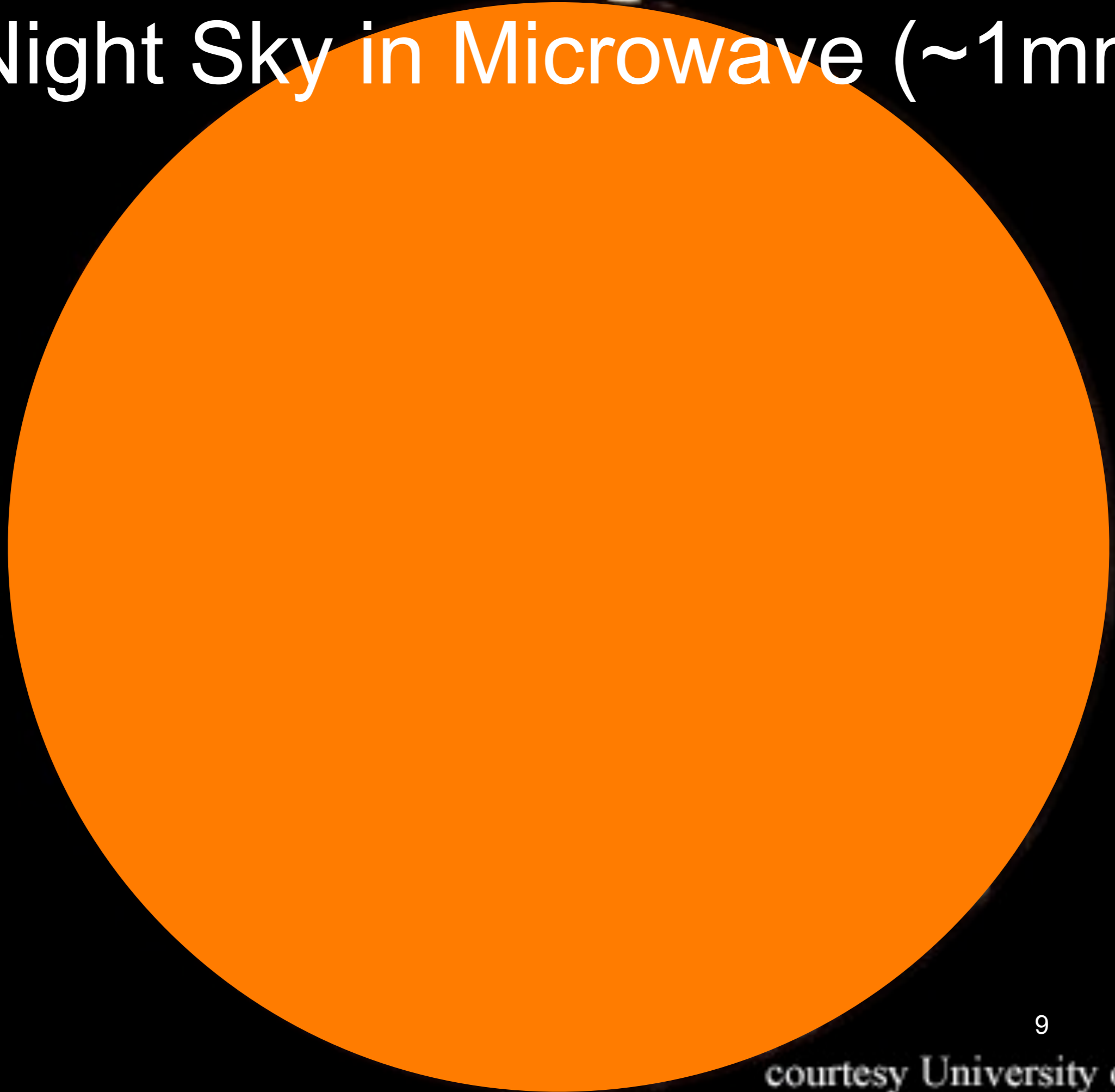
*From "Cosmic Voyage"*

# Night Sky in Optical ( $\sim 0.5\mu\text{m}$ )



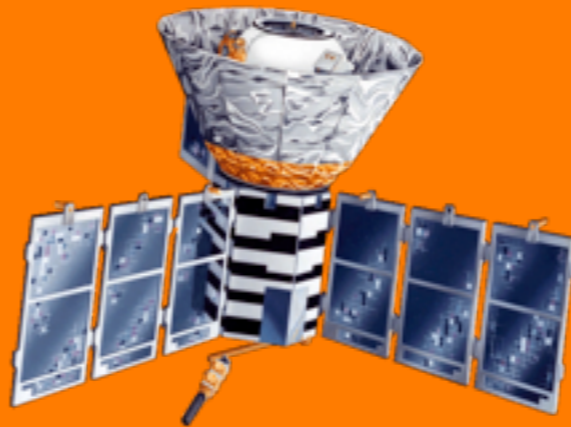


# Night Sky in Microwave ( $\sim 1\text{mm}$ )



# Night Sky in Microwave ( $\sim 1\text{mm}$ )

$T_{\text{today}} = 2.725\text{K}$

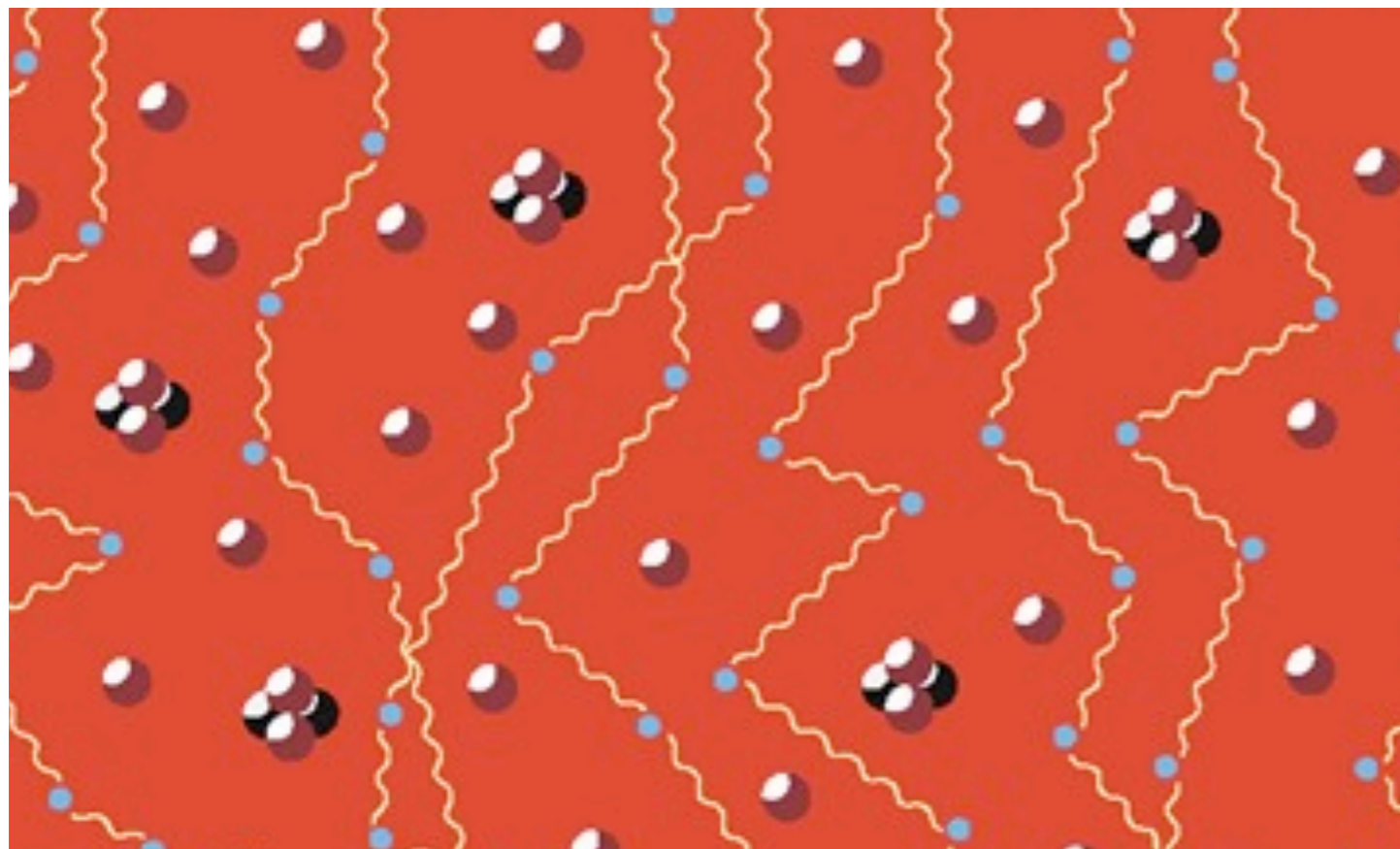


***COBE Satellite, 1989-1993***

# How was CMB created?

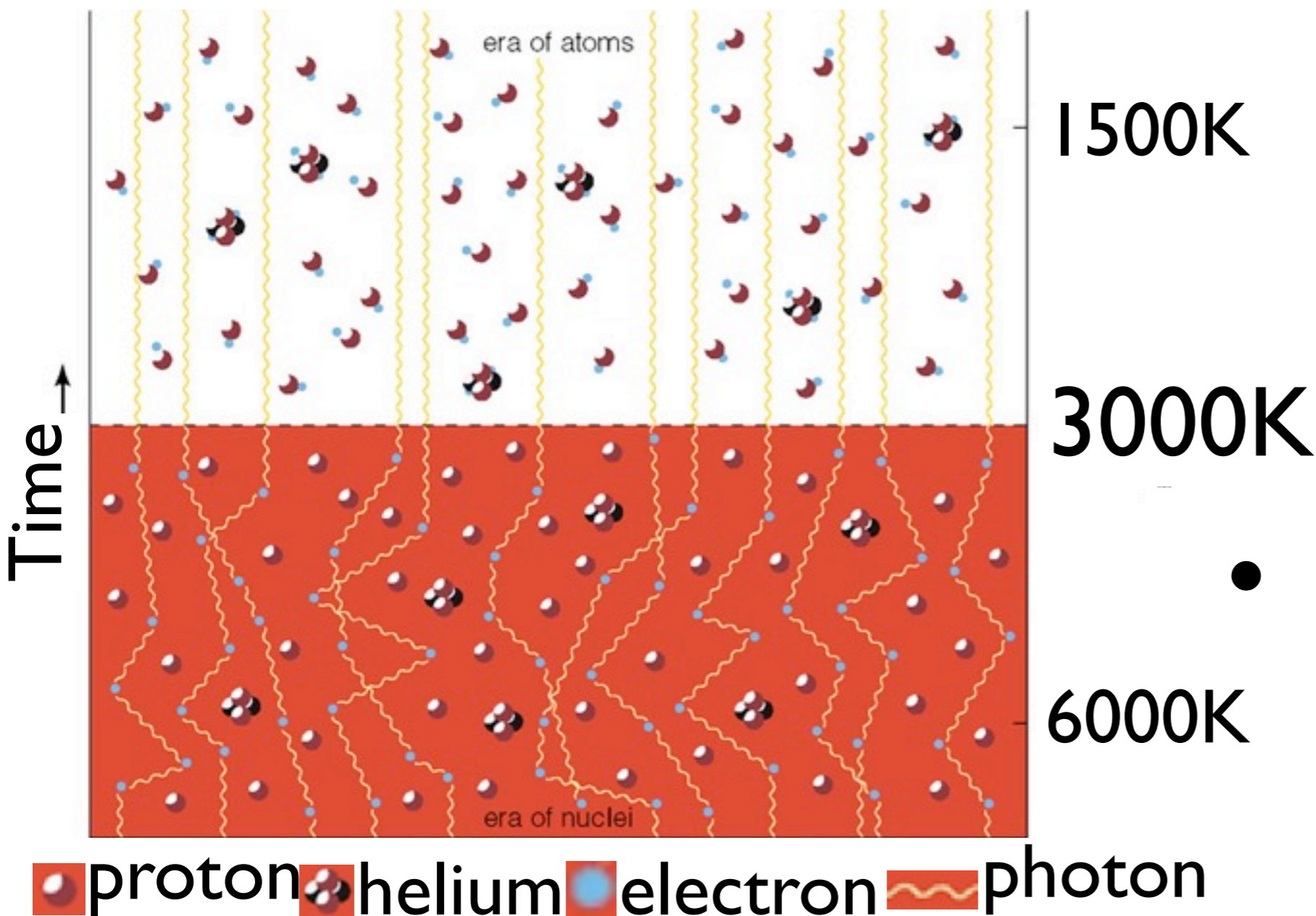
- When the Universe was hot, the Universe was a hot soup made of:
  - Protons, electrons, and helium nuclei
  - Photons and neutrinos
  - Dark matter

# Universe as a hot soup



- Free electrons Thomson-scatter photons efficiently.
- Photons cannot go very far.

# Recombination and Decoupling



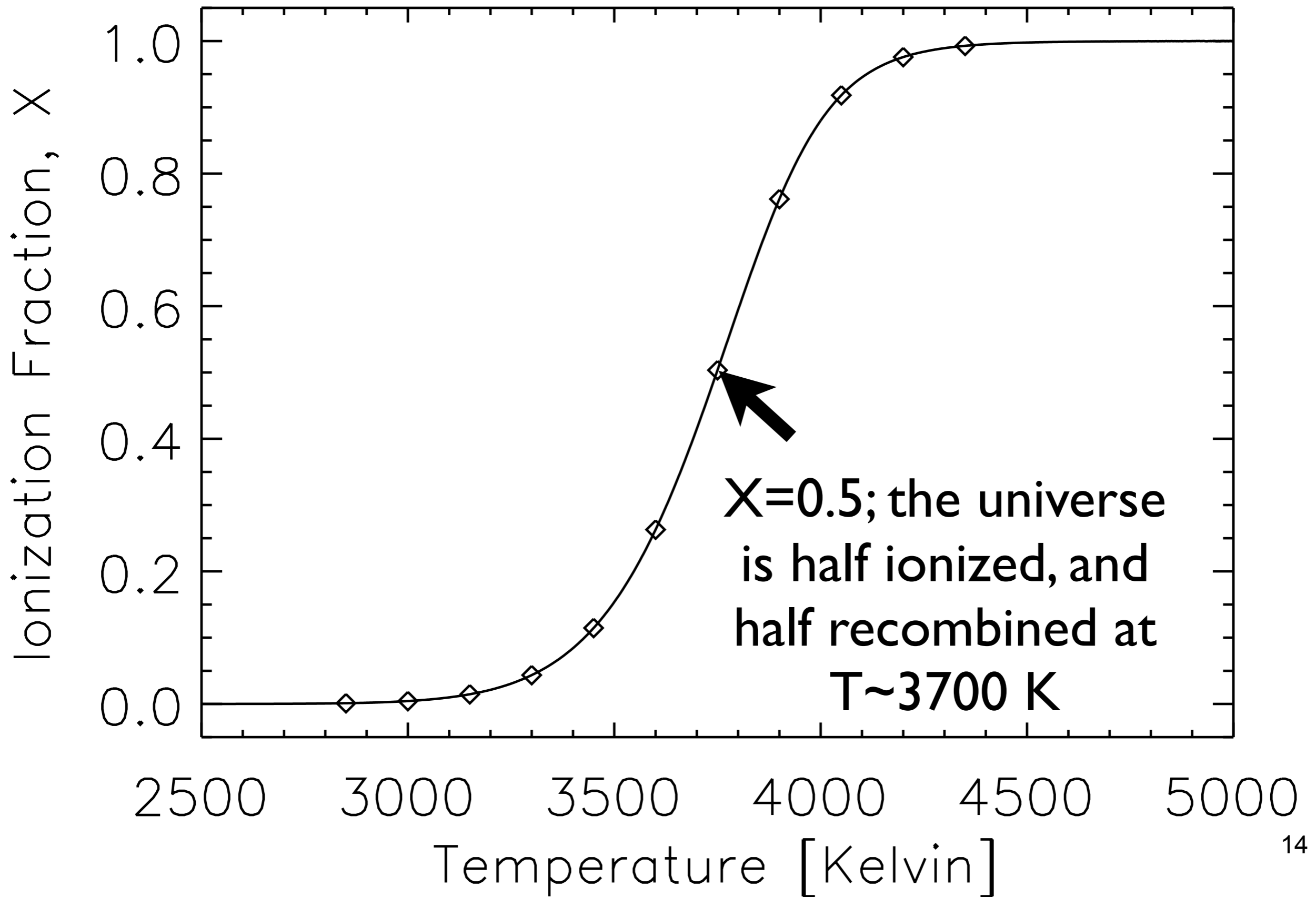
- [**recombination**] When the temperature falls below 3000 K, almost all electrons are captured by protons and helium nuclei.

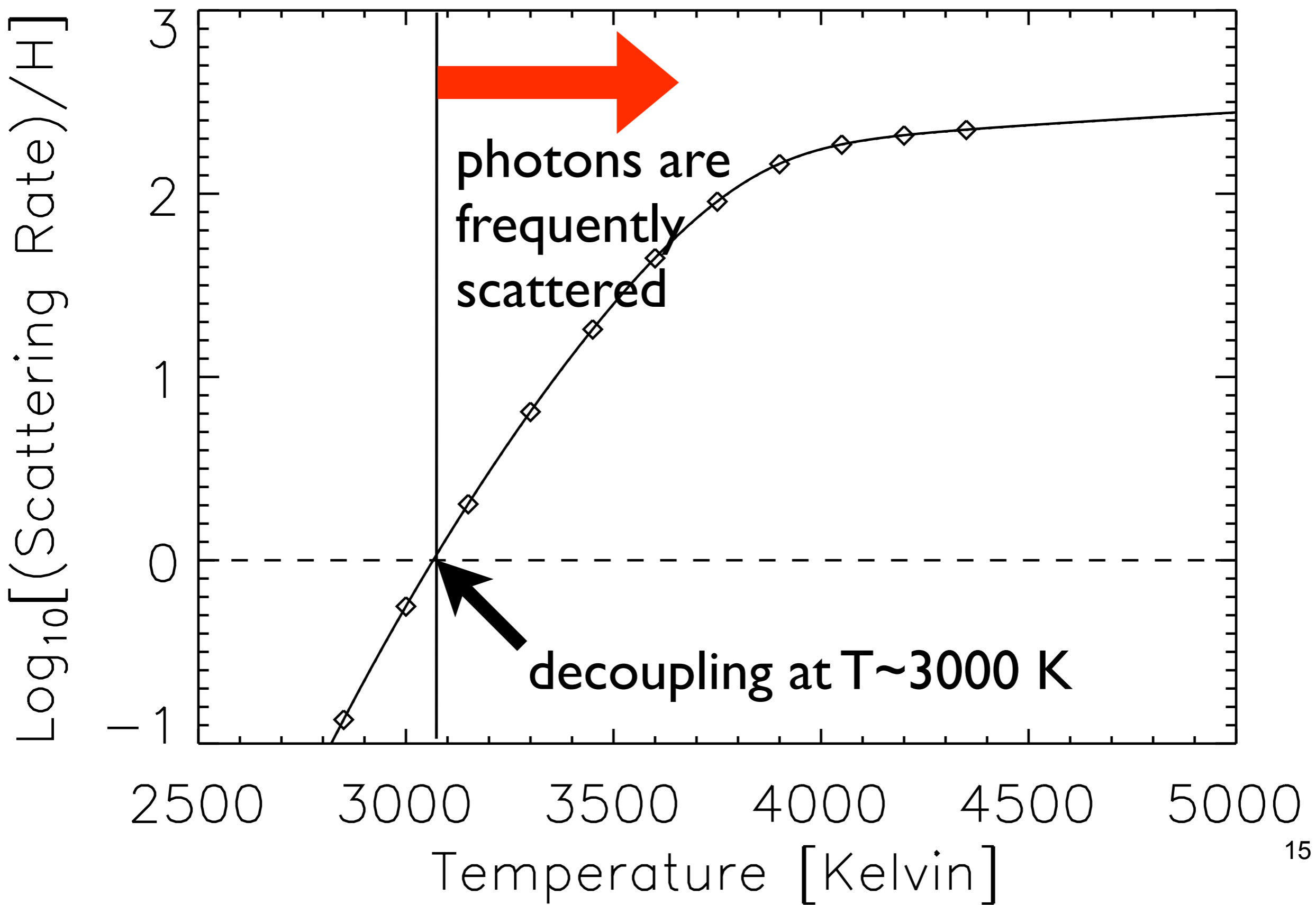
- [**decoupling**] Photons are no longer scattered. I.e., photons and electrons are no longer coupled.

*Ionization*

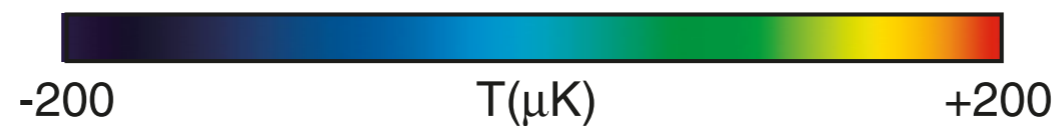
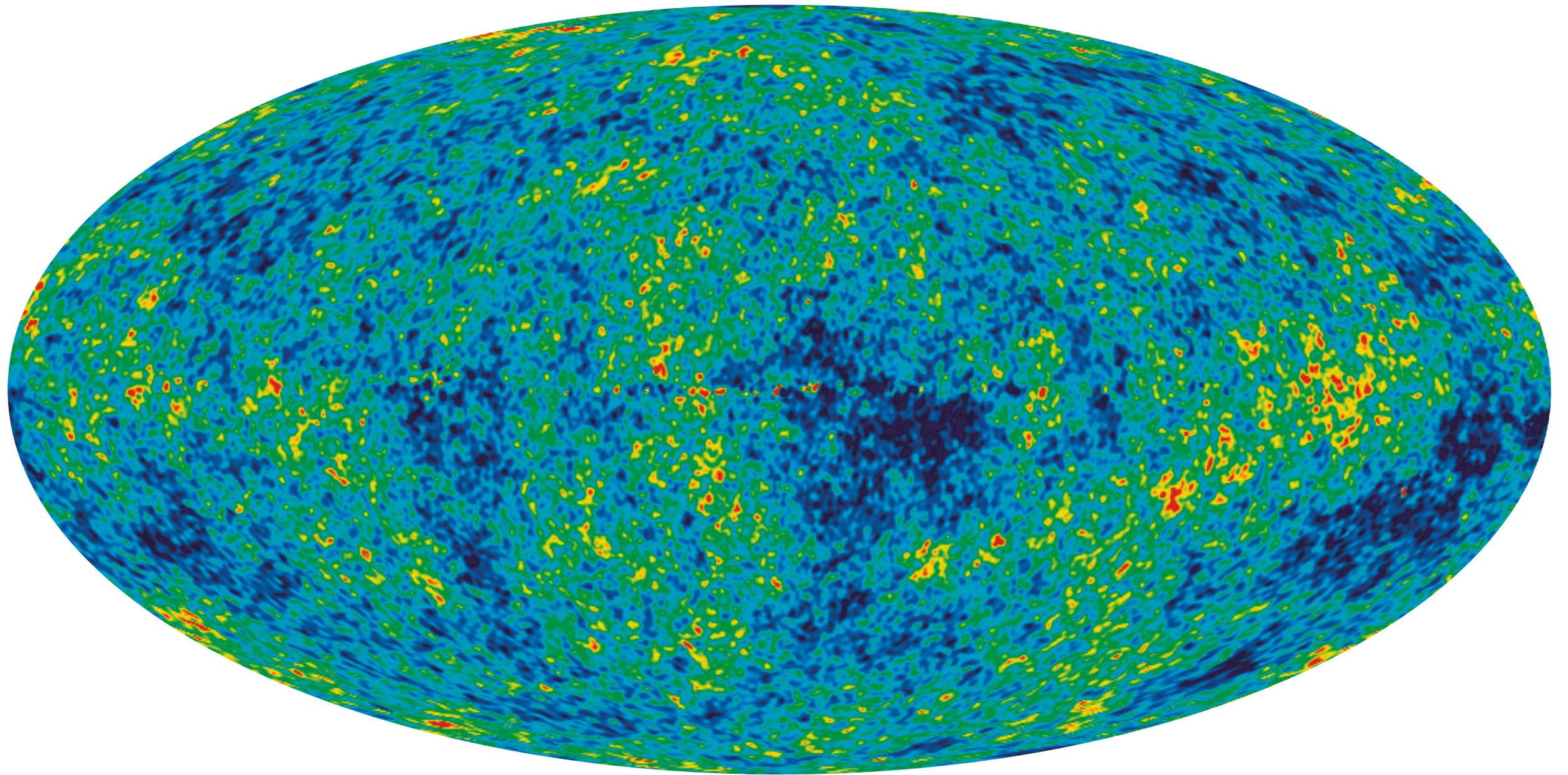


*Recombination*



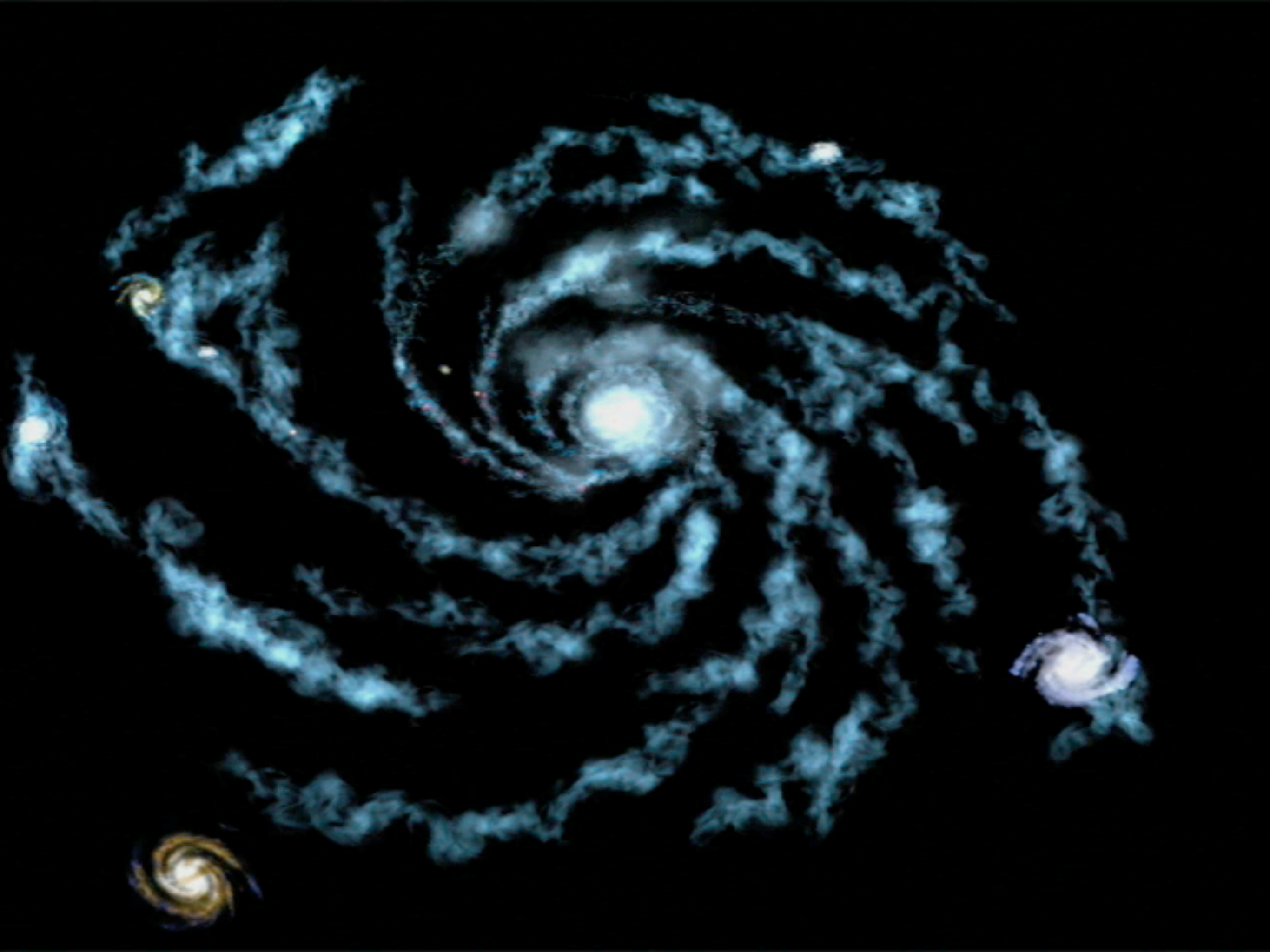


# A direct image of the Universe when it was 3000 K.

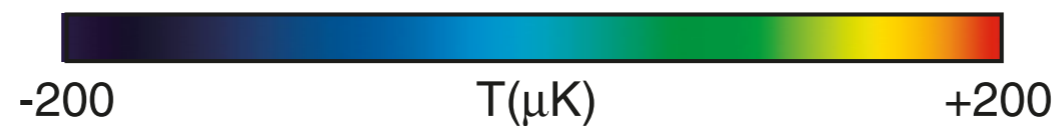
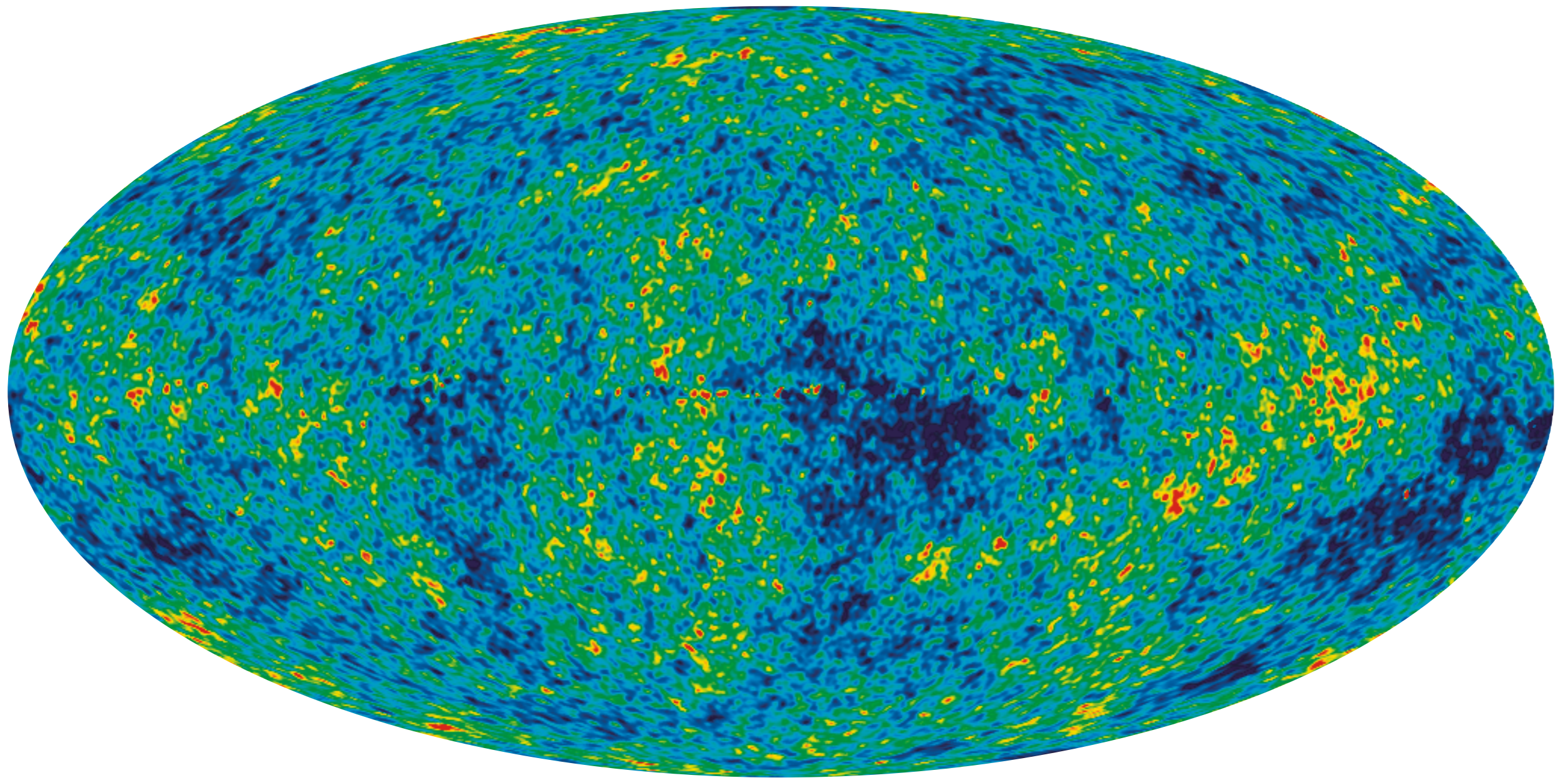


WMAP 5-year





# How were these ripples created?

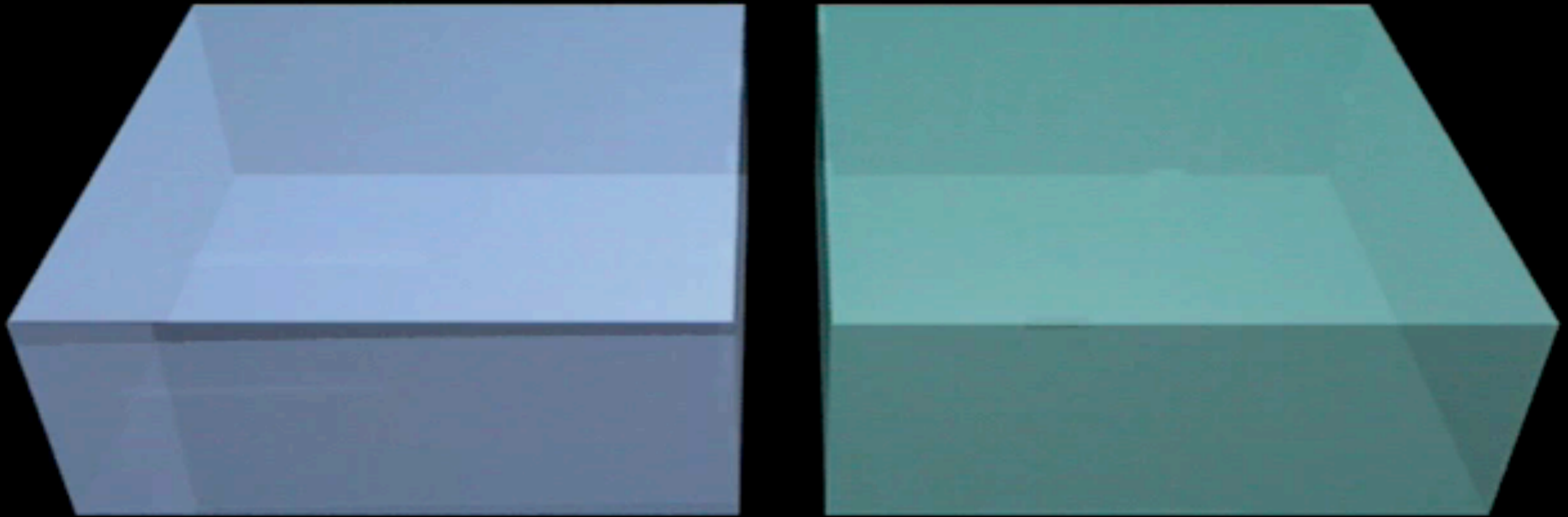


WMAP 5-year

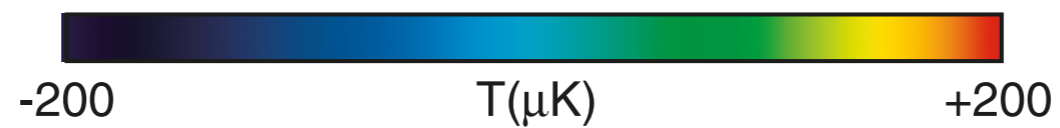
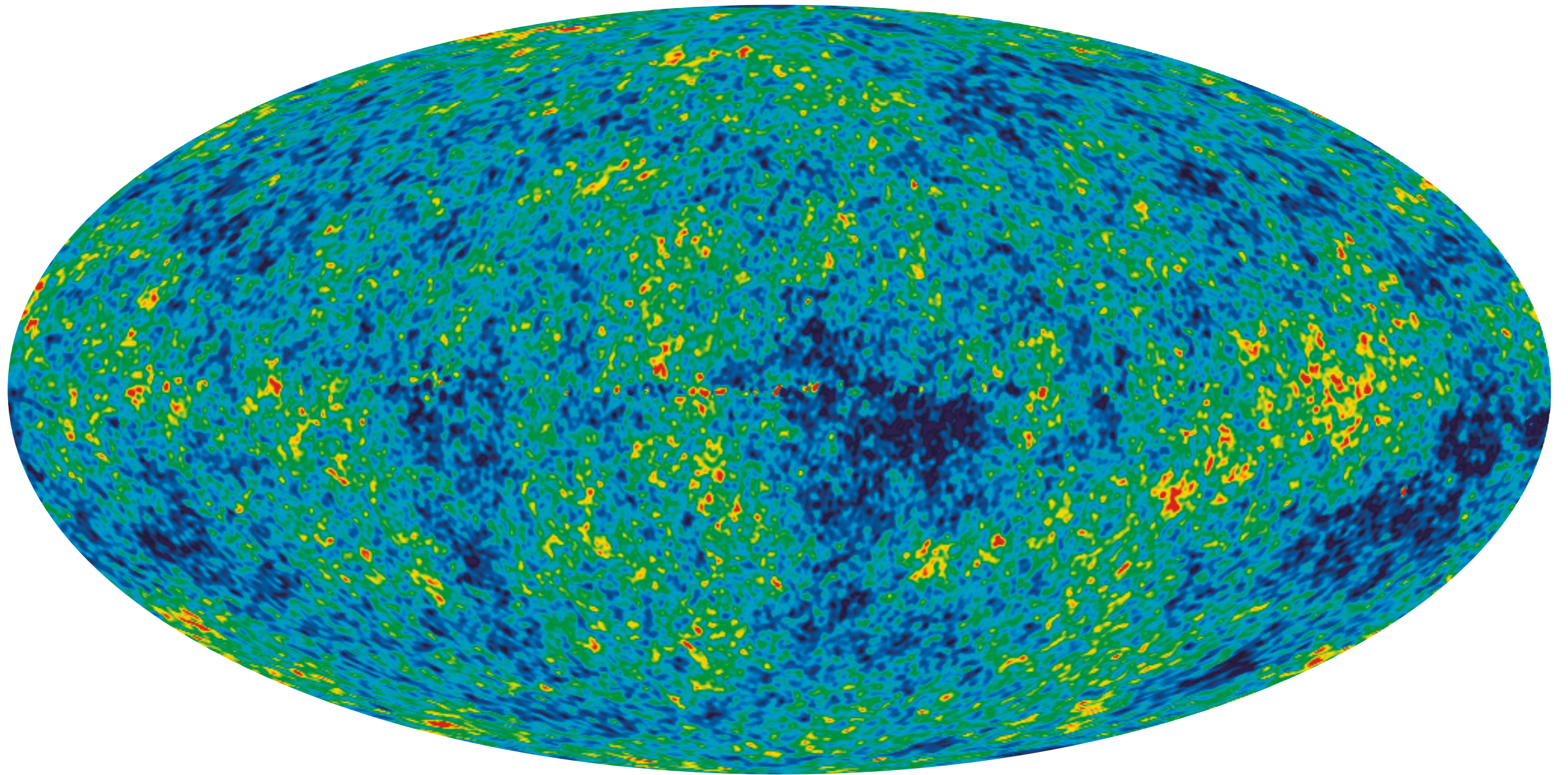
# Have you dropped potatoes in a soup?

- What would happen if you “perturb” the soup?

# The Cosmic Sound Wave

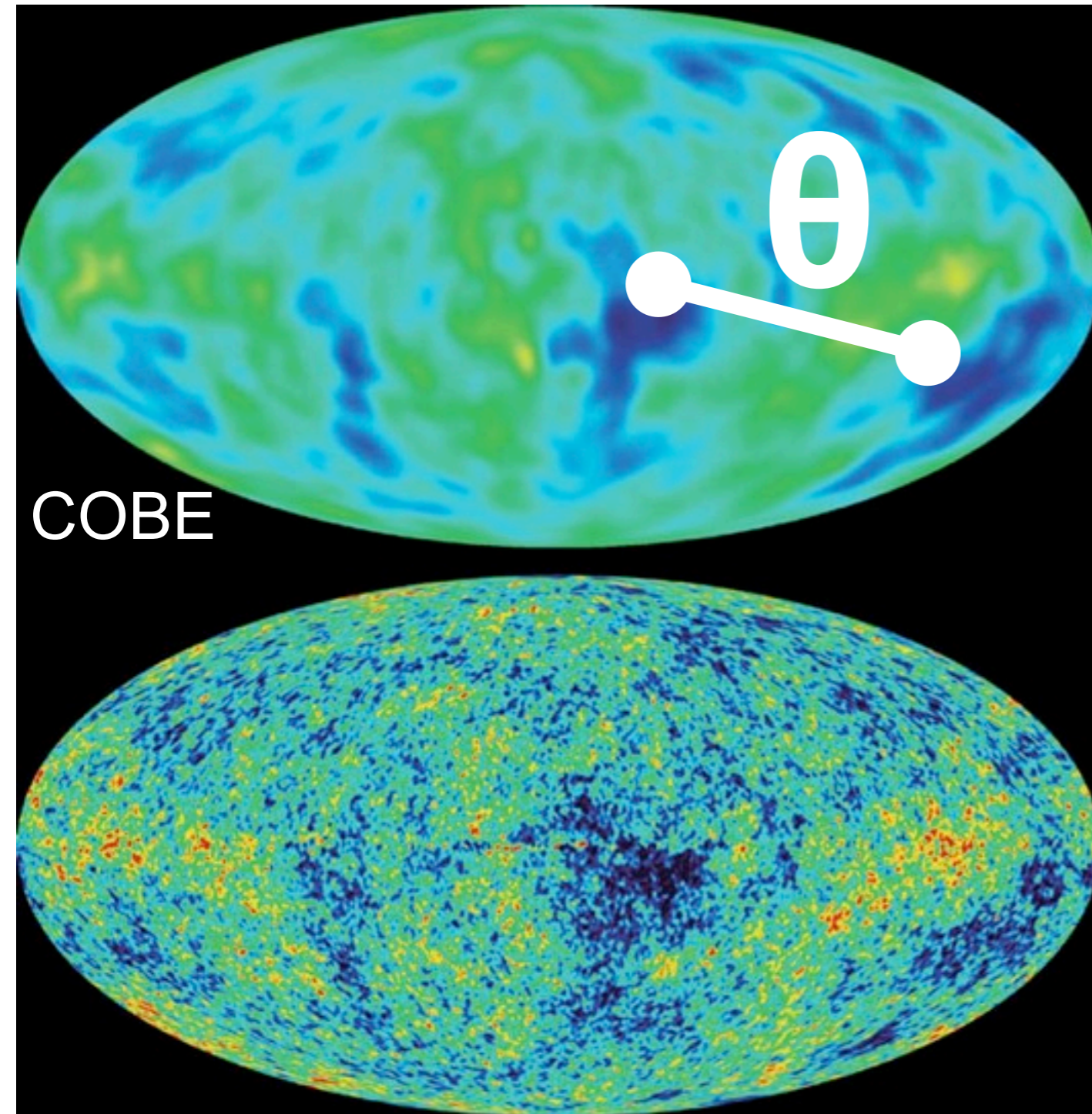


# Can You See the Sound Wave?

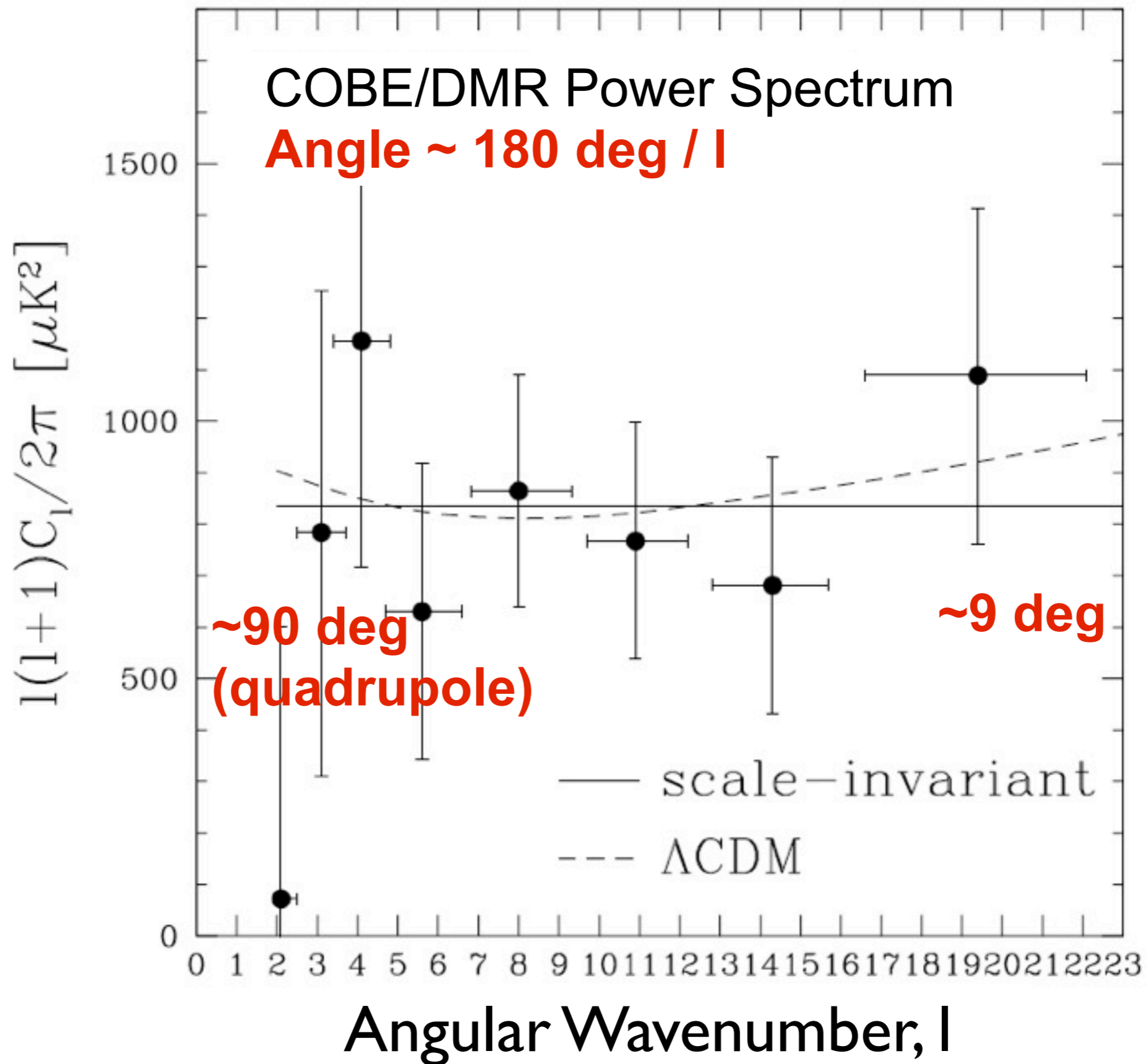


WMAP 5-year

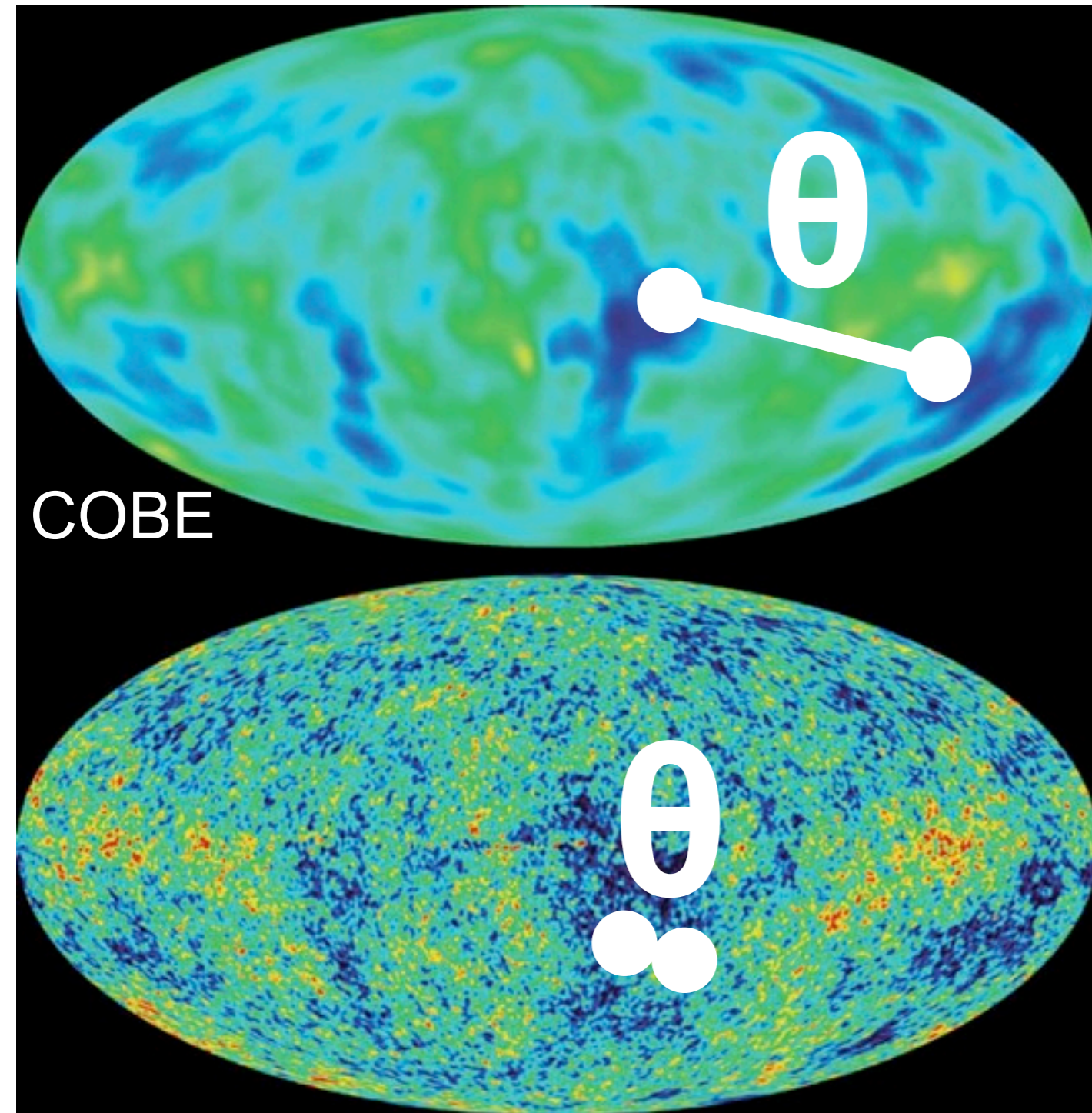
# Analysis: 2-point Correlation



- $C(\theta) = (1/4\pi) \sum (2l+1) C_l P_l(\cos\theta)$
- How are temperatures on two points on the sky, separated by  $\theta$ , are correlated?
- **“Power Spectrum,”**  $C_l$ 
  - How much fluctuation power do we have at a given angular scale?
  - $l \sim 180 \text{ degrees} / \theta$



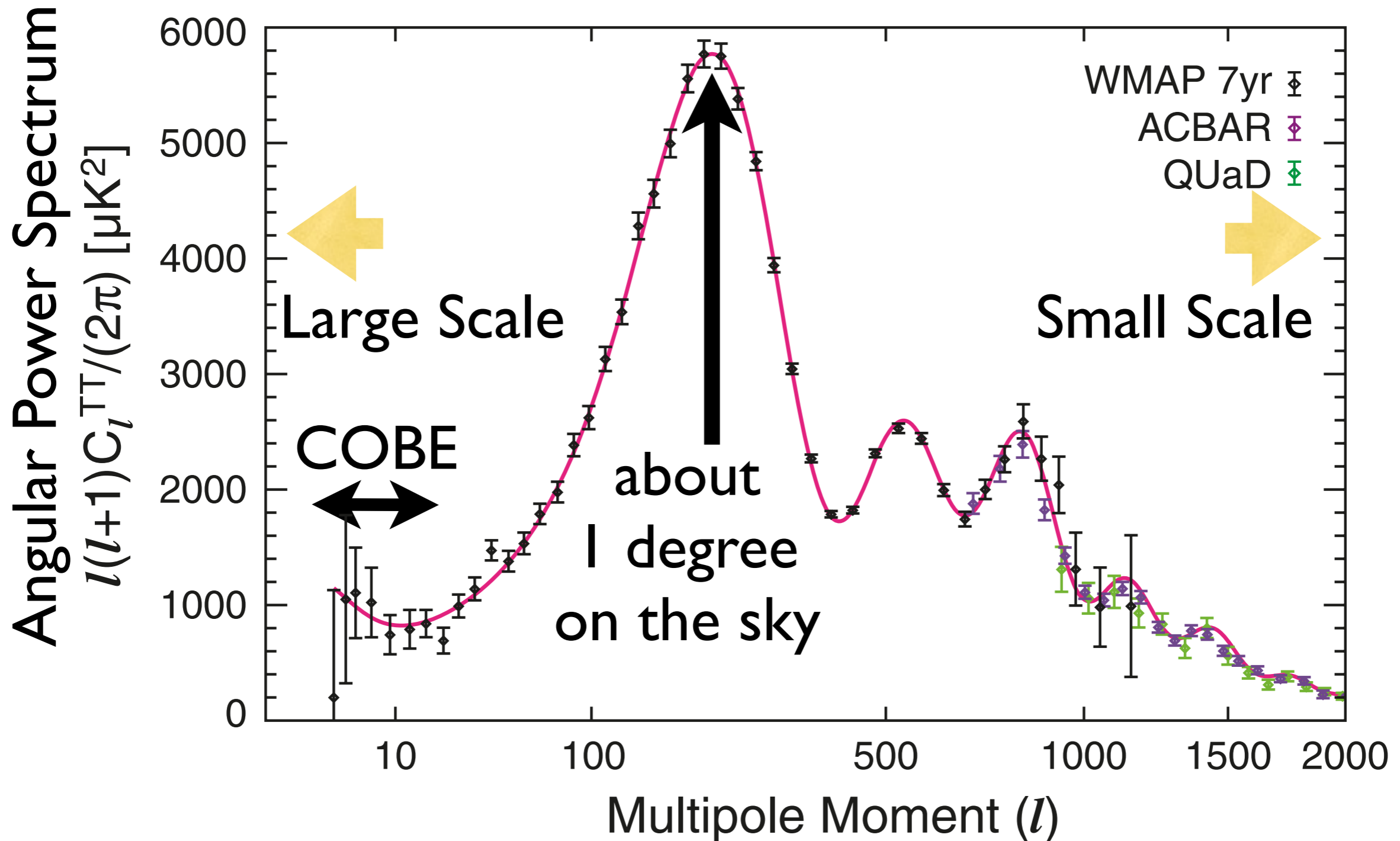
# COBE To WMAP



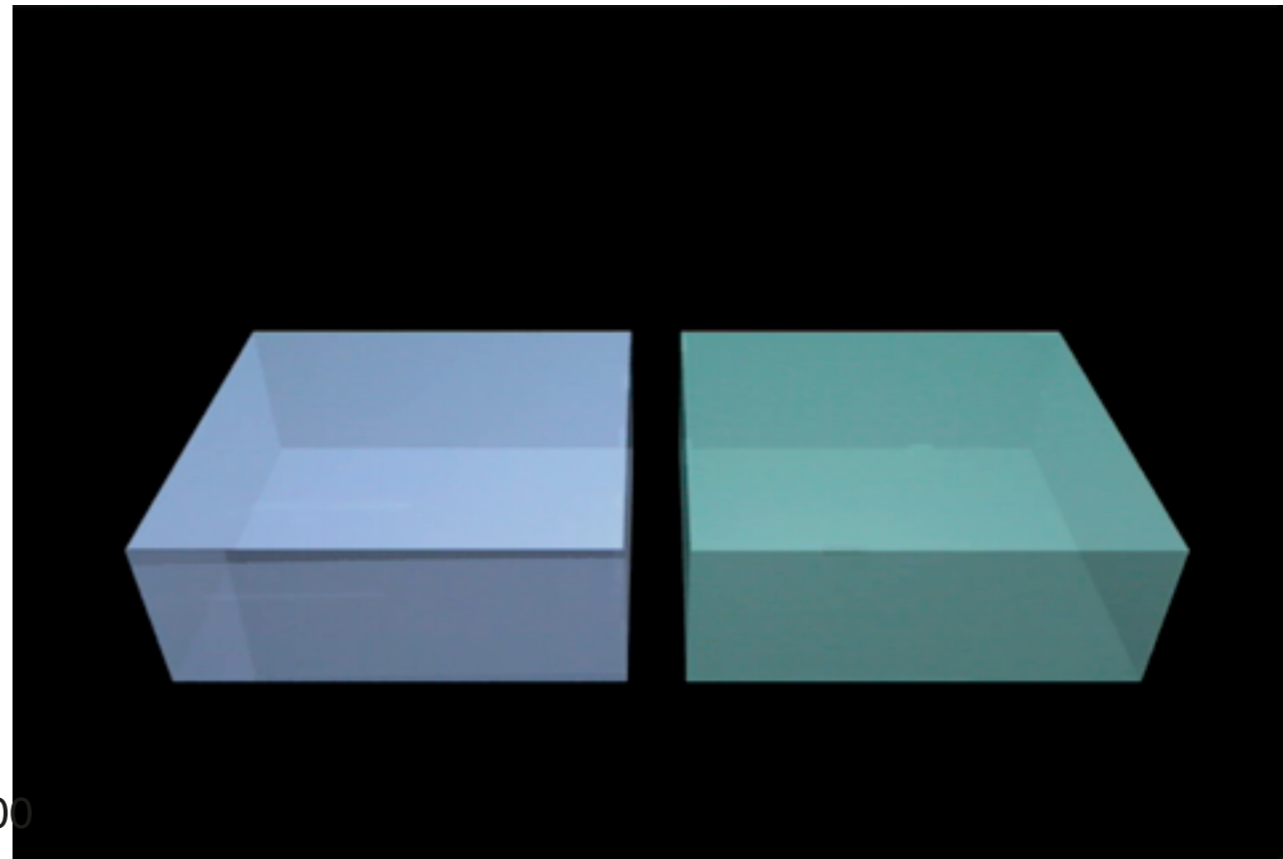
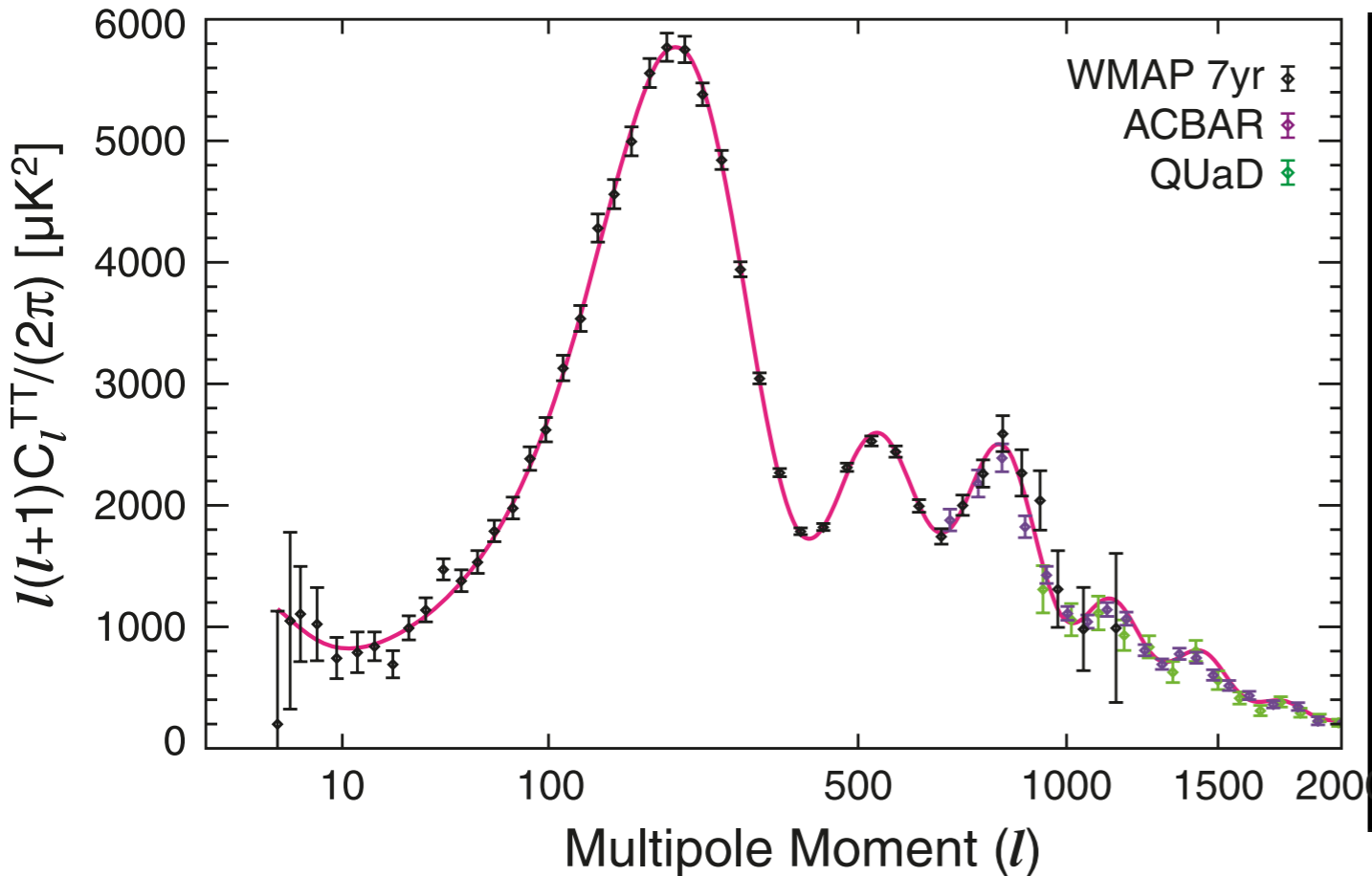
- COBE is unable to resolve the structures below  $\sim 7$  degrees
- WMAP's resolving power is 35 times better than COBE.
- What did WMAP see?



# WMAP Power Spectrum

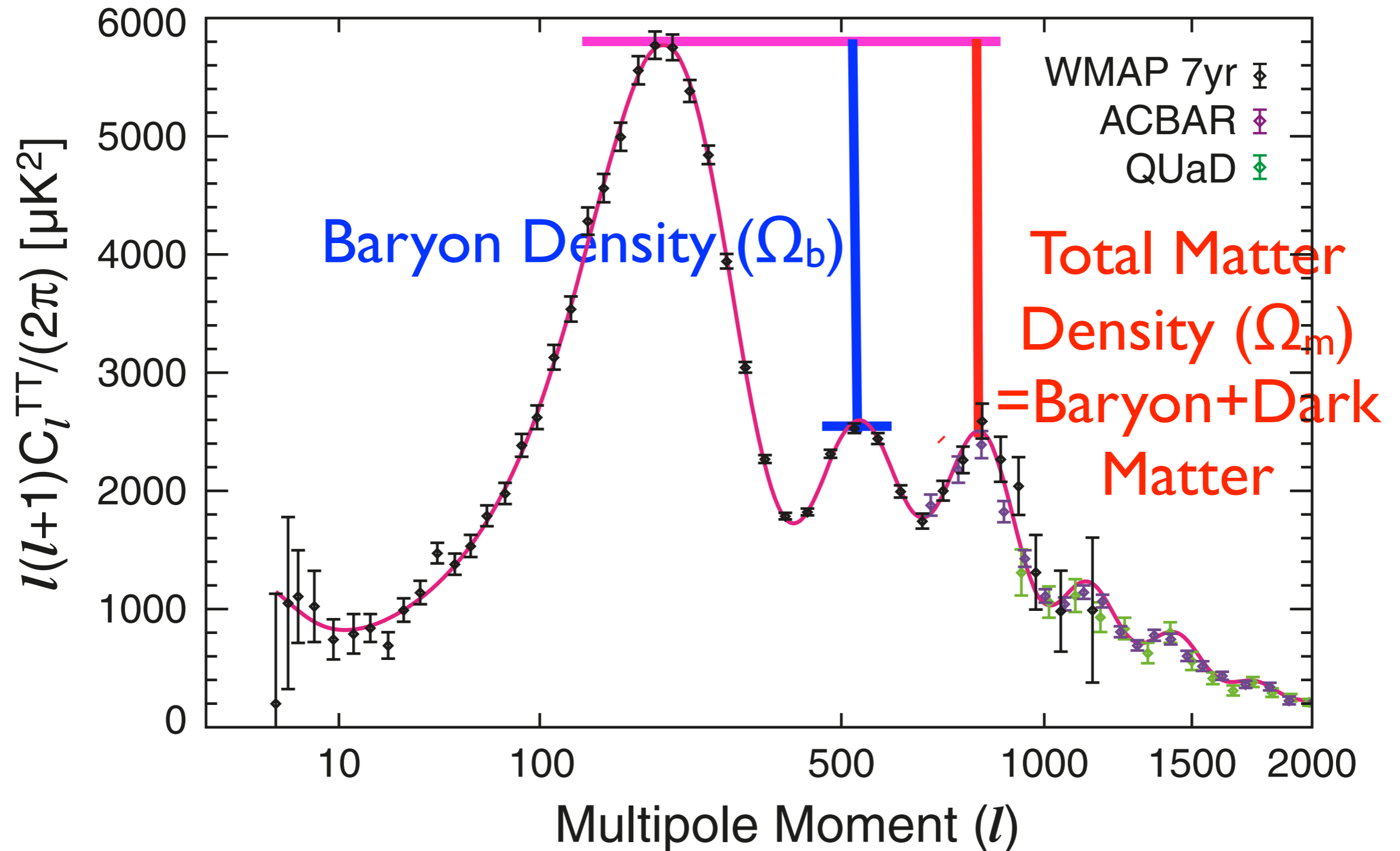


# The Cosmic Sound Wave



- *“The Universe as a potato soup”*
- *Main Ingredients: protons, helium nuclei, electrons, photons*
- *We measure the composition of the Universe by analyzing the wave form of the cosmic sound waves.*

# CMB to Baryon & Dark Matter



By “baryon,” I mean hydrogen and helium.

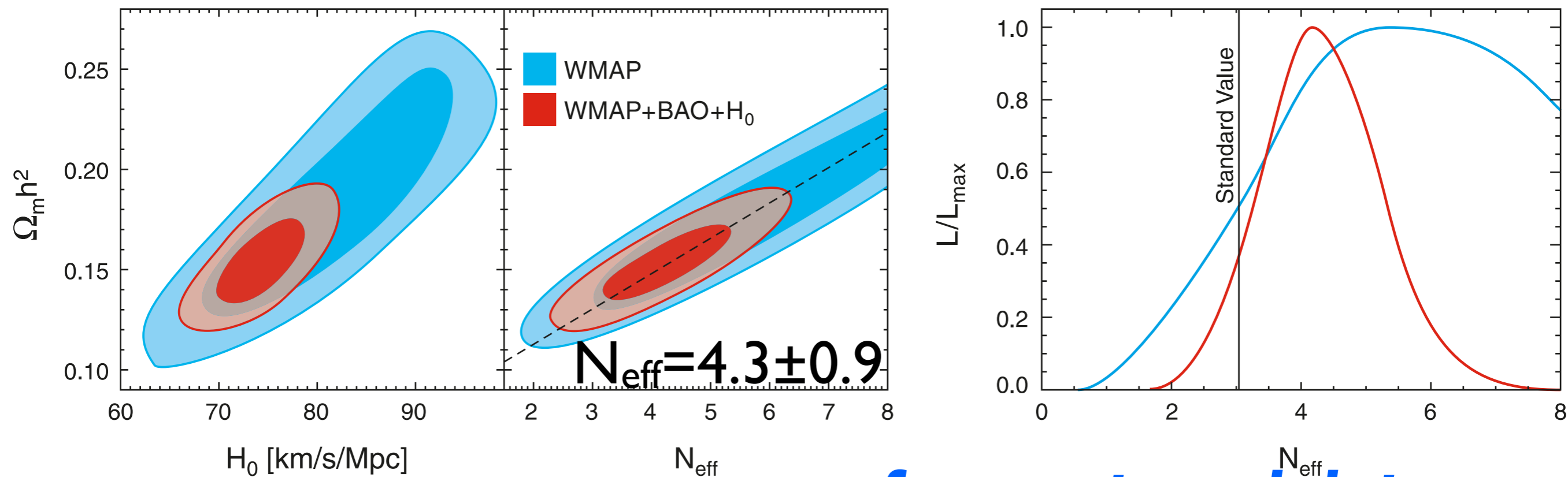
# Fundamental Observables from WMAP

- 1st-to-2nd peak ratio: “baryon-to-photon ratio,”  $\rho_B/\rho_\gamma$
- 1st-to-3rd peak ratio: “matter-to-radiation ratio,”  $\rho_M/\rho_R (=1+z_{EQ})$
- $\rho_M = \rho_B + \rho_{CDM}$
- $\rho_R = \rho_\gamma + \rho_\nu$
- If we assume that we know  $\rho_\nu$ , we can determine  $\rho_{CDM}$  from the 1st-to-3rd peak ratio; however, if we do not, we lose our ability to determine  $\rho_{CDM}$ !

# 3rd-peak “Spectroscopy”

- Total Matter = Baryons (H&He) + Dark Matter
- Total Radiation = Photons + Neutrinos (+new radiation)
  - Neutrino temperature =  $(4/11)^{1/3}$  Photon temperature
- So, for a given assumed value of the number of neutrino species (or the number of new radiation species, i.e., zero), we can measure the dark matter density.
- Or, we can get the dark matter density from elsewhere, and determine the number of radiation species!

# “3rd peak spectroscopy”: Number of Relativistic Species

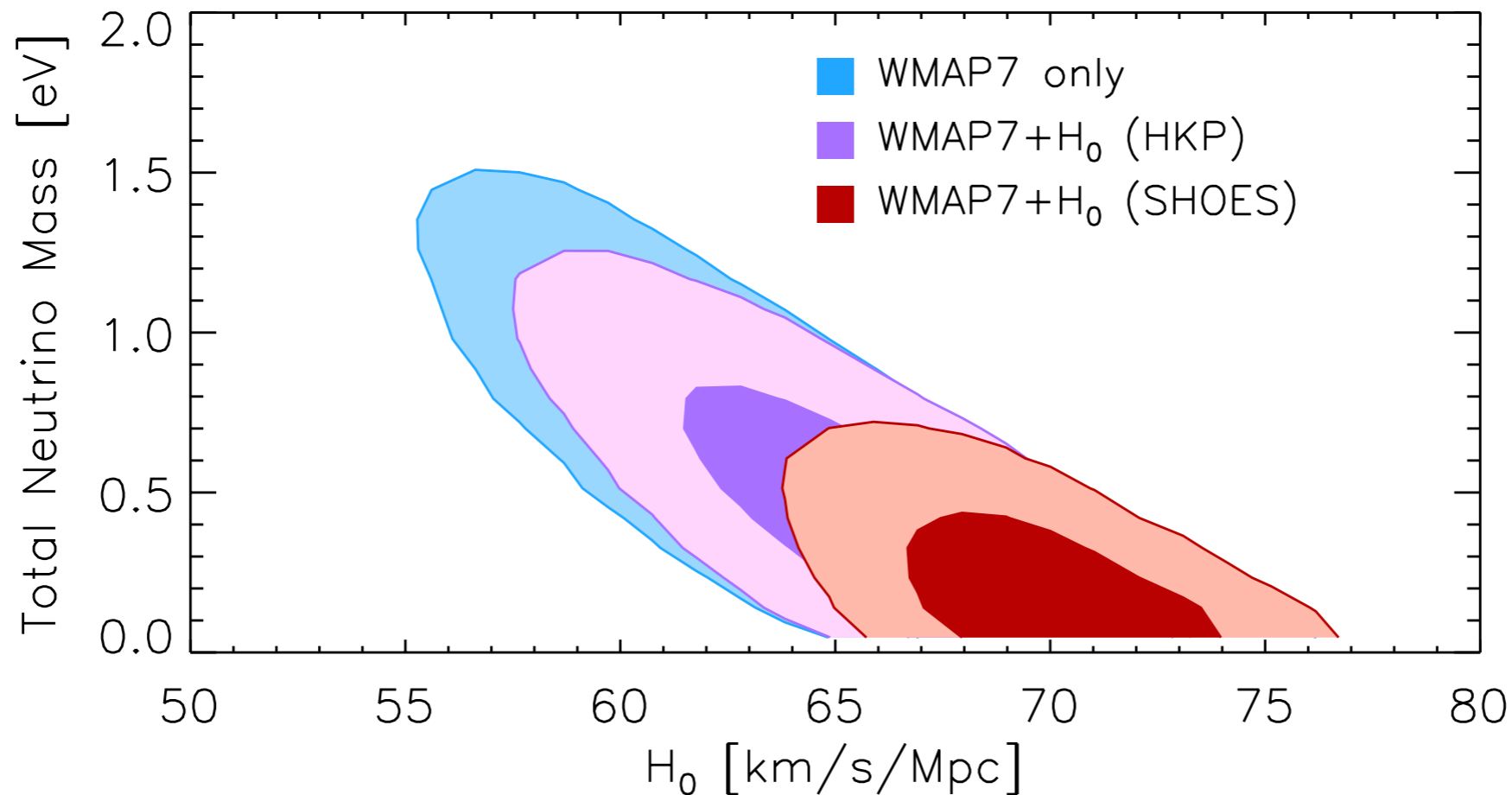


**from external data**

$$N_{\text{eff}} = 3.04 + 7.44 \left( \frac{\Omega_m h^2}{0.1308} \frac{3139}{1 + z_{\text{eq}}} - 1 \right)$$

**from 3rd peak**

# And, the mass of neutrinos



- WMAP data combined with the local measurement of the expansion rate ( $H_0$ ), we get  $\sum m_\nu < 0.6$  eV (95%CL)

# $\Sigma m_\nu$ from CMB alone

- There is a simple limit by which one can constrain  $\Sigma m_\nu$  using the primary CMB from  $z=1090$  alone (ignoring gravitational lensing of CMB by the intervening mass distribution)
- When all of neutrinos were lighter than  $\sim 0.6$  eV, they were still relativistic at the time of photon decoupling at  $z=1090$  (photon temperature  $3000\text{K}=0.26\text{eV}$ ).
  - $\langle E_\nu \rangle = 3.15(4/11)^{1/3} T_{\text{photon}} = 0.58$  eV
- Neutrino masses didn't matter if they were relativistic!
- For degenerate neutrinos,  $\Sigma m_\nu = 3.04 \times 0.58 = 1.8$  eV
  - **If  $\Sigma m_\nu \ll 1.8\text{eV}$ , CMB alone cannot see it**



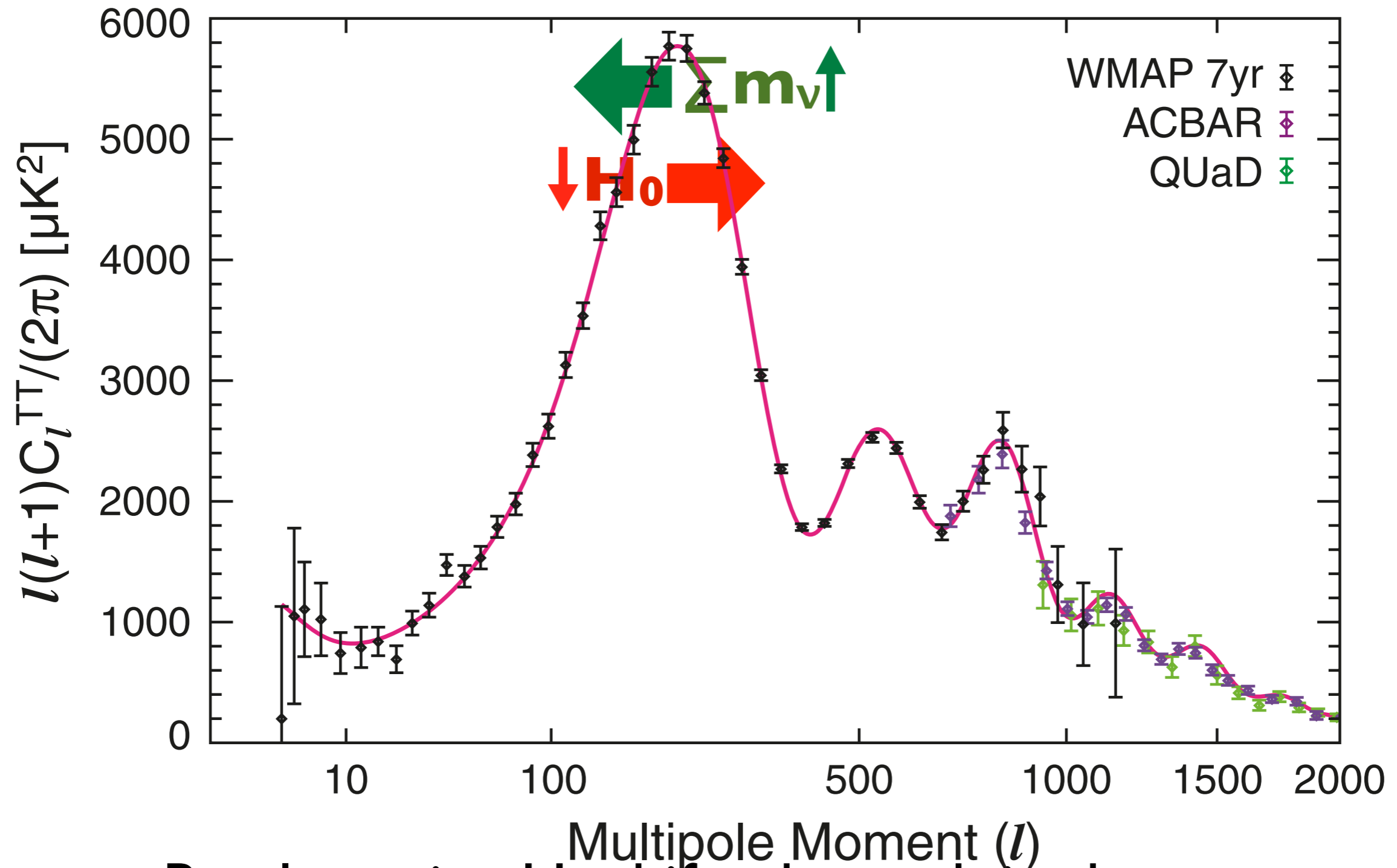
# Neutrino Subtlety

- For  $\sum m_\nu \ll 1.8\text{eV}$ , neutrinos were relativistic at  $z=1090$
- But, we know that  $\sum m_\nu > 0.05\text{eV}$  from neutrino oscillation experiments
- This means that **neutrinos are definitely non-relativistic today!**
- So, today's value of  $\Omega_M$  is the sum of baryons, CDM, and neutrinos:  $\Omega_M h^2 = (\Omega_B + \Omega_{\text{CDM}}) h^2 + 0.0106(\sum m_\nu / 1\text{eV})$

# Matter-Radiation Equality

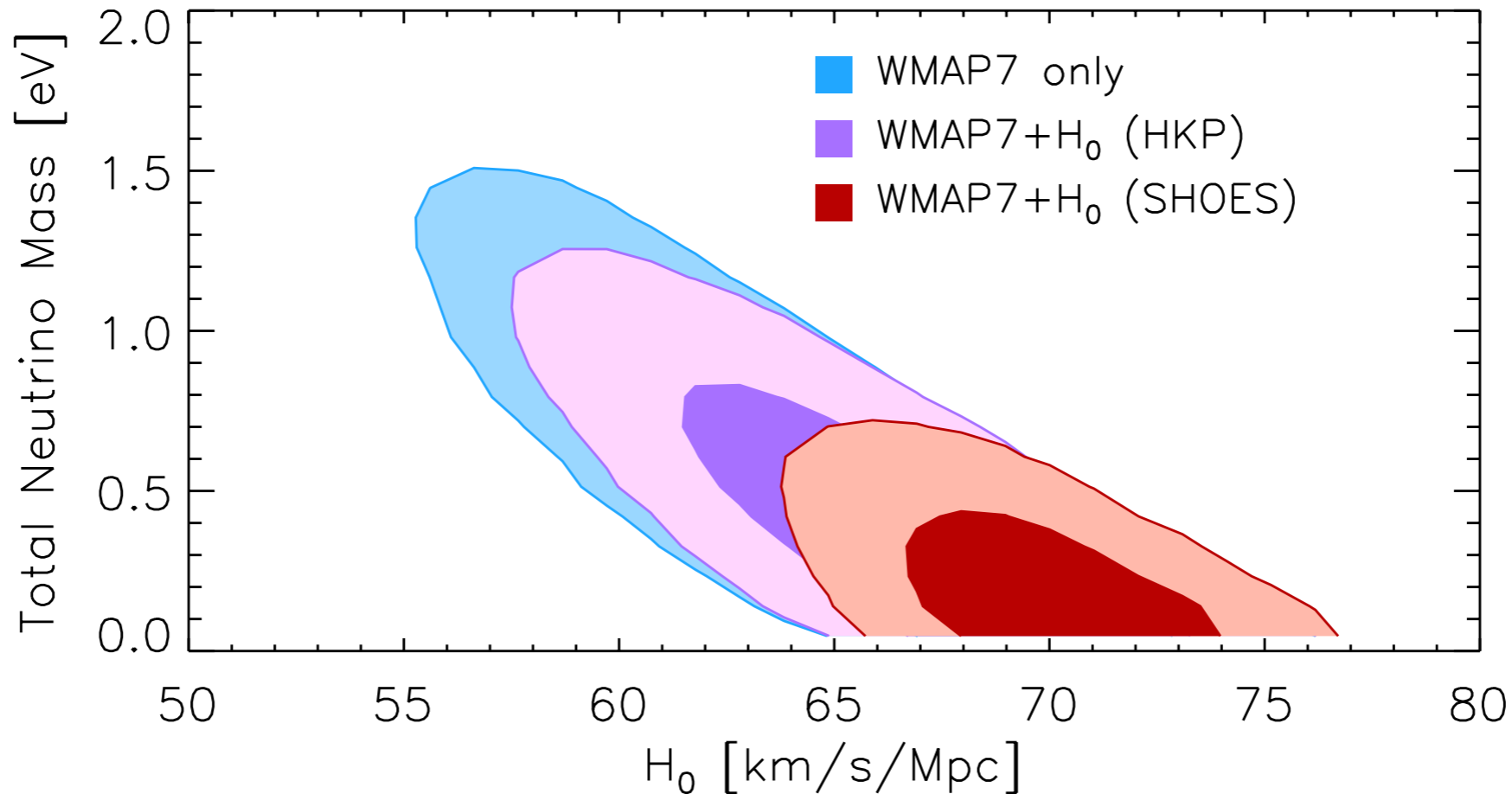
- However, since neutrinos were relativistic before  $z=1090$ , the matter-radiation equality is determined by:
  - $1+z_{\text{EQ}} = (\Omega_B + \Omega_{\text{CDM}}) / \Omega_R$
  - Now, recall  $\Omega_M h^2 = (\Omega_B + \Omega_{\text{CDM}}) h^2 + 0.0106(\sum m_\nu / 1 \text{ eV})$
  - For a given  $\Omega_M h^2$  constrained by the other data, adding  $\sum m_\nu$  makes  $(\Omega_B + \Omega_{\text{CDM}}) h^2$  smaller  $\rightarrow$  smaller  $z_{\text{EQ}}$   $\rightarrow$  **Radiation Era lasts longer**
- **This effect shifts the first peak to a lower multipole**

# $\Sigma m_\nu$ : Shifting the Peak To Low- $l$



- But, lowering  $H_0$  shifts the peak in the opposite direction. So...

# Shift of Peak Absorbed by $H_0$



- Here is a catch:

- **Shift of the first peak to a lower multipole can be canceled by lowering  $H_0$ !  $\sum m_\nu < 0.6$  eV (95%CL)**

# How Do We Test Inflation?

- How can we answer a simple question like this:
  - “*How were primordial fluctuations generated?*”

# Stretching Micro to Macro

$H^{-1}$  = Hubble Size



Quantum fluctuations on microscopic scales



**INFLATION!**

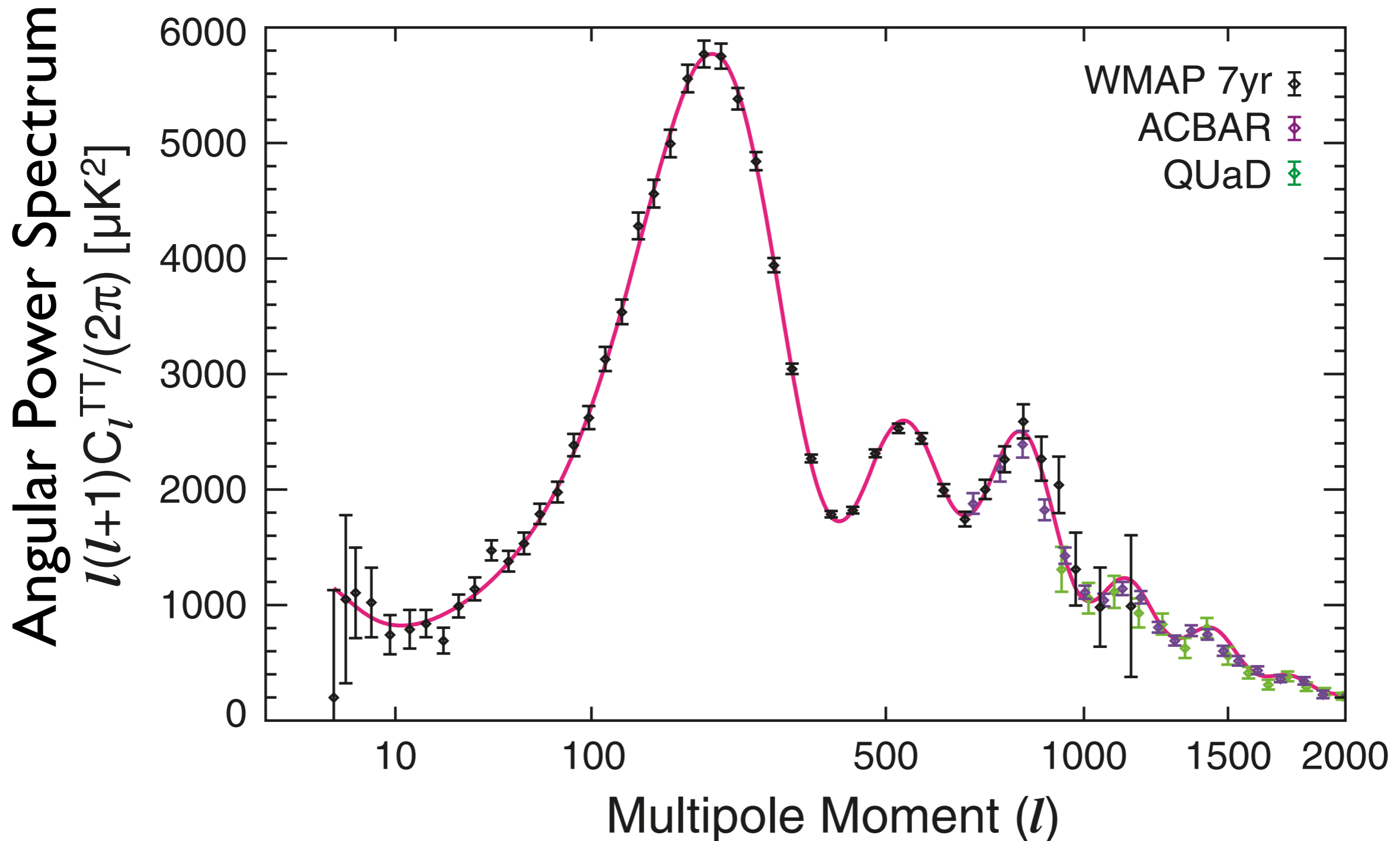


Quantum fluctuations cease to be quantum, and become<sup>38</sup> observable

# Power Spectrum

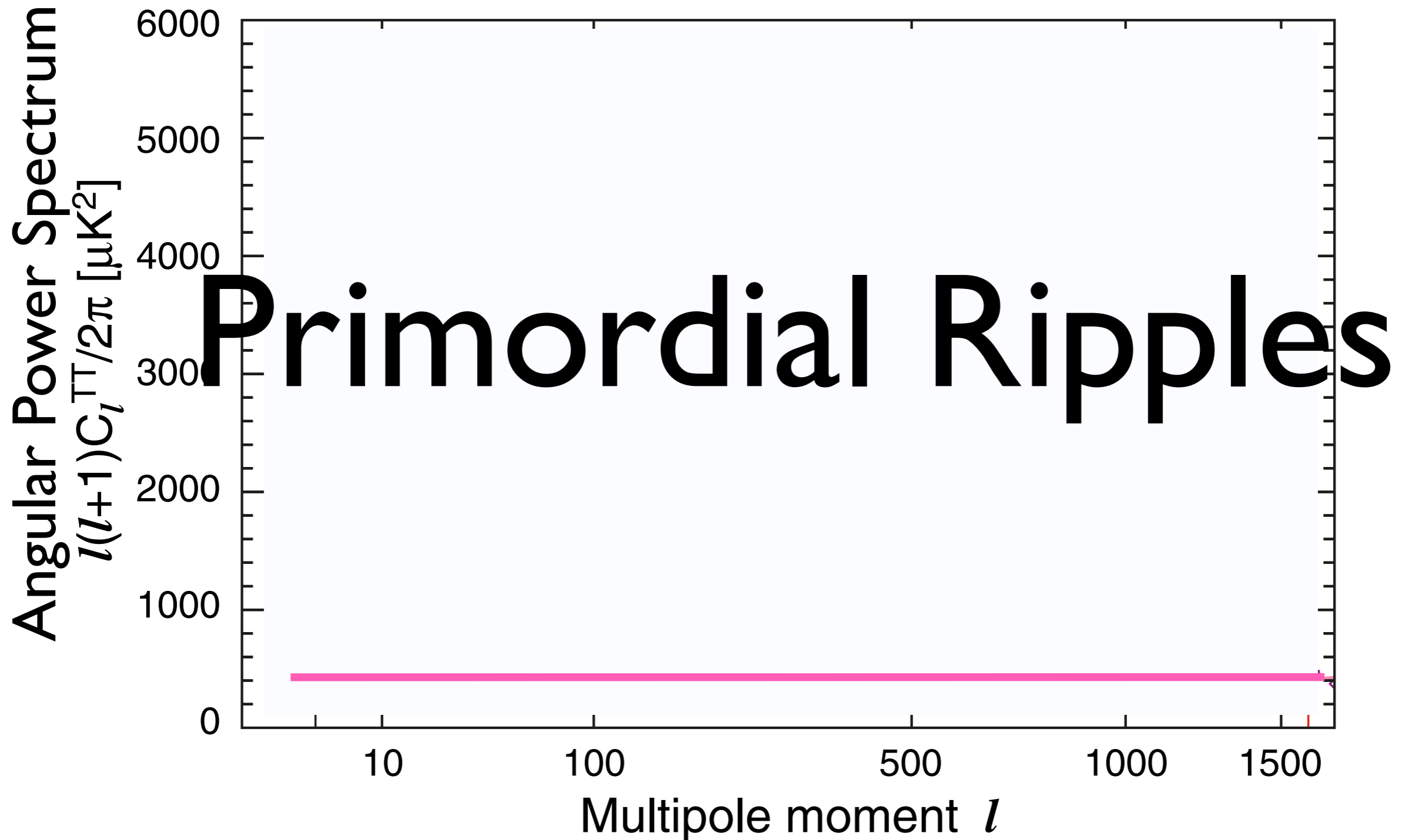
- A very successful explanation (Mukhanov & Chibisov; Guth & Pi; Hawking; Starobinsky; Bardeen, Steinhardt & Turner) is:
- *Primordial fluctuations were generated by quantum fluctuations of the scalar field that drove inflation.*
- The prediction: a nearly scale-invariant power spectrum in the curvature perturbation,  $\zeta = -(\dot{H} dt/d\varphi) \delta\varphi$ 
  - **$P_\zeta(\mathbf{k}) = \langle |\zeta_{\mathbf{k}}|^2 \rangle = A/k^{4-n_s} \sim A/k^3$**
  - where  $n_s \sim 1$  and  $A$  is a normalization.

# WMAP Power Spectrum

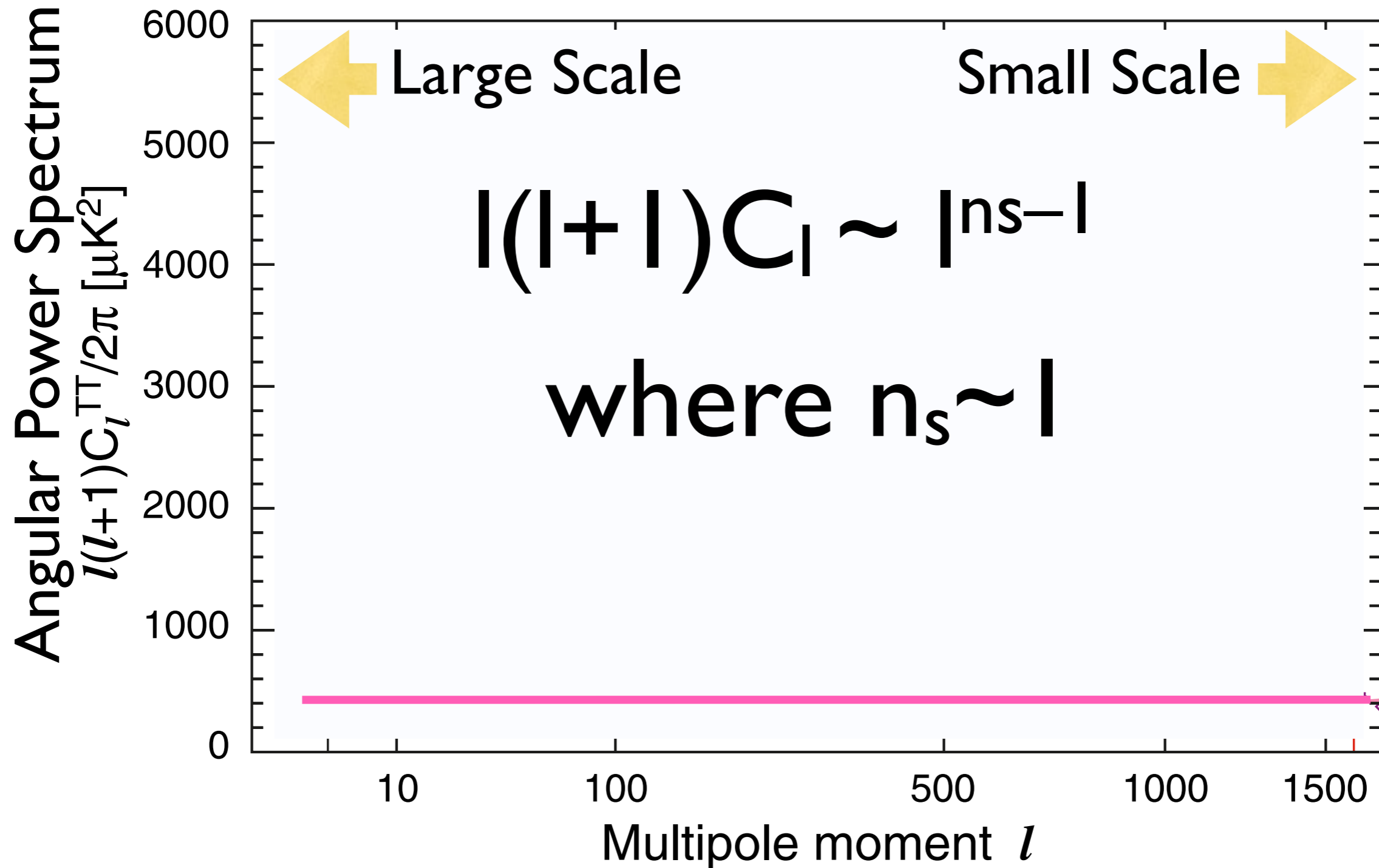




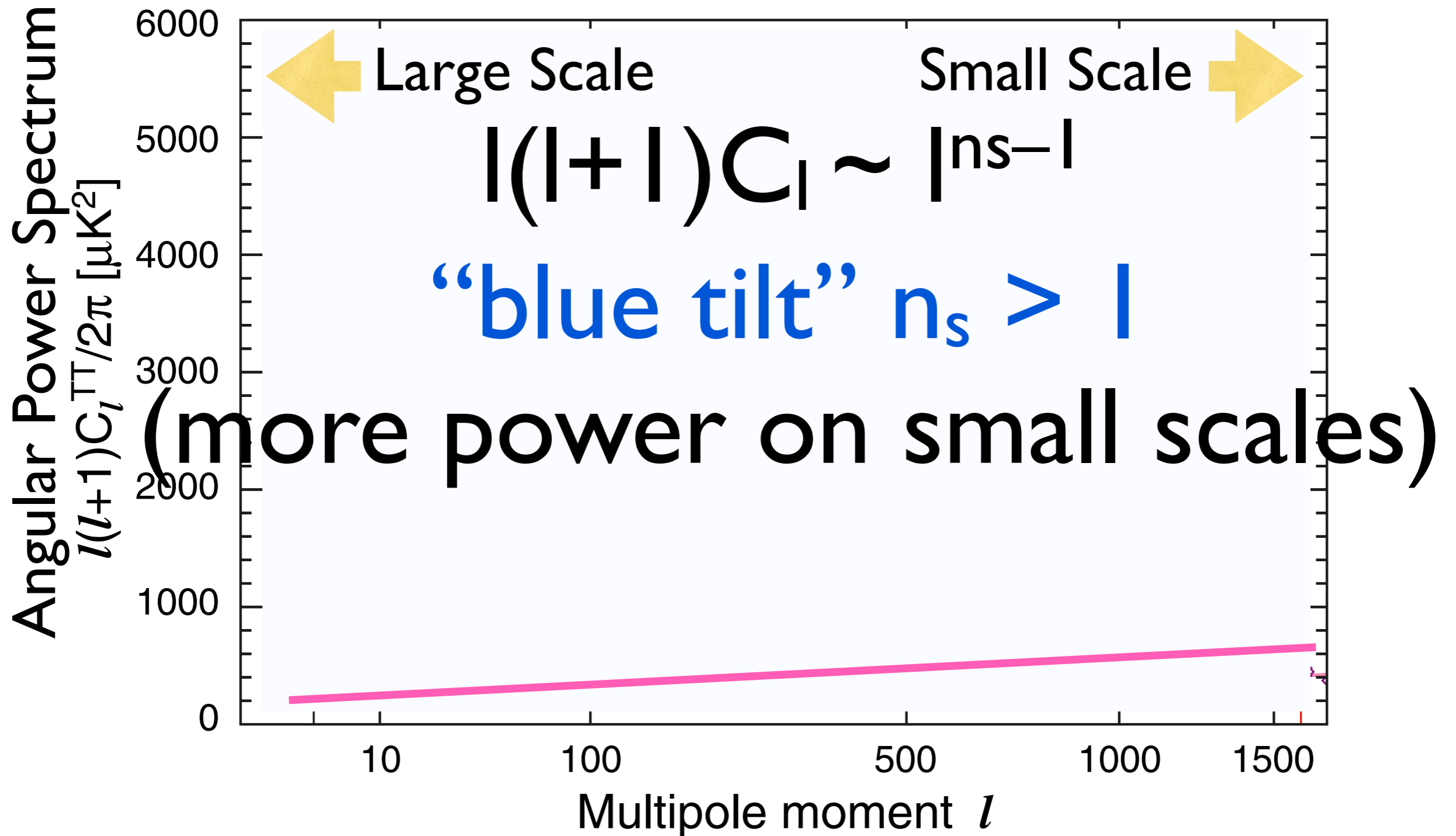
# Getting rid of the Sound Waves



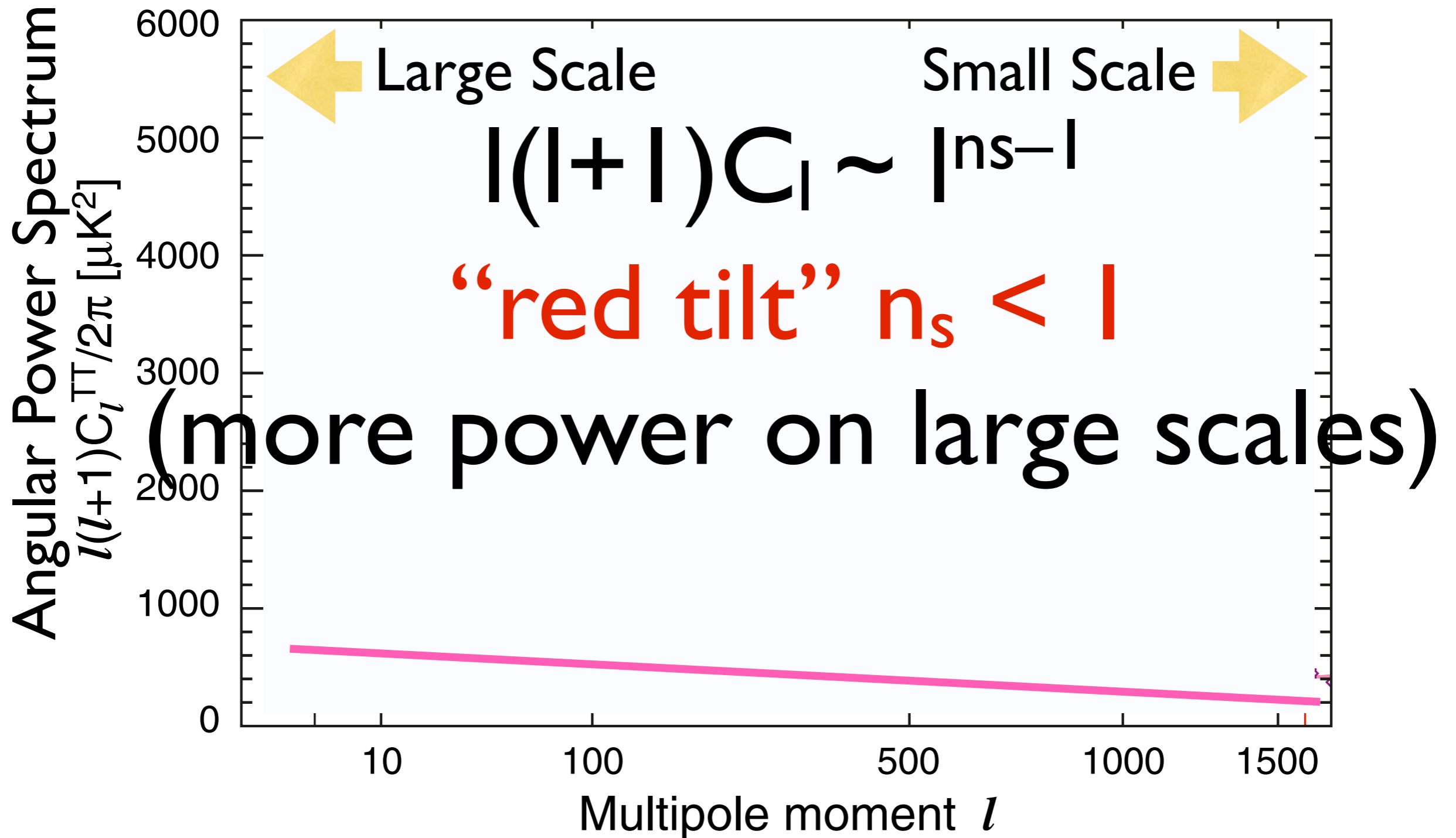
# Inflation Predicts:



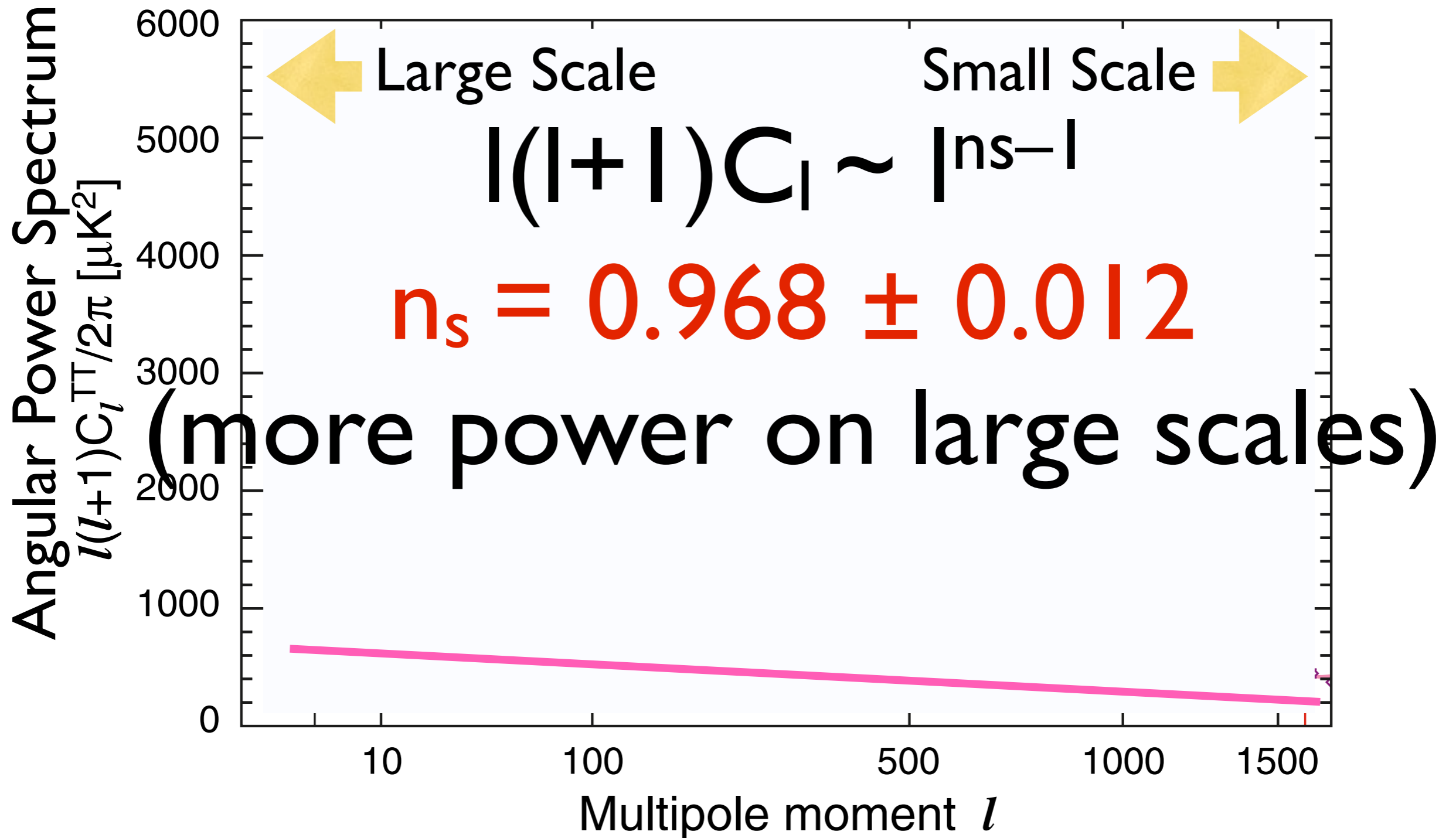
# Inflation may do this



...or this

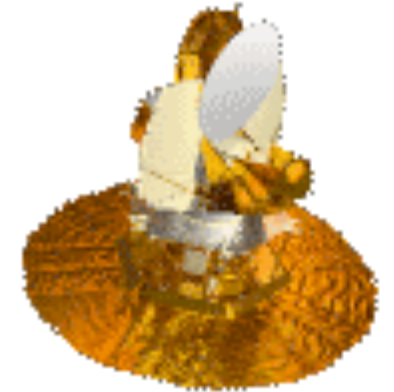


# WMAP 7-year Measurement (Komatsu et al. 2011)



After 9 years of observations...

# WMAP taught us:



- **All of the basic predictions of single-field and slow-roll inflation models are consistent with the data**
  - But, not all models are consistent (i.e.,  $\lambda\phi^4$  is out unless you introduce a non-minimal coupling)

# Testing Single-field by Adiabaticity

- Within the context of single-field inflation, all the matter and radiation originated from a single field, and thus there is a particular relation (adiabatic relation) between the perturbations in matter and photons:

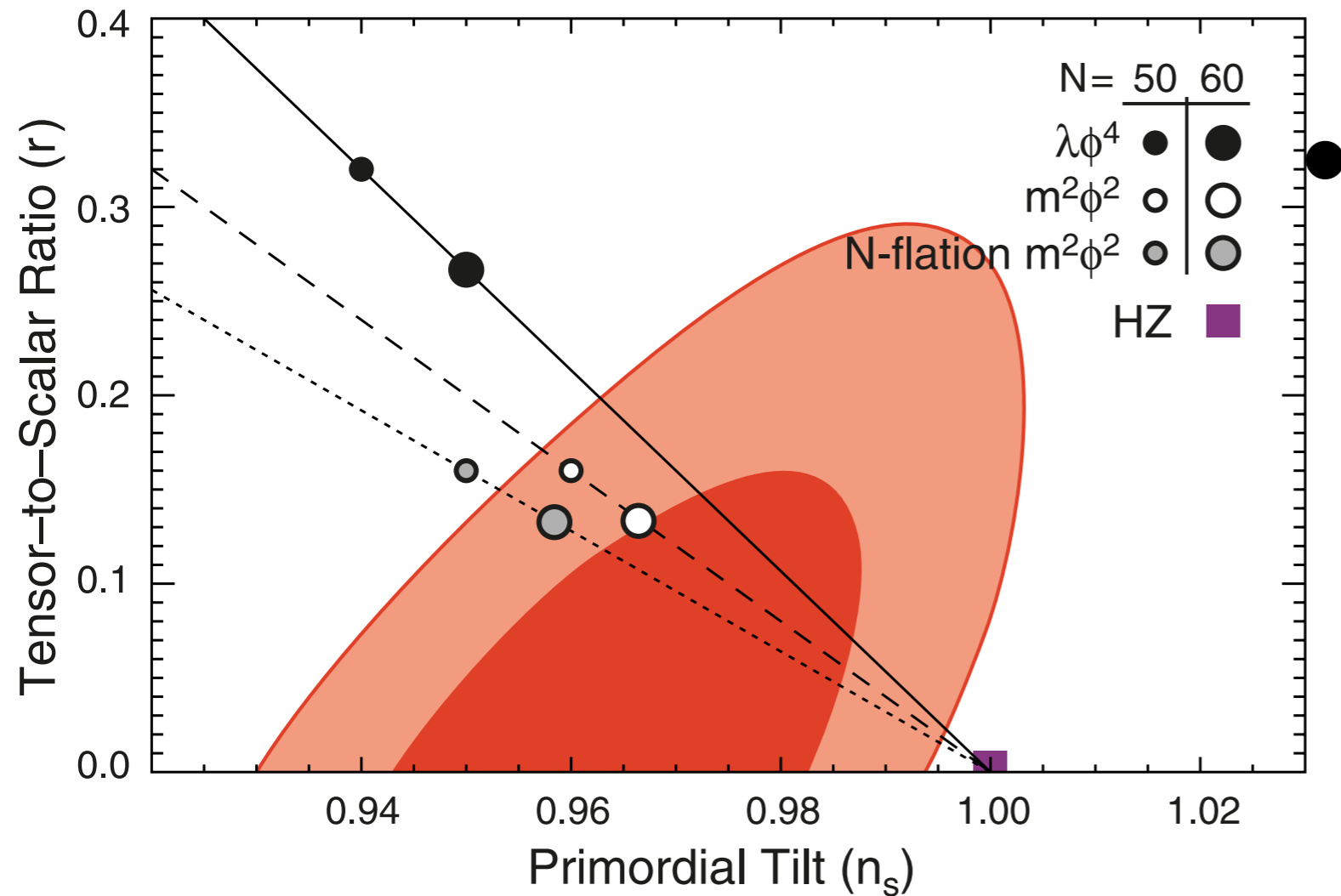
$$S_{c,\gamma} \equiv \frac{\delta\rho_c}{\rho_c} - \frac{3\delta\rho_\gamma}{4\rho_\gamma} = 0$$

The data are consistent with

$S=0$ :

$$\frac{|\delta\rho_c/\rho_c - 3\delta\rho_\gamma/(4\rho_\gamma)|}{\frac{1}{2}[\delta\rho_c/\rho_c + 3\delta\rho_\gamma/(4\rho_\gamma)]} < 0.09 \quad (95\% \text{ CL})$$

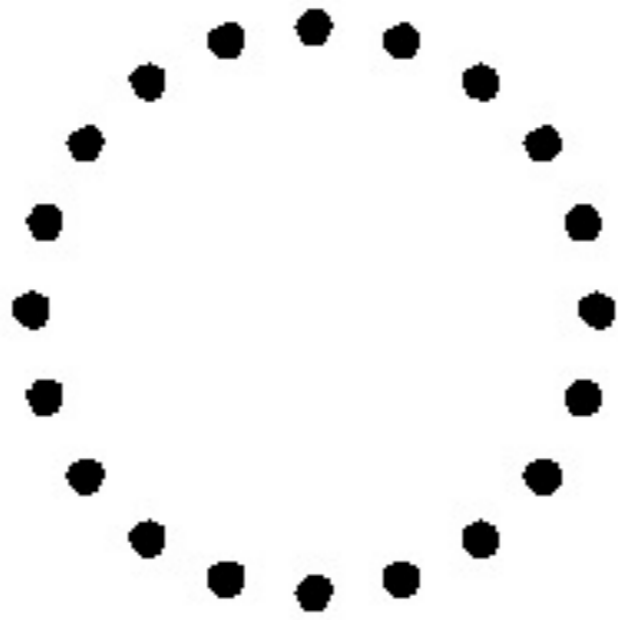
# Inflation looks good



- Joint constraint on the primordial tilt,  $n_s$ , and the tensor-to-scalar ratio,  $r$ .
- **$r < 0.24$  (95%CL; WMAP7+BAO+ $H_0$ )**

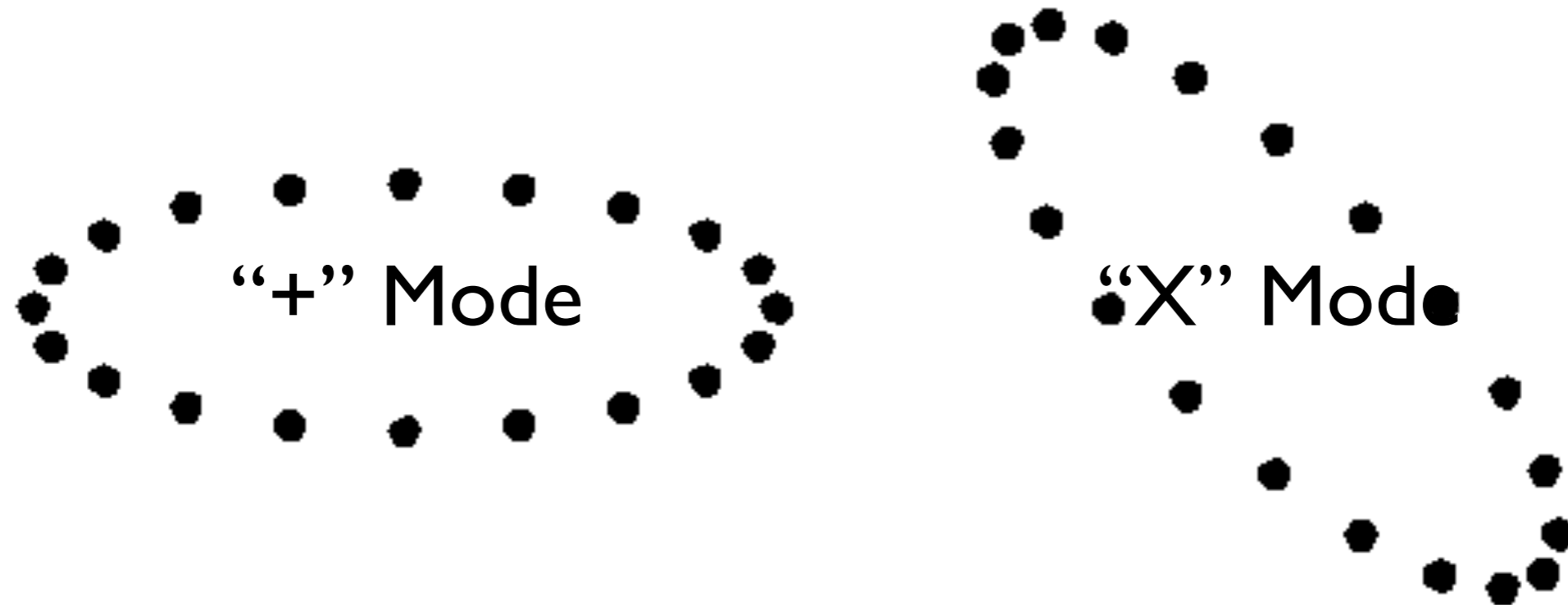


**Gravitational waves are coming toward you... What do you do?**



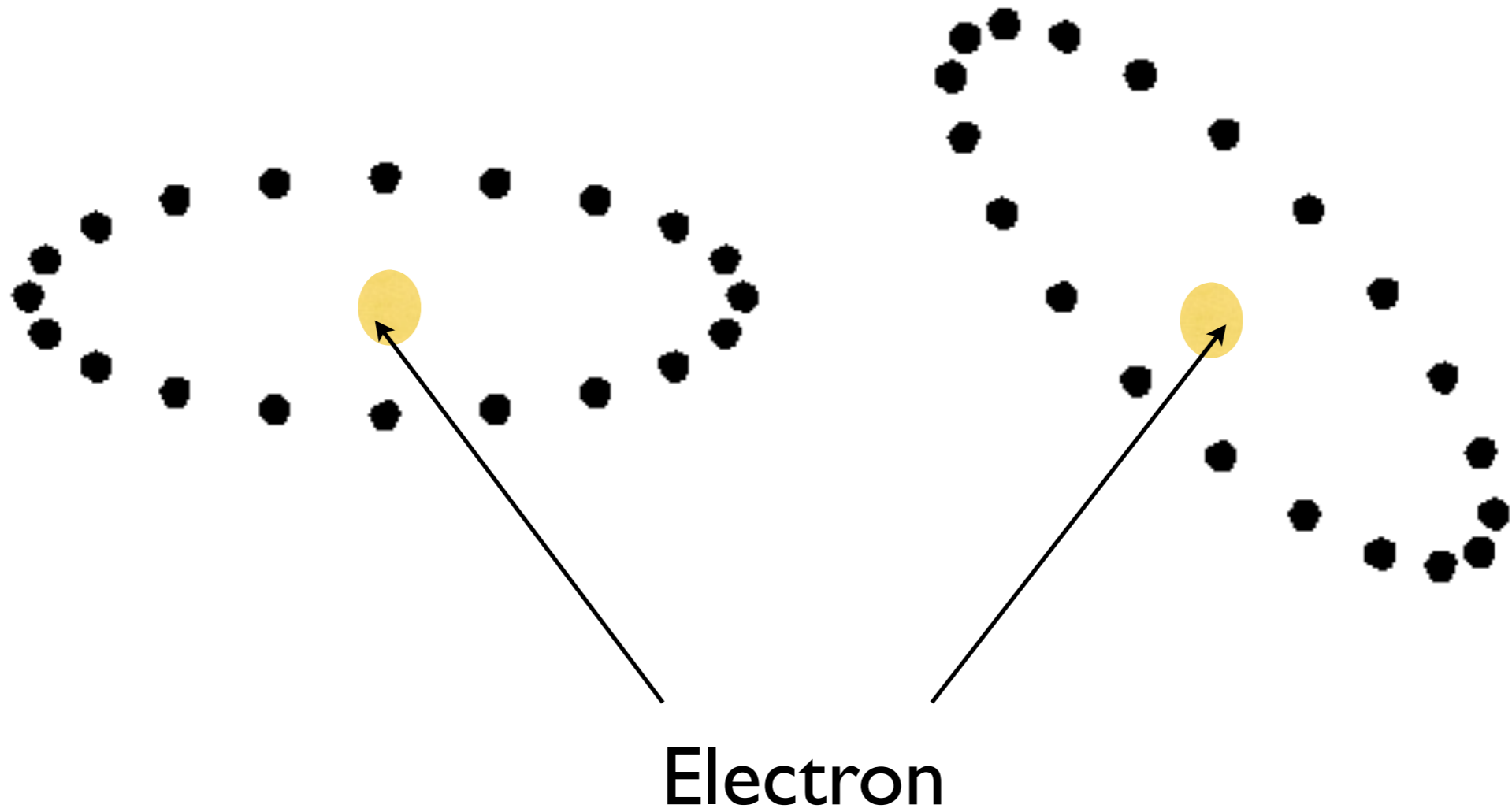
- **Gravitational waves stretch space, causing particles to move.**

# Two Polarization States of GW



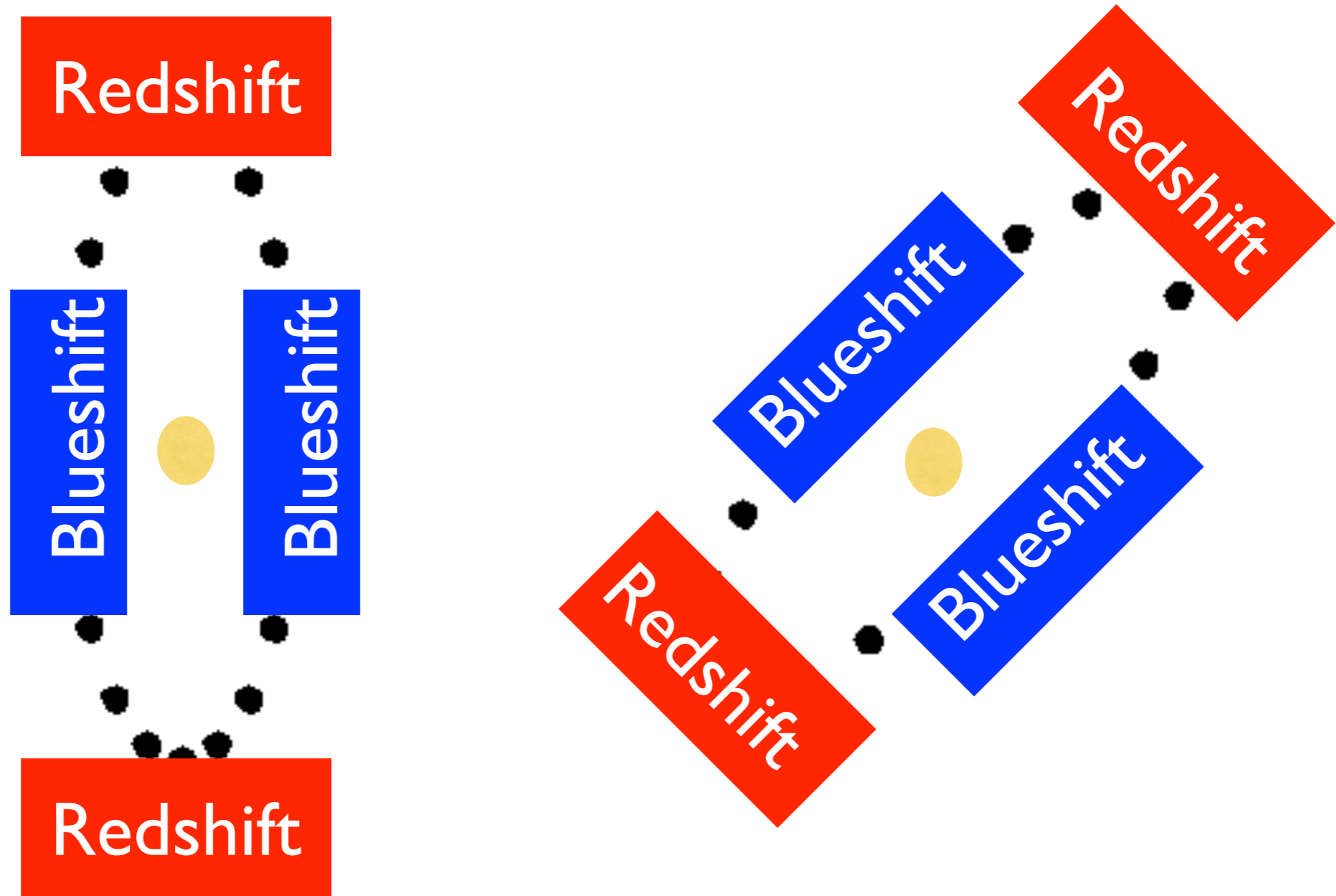
- This is great - this will automatically generate quadrupolar temperature anisotropy around electrons!

# From GW to CMB Polarization

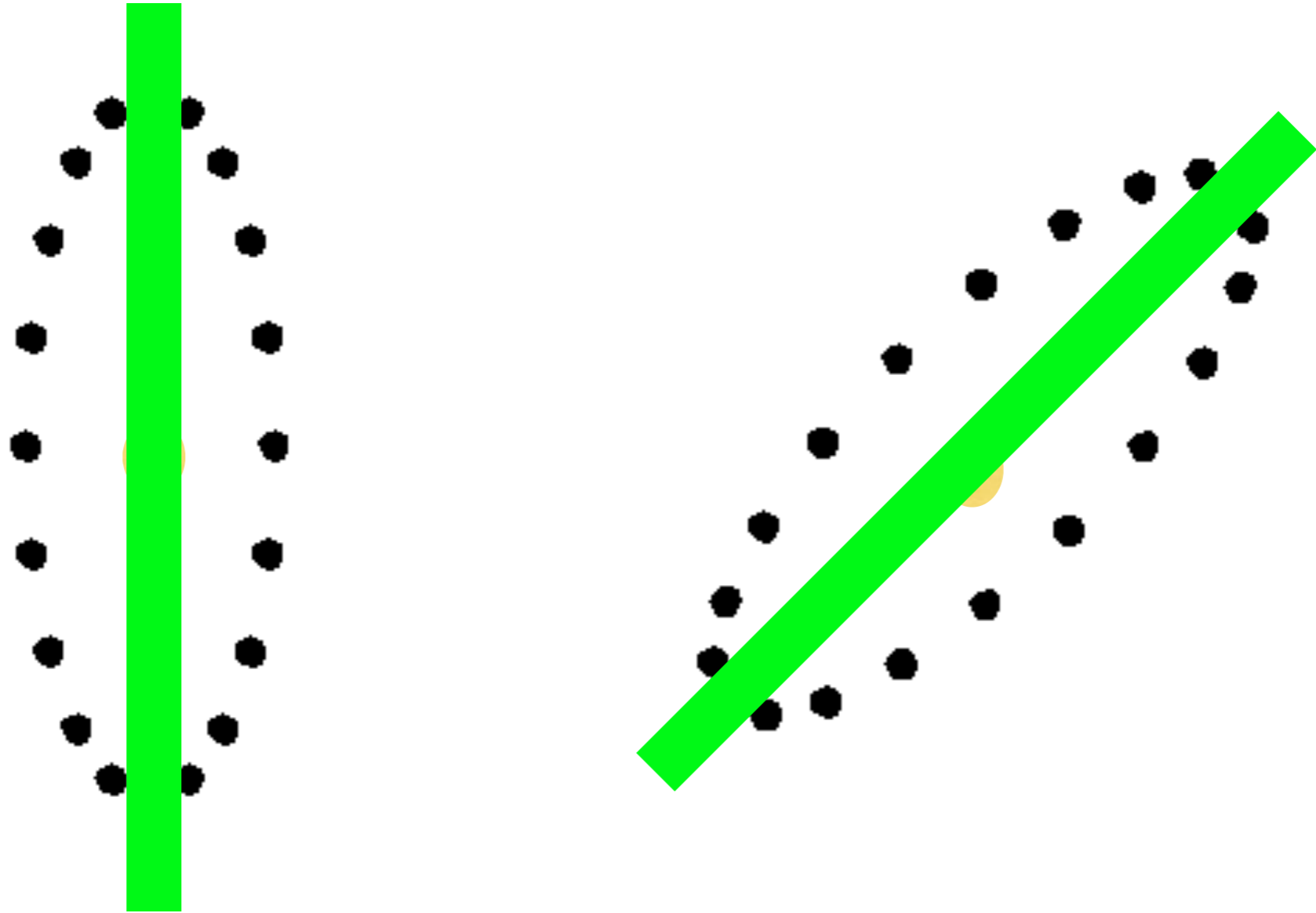


# From GW to CMB

## Polarization



# From GW to CMB Polarization



# “Tensor-to-scalar Ratio,” $r$

$$r \equiv \frac{2\langle |h_{\mathbf{k}}^+|^2 + |h_{\mathbf{k}}^\times|^2 \rangle}{\langle |\zeta_{\mathbf{k}}|^2 \rangle}$$

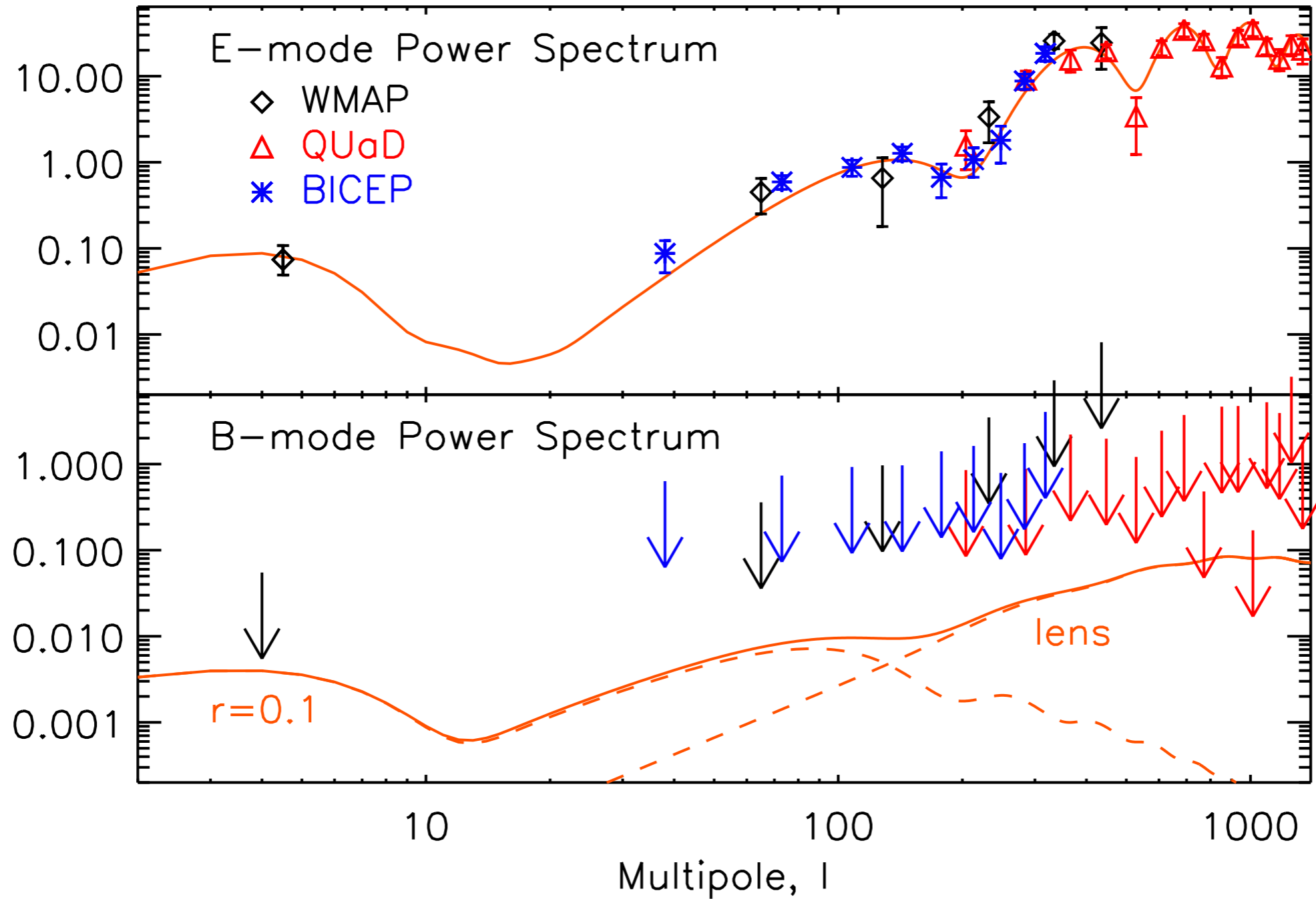
In terms of the slow-roll parameter:

$$r = 16\varepsilon$$

$$\begin{aligned} \text{where } \varepsilon &= -(\dot{H}/H^2) = 4\pi G(\dot{\varphi})^2/H^2 \\ &\approx (16\pi G)^{-1}(\dot{V}/V)^2 \end{aligned}$$

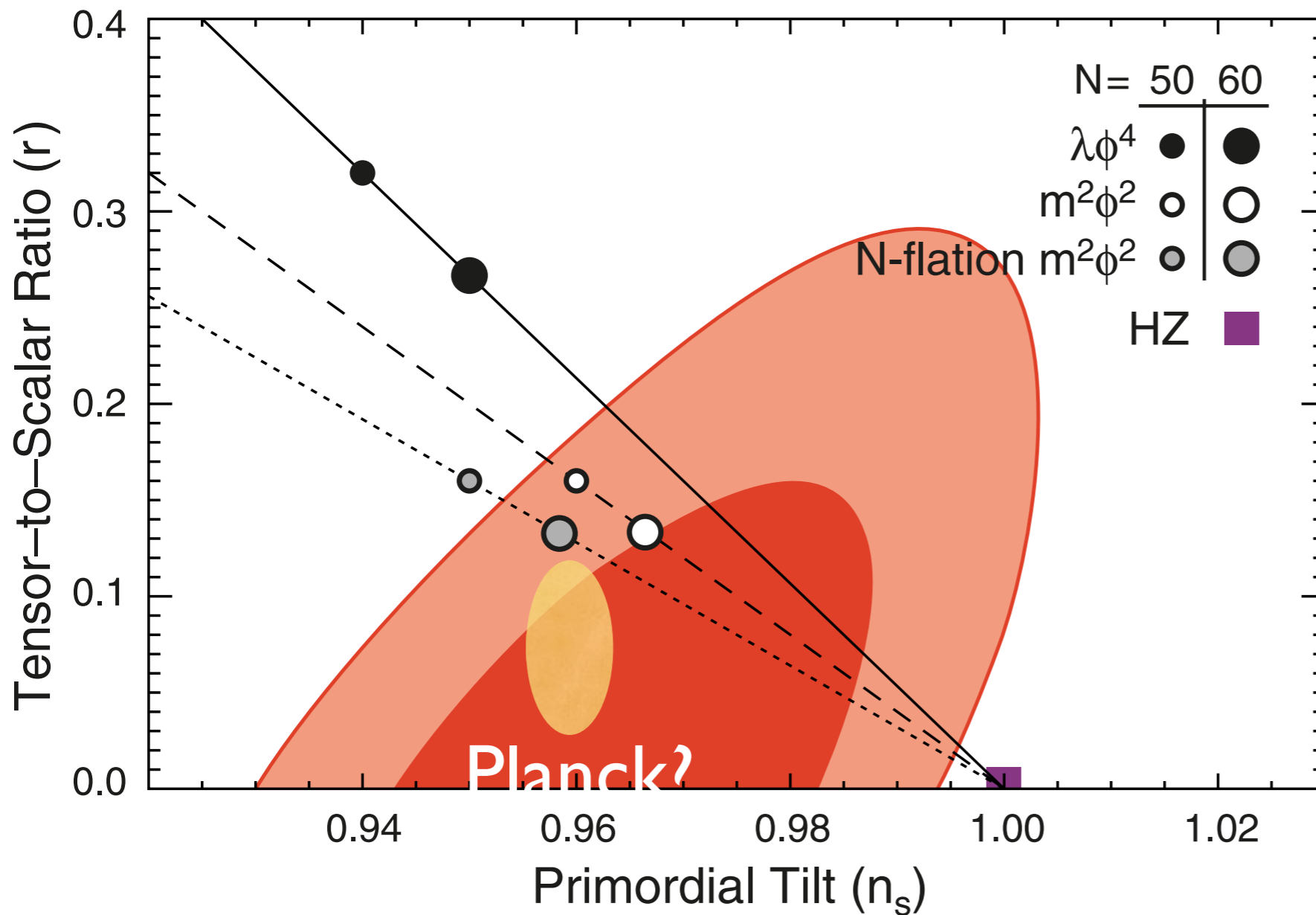
# Polarization Power Spectrum

from  $\zeta$



- No detection of polarization from gravitational waves (B-mode polarization) yet.

# Planck might find gravitational waves (if $r \sim 0.1$ )



If found, this would give us a pretty convincing proof that inflation did indeed happen.



# Bispectrum

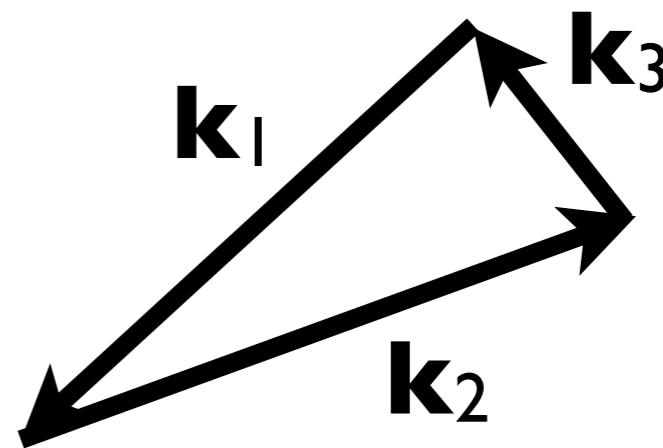
- Three-point function!

- $B_{\zeta}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3)$

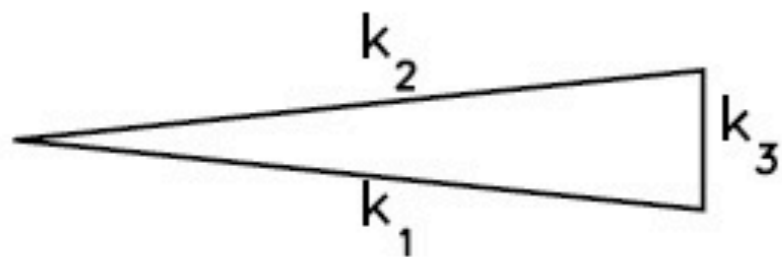
$$= \langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \zeta_{\mathbf{k}_3} \rangle = (\text{amplitude}) \times$$

$$(2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) F(k_1, k_2, k_3)$$

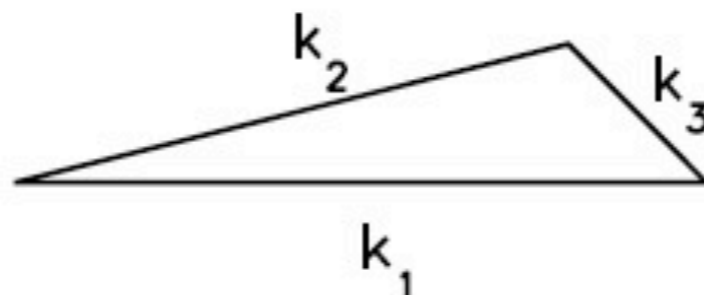
model-dependent function



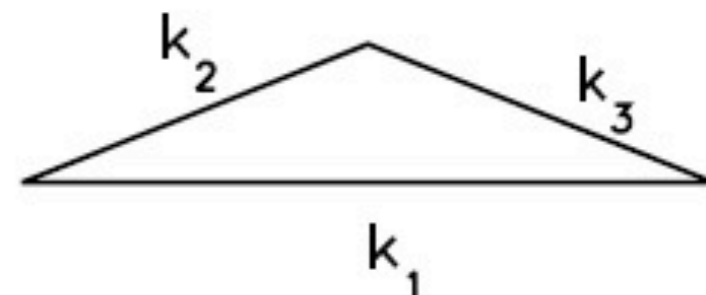
(a) squeezed triangle  
( $k_1 \approx k_2 \gg k_3$ )



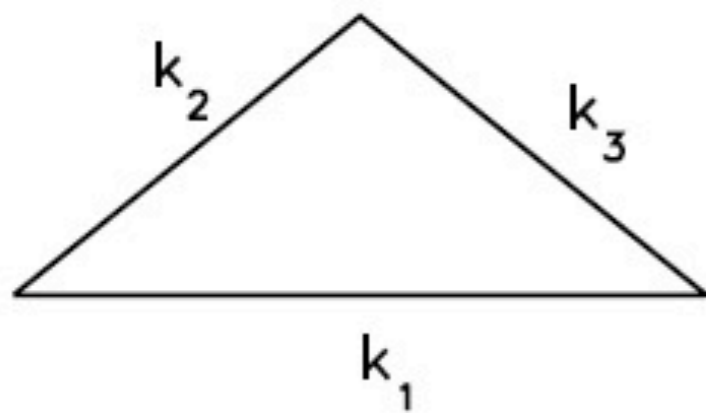
(b) elongated triangle  
( $k_1 = k_2 + k_3$ )



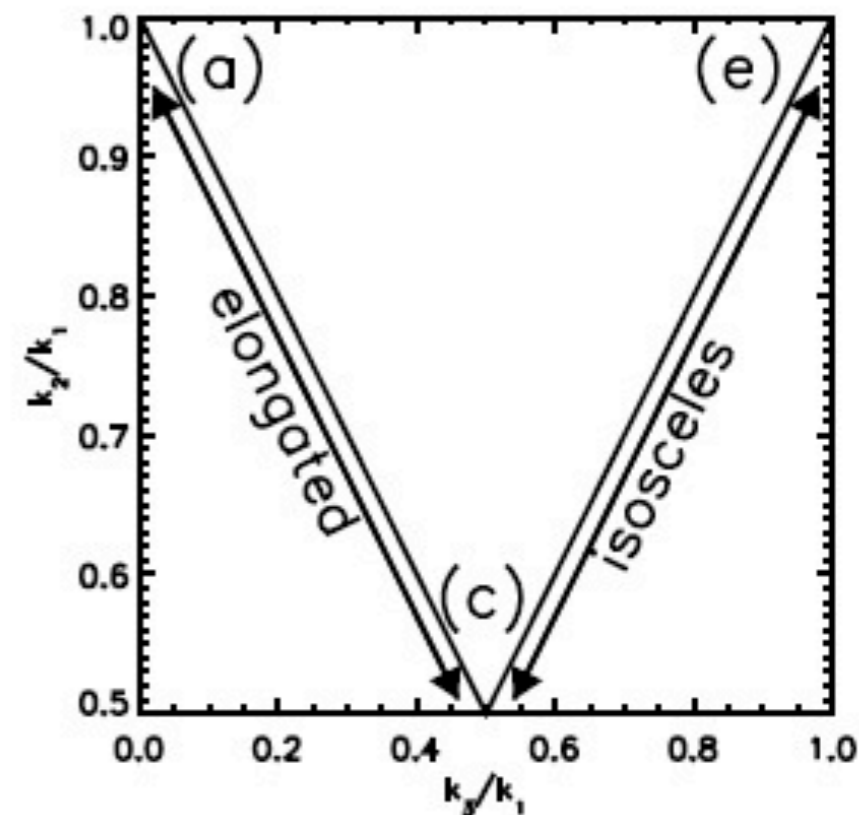
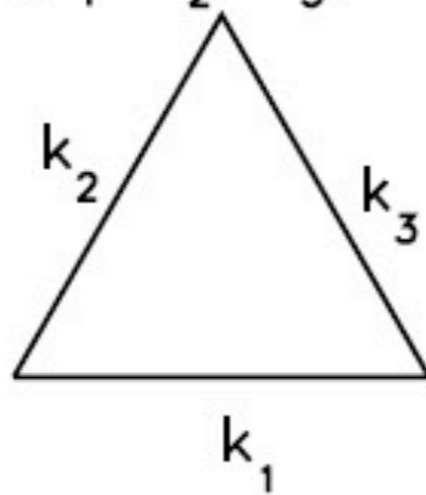
(c) folded triangle  
( $k_1 = 2k_2 = 2k_3$ )



(d) isosceles triangle  
( $k_1 > k_2 = k_3$ )



(e) equilateral triangle  
( $k_1 = k_2 = k_3$ )



**MOST IMPORTANT**

# Ruling-out Inflation (3-point Function)

- Inflation models predict that primordial fluctuations are very close to Gaussian.
- In fact, **ALL SINGLE-FIELD** models predict the squeezed-limit 3-point function to have the amplitude of  $f_{\text{NL}}=0.02$ .
- Detection of  $f_{\text{NL}} > 1$  would rule out **ALL** single-field models!

# Ruling-out Inflation (3-point Function)

- No detection of this form of 3-point function of primordial curvature perturbations. The 95% CL limit is:
  - $-10 < f_{\text{NL}} < 74$
  - 68%CL:  $f_{\text{NL}} = 32 \pm 21$
  - The WMAP data are consistent with the prediction of **simple single-field inflation** models:  $1 - n_s \approx r \approx f_{\text{NL}}$
- Planck will cut the error bar by a factor of four!

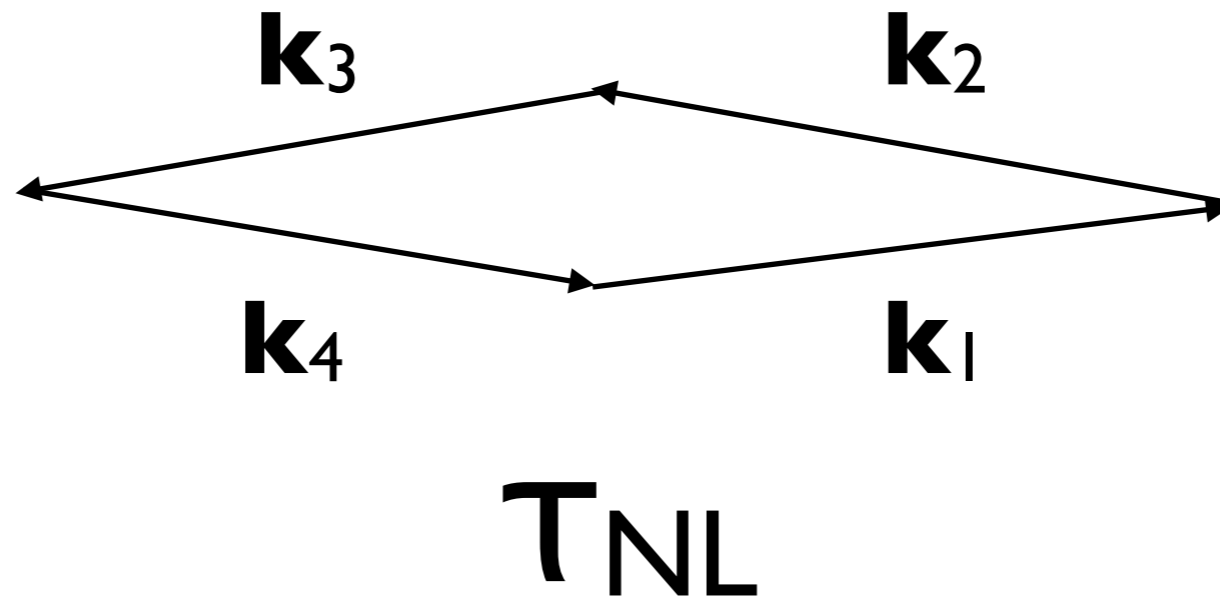
OK, for fun, let us suppose  
that single-field models are  
ruled out by Planck.

Now what?

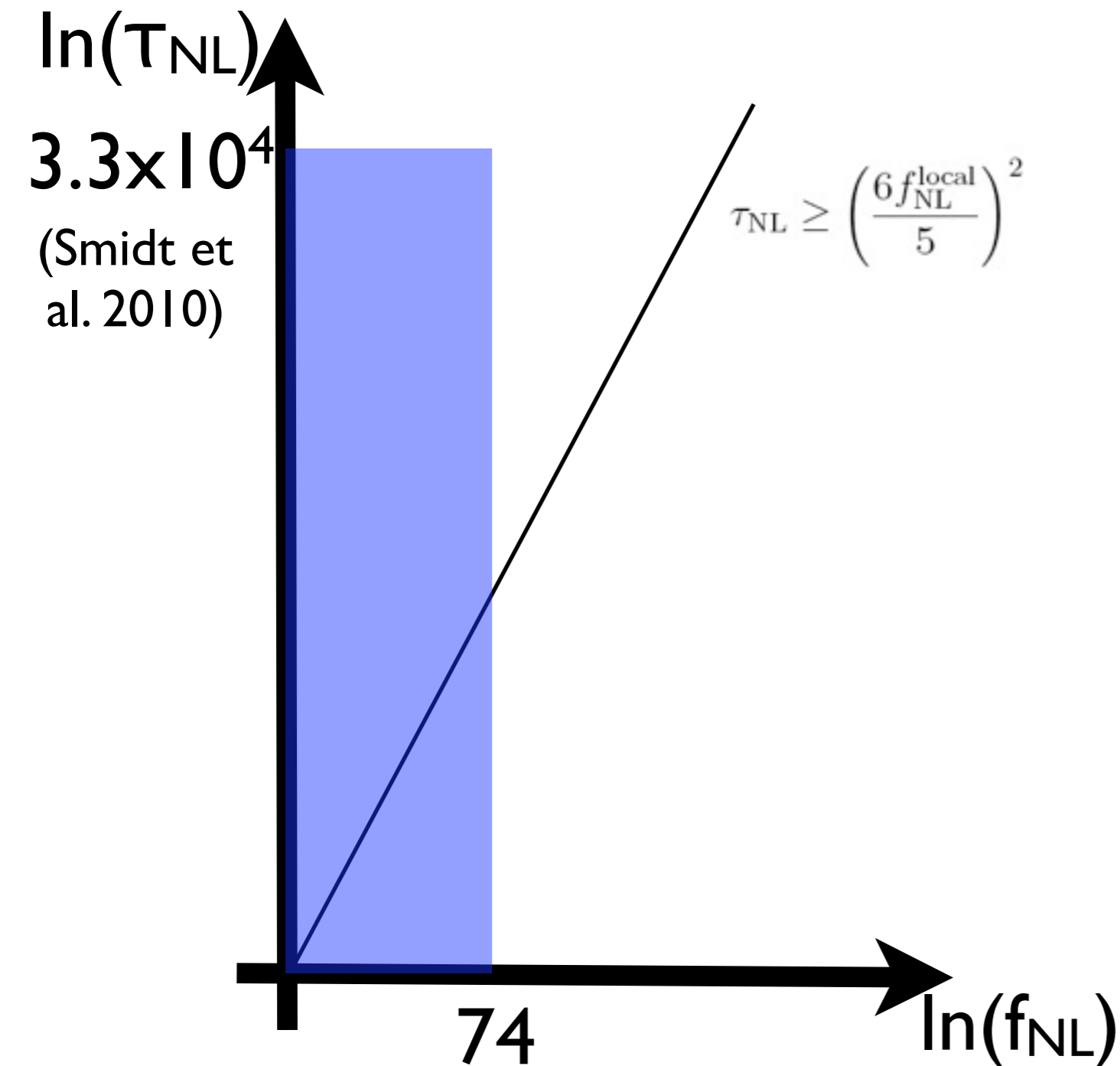
- We just don't want to be thrown into multi-field landscape without any clues...
- What else can we use?
  - Four-point function!

# Trispectrum

- $T_{\zeta}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3, \mathbf{k}_4) = (2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4)$   
 $+ T_{NL} [P_{\zeta}(k_1) P_{\zeta}(k_2) (P_{\zeta}(|\mathbf{k}_1 + \mathbf{k}_3|) + P_{\zeta}(|\mathbf{k}_1 + \mathbf{k}_4|))$   
 $+ \text{cyc.}]$

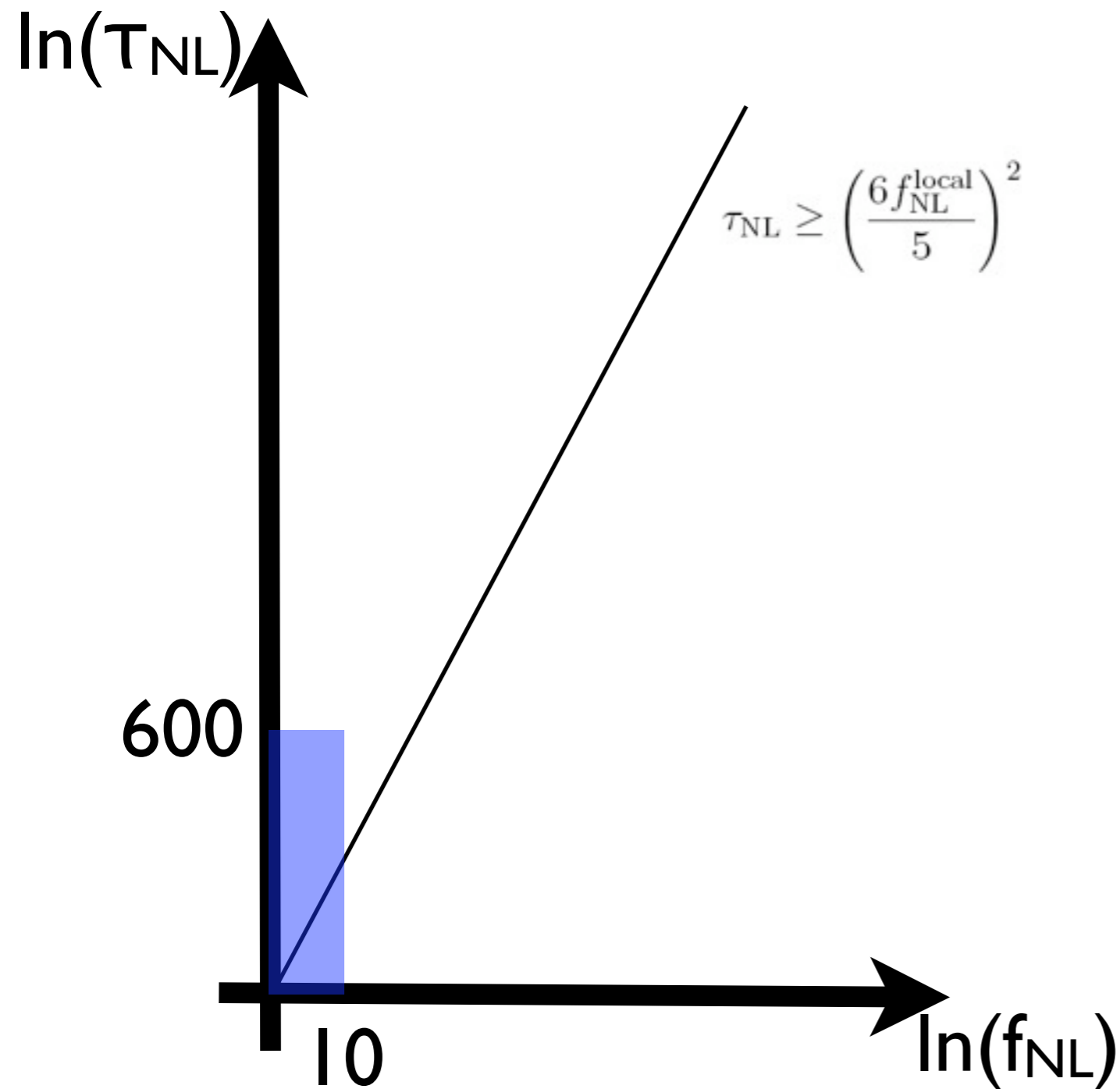


# Testing Inflation Paradigm



- The current limits from WMAP 7-year are consistent with single-field or multi-field models.
- So, let's play around with the future.

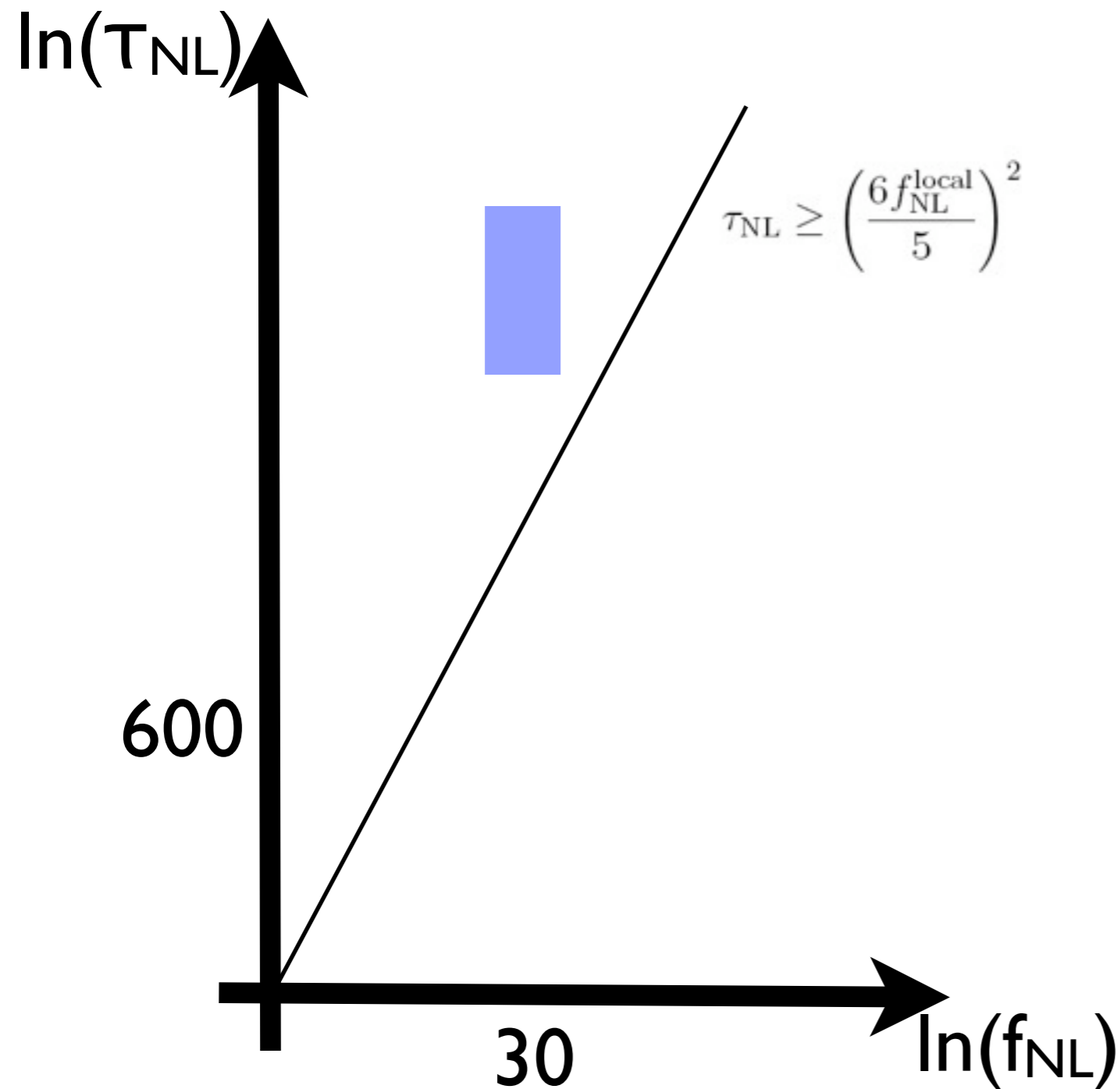
# Case A: Single-field Happiness



- No detection of anything after Planck. Single-field survived the test.



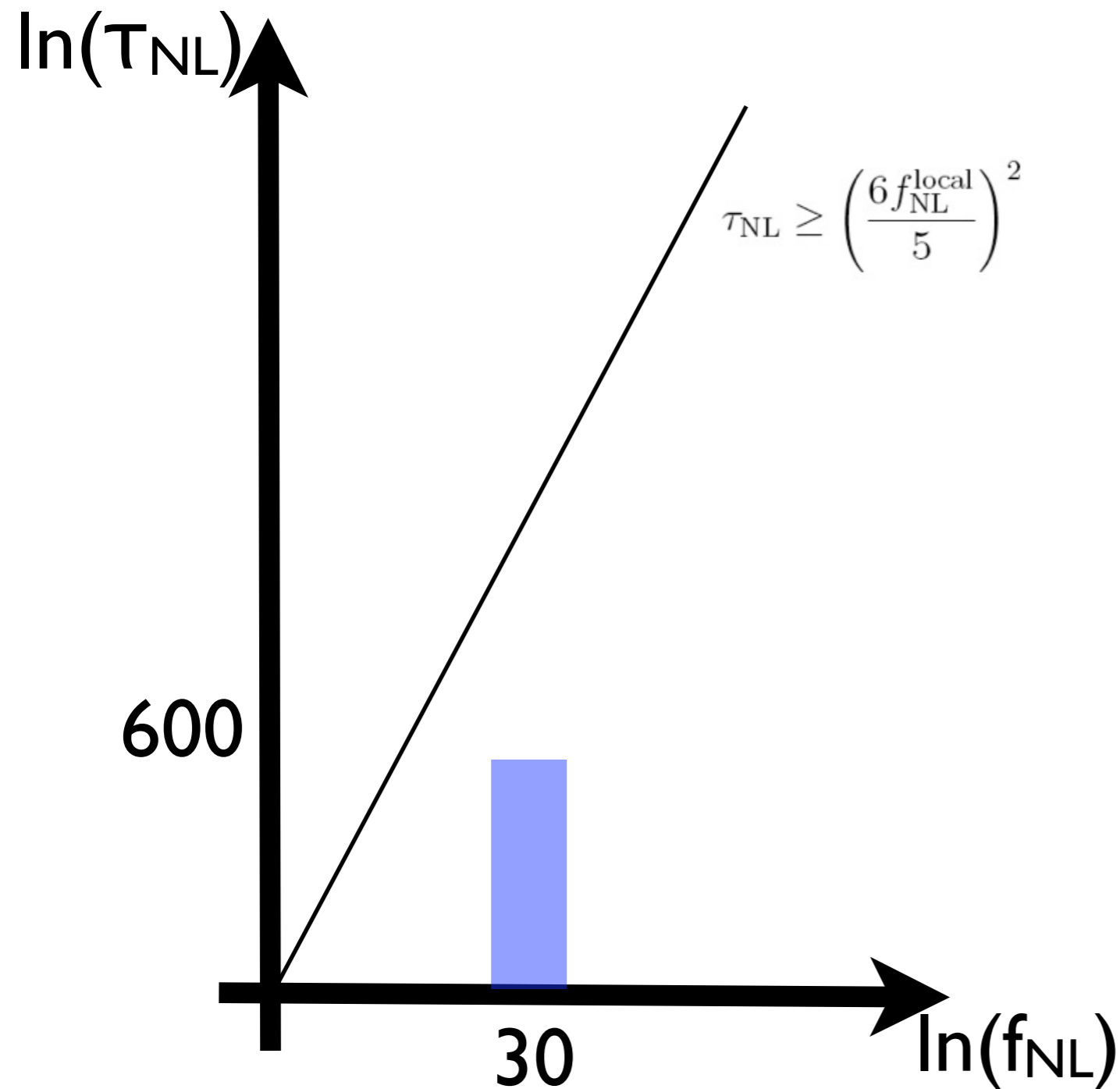
# Case B: Multi-field Happiness



- $f_{\text{NL}}$  is detected. Single-field is dead.
- But,  $\tau_{\text{NL}}$  is also detected, in accordance with the multi-field inflation relation

$$\tau_{\text{NL}} \geq \left(\frac{6f_{\text{NL}}^{\text{local}}}{5}\right)^2$$

# Case C: Madness



- $f_{\text{NL}}$  is detected. Single-field is dead.
- But,  $\tau_{\text{NL}}$  is **not** detected, inconsistent with  $\tau_{\text{NL}} \geq \left(\frac{6f_{\text{NL}}^{\text{local}}}{5}\right)^2$
- BOTH the single-field and multi-field are gone.

# What should you expect in the future?

- Please keep your eyes on:
  - **Year 2013**
    - $N_{\text{eff}}$ : is it 4? If it is 4, Planck will measure  $N_{\text{eff}}=4.0\pm 0.2$
    - $f_{\text{NL}}$ : is it 30? If it is 30, Planck will measure  $f_{\text{NL}}=30\pm 5$ . We should then do the 4-point function test.
  - **Year 2014**
    - $r$ : is it as large as 0.1? If it is 0.1, Planck will measure  $r=0.1\pm 0.05$ .