

12.6 Diagrammatic Expectation Values

$$\mathcal{H}(\phi) = \frac{1}{2}\phi^\dagger D^{-1}\phi - j^\dagger\phi + \frac{1}{4!}\lambda\phi^4$$

Logarithm of partition function:

$$\ln \mathcal{Z}(j) = \ln \mathcal{Z}_G(0) + \dots + \begin{array}{c} \bullet \\ | \\ | \\ | \\ \bullet \end{array} + \begin{array}{c} \bullet \\ \nearrow \\ \circlearrowleft \\ \circlearrowright \\ \bullet \end{array} + \begin{array}{c} \bullet \\ \circlearrowleft \\ \circlearrowright \\ \bullet \end{array} + \mathcal{O}(\lambda^2)$$

$$\Rightarrow \ln \mathcal{Z}(j) = \ln \mathcal{Z}_G(0) + \frac{1}{2}j^\dagger Dj - \frac{\lambda^\dagger}{4!}(Dj)^4 - \frac{1}{4}\lambda^\dagger(Dj)^2\widehat{D} - \frac{1}{8}\lambda^\dagger\widehat{D}^2 + \mathcal{O}(\lambda^2)$$

12.6 Diagrammatic Expectation Values

$$\ln \mathcal{Z}(j) = \ln \mathcal{Z}_G(0) + \frac{1}{2} j^\dagger D j - \frac{\lambda^\dagger}{4!} (D j)^4 - \frac{1}{4} \lambda^\dagger (D j)^2 \hat{D} - \frac{1}{8} \lambda^\dagger \hat{D}^2 + \mathcal{O}(\lambda^2)$$

Expectation value:

$$\langle \phi \rangle = \frac{\delta \ln \mathcal{Z}}{\delta j}$$

$$= D j - \frac{\lambda}{3!} D(D j)^3 - \frac{1}{2} D \lambda (D j) \hat{D} + \mathcal{O}(\lambda^2)$$

$$\langle \phi^x \rangle = D^{xy} j_y - \frac{\lambda}{3!} \int dy D^{yx} (D^{yz} j_z)^3 - \frac{\lambda}{2} \int dy D^{yx} (D^{yz} j_z) D^{yy} + \mathcal{O}(\lambda^2)$$

$$\langle \phi^x \rangle = \text{---} \bullet + \text{---} \begin{array}{c} \bullet \\ | \\ | \\ \bullet \end{array} + \text{---} \circlearrowleft \text{---} \circlearrowright + \mathcal{O}(\lambda^2)$$

12.6 Diagrammatic Expectation Values

$$\ln \mathcal{Z}(j) = \ln \mathcal{Z}_G(0) + \frac{1}{2} j^\dagger D j - \frac{\lambda^\dagger}{4!} (D j)^4 - \frac{1}{4} \lambda^\dagger (D j)^2 \hat{D} - \frac{1}{8} \lambda^\dagger \hat{D}^2 + \mathcal{O}(\lambda^2)$$

Covariance:

$$\begin{aligned}\langle \phi \phi^\dagger \rangle^c &= \frac{\delta \ln \mathcal{Z}}{\delta j \delta j^\dagger} \\ &= D - \frac{\lambda}{2} D (D j)^2 D - \frac{\lambda}{2} D (\hat{D}) D + 0 + \mathcal{O}(\lambda^2) \\ \langle \phi^x \phi^y \rangle^c &= D^{xy} - \frac{\lambda}{2} \int dz D^{zx} (D^{zu} j_u)^2 D^{zy} - \frac{\lambda}{2} \int dz D^{xz} D^{zz} D^{zy} + \mathcal{O}(\lambda^2)\end{aligned}$$

$$\langle \phi^x \phi^y \rangle^c = \text{---} + \text{---} + \text{---} + \mathcal{O}(\lambda^2)$$

12.7 Log-Normal Poisson Model Diagrammatically

Log-normal Poisson model (*repetition*)

$$\begin{aligned}\mathcal{H}(d, s) &\stackrel{\cong}{=} \frac{1}{2}s^\dagger S^{-1}s - d^\dagger s + \kappa^\dagger e^s \\ &\stackrel{\cong}{=} \frac{1}{2}s_x((S^{-1})^{xy} + \kappa\delta^{xy})s^y - (\textcolor{red}{d}_x - \kappa_x)s^x + \kappa_x \sum_{n=3}^{\infty} \frac{1}{n!}(s^n)^x \\ &\stackrel{\cong}{=} \frac{1}{2}s^\dagger(S^{-1} + \widehat{\kappa})s - \textcolor{red}{j}^\dagger s + \kappa^\dagger \sum_{n=3}^{\infty} \frac{1}{n!}s^n \\ &\stackrel{\cong}{=} \frac{1}{2}s^\dagger D^{-1}s - j^\dagger s + \sum_{n=3}^{\infty} \frac{1}{n!} \int dx_1 \dots dx_n \kappa_{x_1} \delta(x_1 - x_2) \dots \delta(x_1 - x_n) s^{x_1} \dots s^{x_n} \\ &\stackrel{\cong}{=} \frac{1}{2}s^\dagger D^{-1}s - j^\dagger s + \sum_{n=0}^{\infty} \frac{1}{n!} \Lambda_{x_1 \dots x_n}^{(n)} s^{x_1} \dots s^{x_n}\end{aligned}$$

12.7 Log-Normal Poisson Model Diagrammatically

$$\mathcal{H}(d, s) \quad \hat{=} \quad \frac{1}{2} s^\dagger D^{-1} s - j^\dagger s + \kappa^\dagger \sum_{n=3}^{\infty} \frac{1}{n!} s^n$$

$$\begin{aligned}\textbf{MAP: } \frac{\delta H}{\delta s} \Big|_{s=m} &\stackrel{!}{=} 0 \\ &= D^{-1} s - j + \kappa \sum_{n=3}^{\infty} \frac{s^{n-1}}{(n-1)!} \Big|_{s=m} \\ \Rightarrow m &= D \left(j - \kappa \sum_{n=2}^{\infty} \frac{m^n}{n!} \right)\end{aligned}$$

Iteration:

- ▶ $i = 0: m_0 = 0$
- ▶ $i = 1: \quad m_1 = Dj = \text{_____} \bullet$

12.7 Log-Normal Poisson Model Diagrammatically

► $i = 2 :$

$$m_2 = D \left(j - \kappa \sum_{n=2}^{\infty} \frac{(Dj)^n}{n!} \right) = \underbrace{\text{---}}_{n=1} \bullet + \underbrace{\text{---} \bullet \begin{array}{c} \bullet \\ \backslash \end{array}}_{n=2} + \underbrace{\text{---} \bullet \begin{array}{c} \bullet \\ | \\ \bullet \end{array}}_{n=3} + \dots$$

► $i = 3$

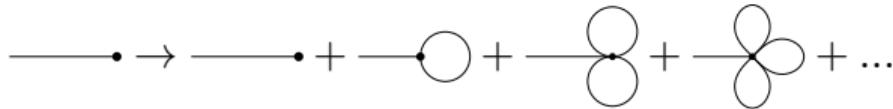
$$\begin{aligned} m_3 = D \left(j - \kappa \sum_{n=2}^{\infty} \frac{(m_2)^n}{n!} \right) &= \text{---} \bullet + \text{---} \bullet \begin{array}{c} \bullet \\ \backslash \end{array} + \text{---} \bullet \begin{array}{c} \bullet \\ | \\ \bullet \end{array} + \dots \\ &+ \text{---} \bullet \begin{array}{c} \bullet \\ / \end{array} + \text{---} \bullet \begin{array}{c} \bullet \\ / \\ \bullet \end{array} + \text{---} \bullet \begin{array}{c} \bullet \\ | \\ / \end{array} + \dots \end{aligned}$$

⇒ The classical/ MAP estimate m_∞ is always given by the sum of all tree diagrams with one external point.

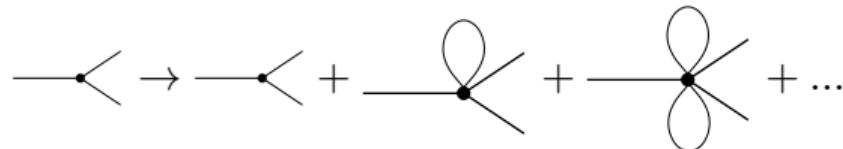
12.7.1 Consideration of uncertainty loops

$$\langle s \rangle_{(s|d)} = \underbrace{\sum \text{tree diagrams}}_{\text{MAP}} + \underbrace{\sum \text{loop diagrams}}_{\text{uncertainty corrections}}$$

- ▶ source loops:

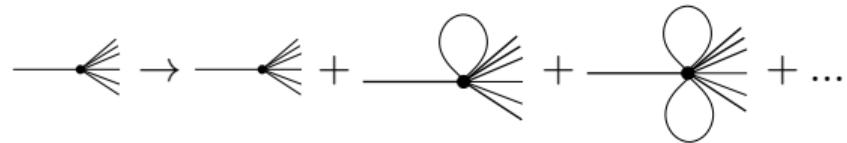


- ▶ 3-vertex:



12.7.1 Consideration of uncertainty loops

► n-vertex:



$$\Rightarrow -\kappa_x \rightarrow - \left[\kappa_x + \frac{1}{2} \kappa_x D^{xx} + \frac{1}{8} \kappa_x (D^{xx})^2 + \dots + \frac{1}{n!2^n} \kappa_x (\hat{D}^x)^n \right] \text{ no summation}$$
$$-\kappa \rightarrow -\kappa e^{\hat{D}/2}$$

12.7.1 Consideration of uncertainty loops

Loop-normalized MAP:

$$\begin{aligned} m &= S(d - \kappa e^m) \\ &= S(d - \kappa_m) \end{aligned}$$

$$\kappa_m \rightarrow \kappa_m e^{\hat{D}/2} = \kappa_{m+\hat{D}/2}:$$

loop normalized solution:

$$\begin{aligned} m &= S \left(d - \kappa_{m+\hat{D}/2} \right) \\ D &= \left(S^{-1} + \widehat{\kappa}_{m+\hat{D}/2} \right) \end{aligned}$$

13 Thermodynamical Inference

Tempered posterior: $T = \frac{1}{\beta}$

$$\begin{aligned}\mathcal{P}(s|d, T, J) &= \frac{e^{-\beta(\mathcal{H}(d, s) + J^\dagger s)}}{\mathcal{Z}(d, \beta, J)} \\ &= \underbrace{\frac{\left(\mathcal{P}(d, s)e^{-J^\dagger s}\right)^\beta}{\int \mathcal{D}s \left(\mathcal{P}(d, s)e^{-J^\dagger s}\right)^\beta}}_{=\mathcal{Z}(d, \beta, J)}\end{aligned}$$

- ▶ $T = \beta = 1$: usual inference
- ▶ $T \rightarrow 0, \beta \rightarrow \infty$: enlarged contrast $\Rightarrow \mathcal{P}(s|d, T) \rightarrow \delta(s - s_{\text{MAP}})$
- ▶ $T \rightarrow \infty, \beta \rightarrow 0$: weaker contrast $\Rightarrow \mathcal{P}(s|d, T) \rightarrow \text{const}$

Boltzmann Entropy

$$\begin{aligned}\mathcal{P}(s|d, T, J) &= \frac{e^{-\beta(\mathcal{H}(d, s) + J^\dagger s)}}{\mathcal{Z}(d, \beta, J)} \\ S_B &= - \int \mathcal{D}s \mathcal{P}(s|d, T, J) \ln \left(\frac{\mathcal{P}(s|d, T, J)}{\text{const}} \right) \\ \Delta S_B &= \int \mathcal{D}s \mathcal{P}(s|d, T, J) \left[\beta (\mathcal{H}(d, s) + J^\dagger s) + \ln \mathcal{Z}(d, \beta, J) \right] \\ &= \beta \left[\langle \mathcal{H}(d, s) \rangle_{(s|d, T, J)} + J^\dagger \langle s \rangle_{(s|d, T, J)} + \frac{1}{\beta} \ln \mathcal{Z}(d, \beta, J) \right] \\ &= \beta [U(d, T, J) + J^\dagger m(d, T, J) - F(d, \beta, J)]\end{aligned}$$

Boltzmann Entropy

$$T\Delta S_B = U(d, T, J) + J^\dagger m(d, T, J) - F(d, \beta, J)$$

- ▶ internal energy: $U(d, T, J) = \langle \mathcal{H}(d, s) \rangle_{(s|d, T, J)}$
- ▶ Helmholtz free energy: $F(d, \beta, J) = -\frac{1}{\beta} \ln \mathcal{Z}(d, \beta, J)$
- ▶ mean field:

$$m(d, 1, J) = \langle s \rangle_{(s|d, 1, J)} = -\frac{\partial}{\partial J} \ln \mathcal{Z}(d, \beta = 1, J) \Big|_{J=0} = \frac{\partial F(d, \beta, J)}{\partial J} \Big|_{J=0, \beta=1}$$

13. Thermodynamical Inference

Ansatz:

$$\begin{aligned}\mathcal{P}(s|d, T, J) &\approx \tilde{\mathcal{P}}(s|m, D) = \mathcal{G}(s - m, D) \\ T\Delta \tilde{S}_B(d, T, J) &= \tilde{U}(d, T, J) + J^\dagger m(d, T, J) - \tilde{F}(d, \beta, J) \\ \tilde{U}(d, T, J) &= \langle \mathcal{H}(d, s) \rangle_{\mathcal{G}(s-m, D)}\end{aligned}$$

$$\begin{aligned}\Rightarrow \tilde{S}_B(d, T, J) &= -\langle \ln \tilde{\mathcal{P}} \rangle_{\tilde{\mathcal{P}}} \\ &= + \int \mathcal{D}s \mathcal{G}(\textcolor{blue}{s} - \textcolor{blue}{m}, D) \left[\frac{1}{2} (\textcolor{blue}{s} - \textcolor{blue}{m})^\dagger D^{-1} (\textcolor{blue}{s} - \textcolor{blue}{m}) + \frac{1}{2} \ln |2\pi D| \right] \\ &= \frac{1}{2} \left[\int \mathcal{D}\varphi \left(\mathcal{G}(\varphi, D) \operatorname{Tr}(\varphi \varphi^\dagger D^{-1}) \right) + \ln |2\pi D| \right]\end{aligned}$$

13. Thermodynamic Inference

$$\begin{aligned}\Rightarrow \tilde{S}_B(d, T, J) &= \frac{1}{2} \left[\int \mathcal{D}\varphi \left(\mathcal{G}(\varphi, D) \text{Tr}(\varphi\varphi^\dagger D^{-1}) \right) + \ln |2\pi D| \right] \\ &= \frac{1}{2} \text{Tr} \left(\langle \varphi\varphi^\dagger \rangle_{\mathcal{G}(\varphi, D)} D^{-1} \right) + \frac{1}{2} \ln |2\pi D| \\ &= \frac{1}{2} \text{Tr}(DD^{-1}) + \frac{1}{2} \text{Tr}(\ln(2\pi D)) \\ &= \frac{1}{2} \text{Tr}(\mathbb{1} + \ln(2\pi D)) \\ &= \tilde{S}_B(D)\end{aligned}$$

$$\Rightarrow \tilde{F}(d, \beta, J) = \tilde{U}(m_J, D_J) - \tilde{T} \tilde{S}_B(D_J) + J^\dagger m_J$$

Legendre Transformation

$$F(J) = F(J_0) + \frac{\partial F}{\partial J} \Big|_{J_0}^\dagger (J - J_0) + \dots$$

$$G = F(J_0) - \frac{\partial F}{\partial J} \Big|_{J_0}^\dagger J_0$$

F convex $\Rightarrow m_J = \frac{\partial F}{\partial J} \Rightarrow F$ can be reconstructed from $G(m)$

Gibbs Free Energy

$$\begin{aligned} G &= F - \frac{\partial F^\dagger}{\partial J} J \\ &= U - TS_B + J^\dagger m - J^\dagger m \\ \Rightarrow \tilde{G}(d, \beta, m, D) &= \tilde{U}(d, \beta, m, D) - T \tilde{S}_B(D) \end{aligned}$$

Mean Field From Minimal Gibbs Free Energy:

$$\frac{\delta G(d, m, D)}{\delta m} = 0 \quad \Rightarrow \quad m = \langle s \rangle_{(s|d)} \Big|_{T=1}$$

Proof:

$$\begin{aligned}\frac{\delta G}{\delta m} &= \frac{\delta}{\delta m} \left(F(d, J(m)) - J^\dagger(m)m \right) \\ &= \frac{\delta J(m)}{\delta m}^\dagger \frac{\delta F(d, J)}{\delta J} - \frac{\delta J^\dagger}{\delta m}m - J \\ &= \frac{\delta J(m)}{\delta m}^\dagger \textcolor{blue}{m} - \frac{\delta J^\dagger}{\delta m}m - J \\ &= -J \stackrel{!}{=} 0 \\ J = 0 \quad \Rightarrow \quad m &= \frac{\partial F}{\partial J} \Big|_{J=0} = \langle s \rangle_{(s|d)}\end{aligned}$$

Uncertainty Dispersion From Minimal Gibbs Free Energy:

$$\left(\frac{\delta^2 G}{\delta m \delta m^\dagger} \right)^{-1} \Big|_{m=\langle s \rangle_{(s|d)}} = \left. \frac{-\delta^2 F}{\delta J \delta J^\dagger} \right|_{J=0} = \beta D$$

Proof:

$$\begin{aligned} \left(\frac{\delta^2 G}{\delta m \delta m^\dagger} \right)^{-1} \Big|_{m=\langle s \rangle_{(s|d)}} &= \left(-\frac{\delta J}{\delta m} \right)^{-1} \Big|_{m=\langle s \rangle_{(s|d)}} \\ &= - \left(\frac{\delta m(J)}{\delta J} \right) \Big|_{J=0} \\ &= - \left. \frac{\delta^2 F(J)}{\delta J \delta J^\dagger} \right|_{J=0} \\ &= \frac{1}{\beta} \frac{\delta^2}{\delta J \delta J^\dagger} \ln \mathcal{Z}(d, J, \beta) \\ &= \beta D \end{aligned}$$